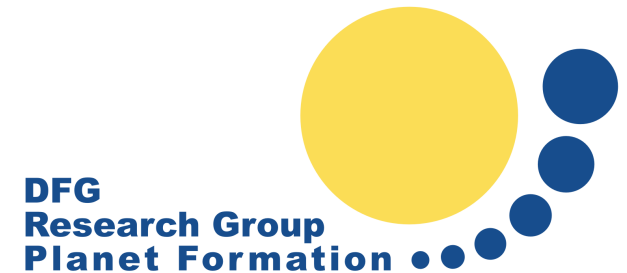


The Formation and orbital Evolution of Planets

Wilhelm Kley
Institut für Astronomie & Astrophysik
& Kepler Center for Astro and Particle Physics Tübingen

EBERHARD KARLS
UNIVERSITÄT
TÜBINGEN

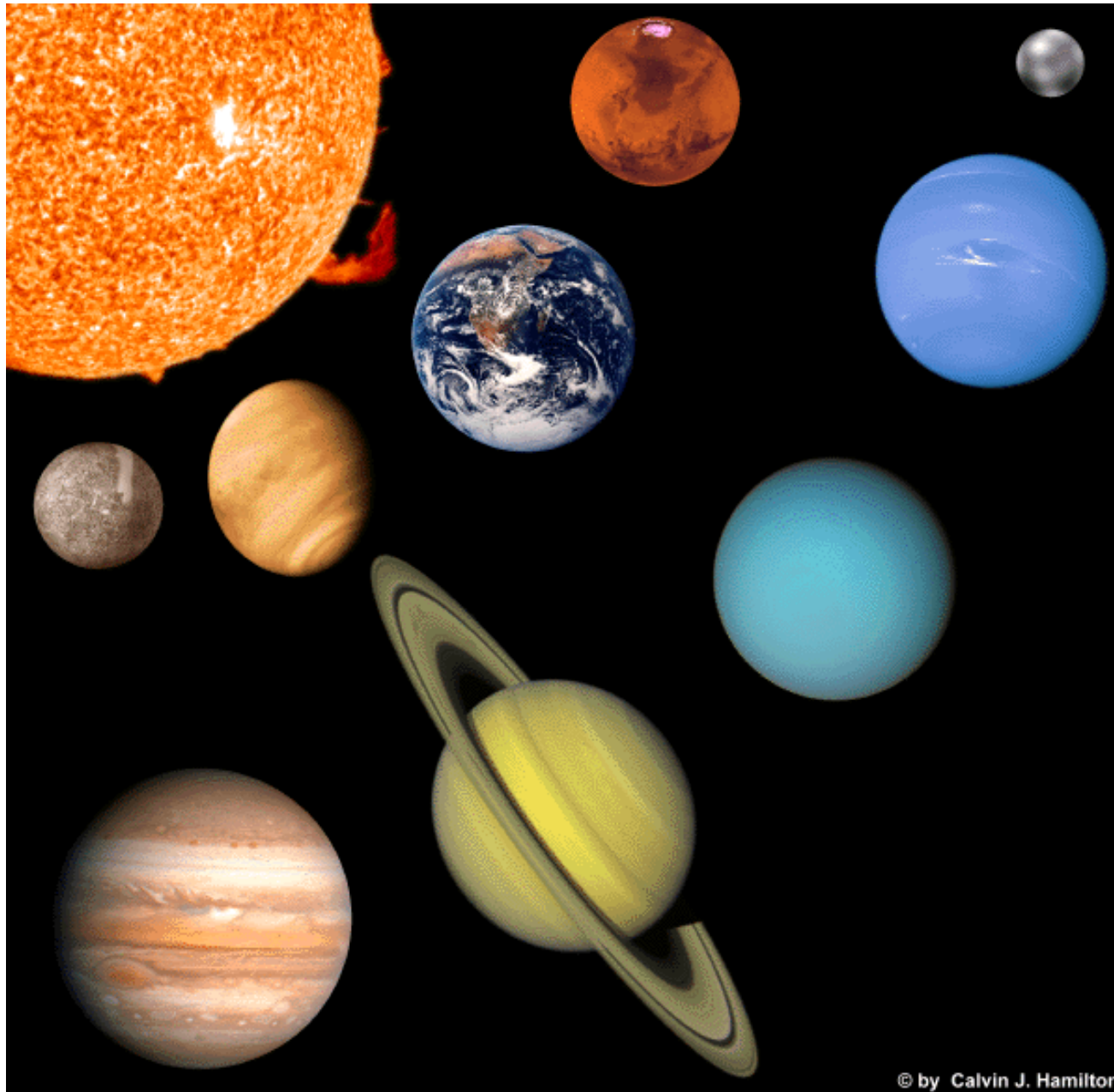


5. September, 2012



- The Solar System
 - Characteristics
 - Formation
- Extrasolar Planets
 - Planet-disk interaction
 - Dynamical evolution
- Gravitational instability
- Planet Formation in binaries
- Summary

(A. Crida)



Sun

Mercury

Venus

Earth

Mars

Jupiter

Saturn

Uranus

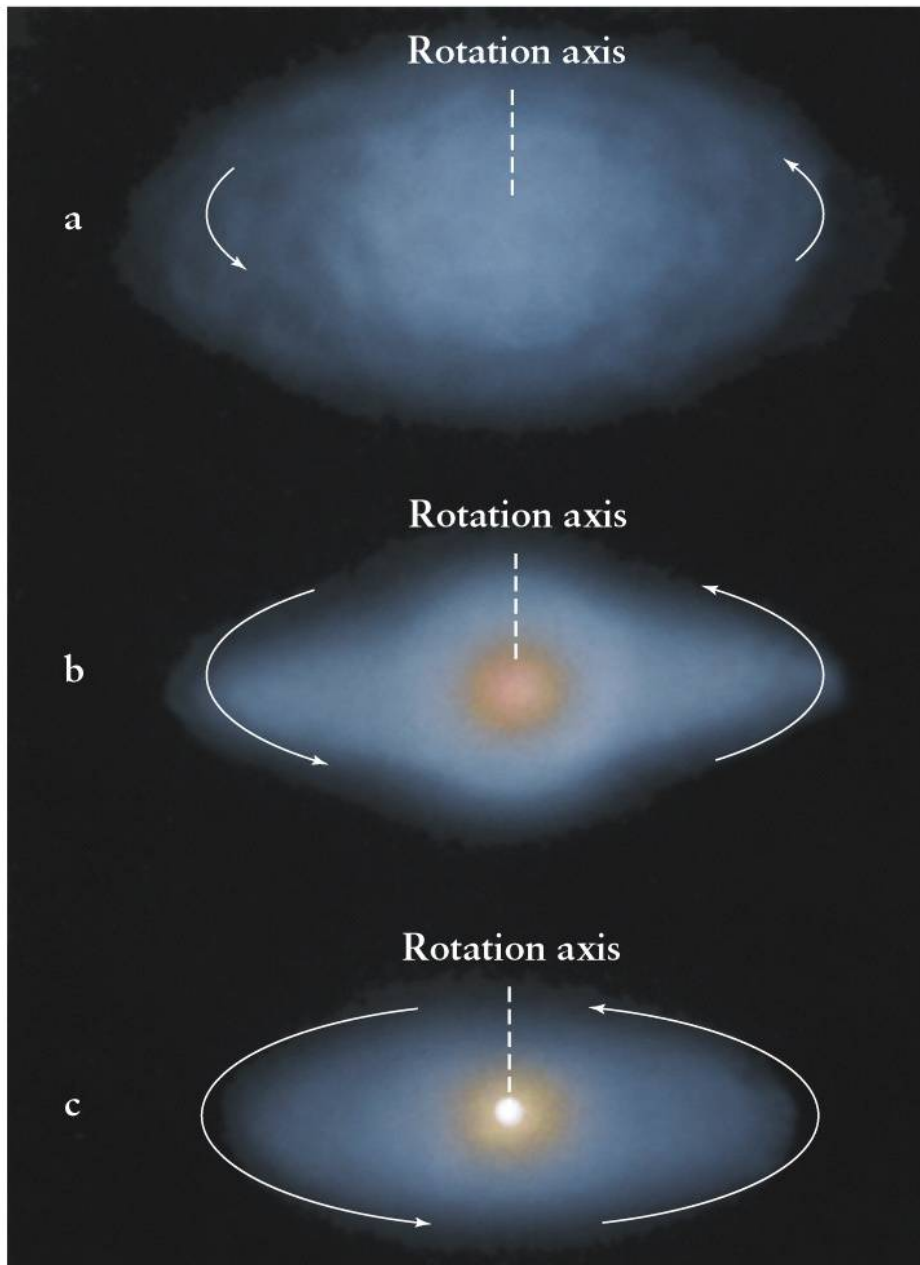
Neptun

(Pluto)



8 Planets:	Mercury to Neptune
5 Dwarf Planets:	Ceres, Pluto, Eris, Makemake, Haumea
Minor bodies:	TNO, asteroids, comets
Tiny bodies:	meteorites, dust

- coplanar, circular, uniform orbits (cf. Kepler candidates)
- Solid and gaseous planets (with Cores)
- prograde rotation (with exceptions)
- 99% of mass in Sun
- 99% of angular momentum in planets
- Age: about 4.5 billion years



Historic View:

(Leukippos, 480-420 BC)

“The worlds form in such a way, that the bodies sink into the empty space and connect to each other.”

Modern View:

Collaps of an interstellar
Molecular Cloud

Slight rotation \Rightarrow Flattening

Protosun in center / disk formation
(based on Kant & Laplace, 1750s)

Planets form in protoplanetary disks

\equiv Accretion Disks (99% Gas, 1% Dust)



Flat system, uniform rotation, circular orbits

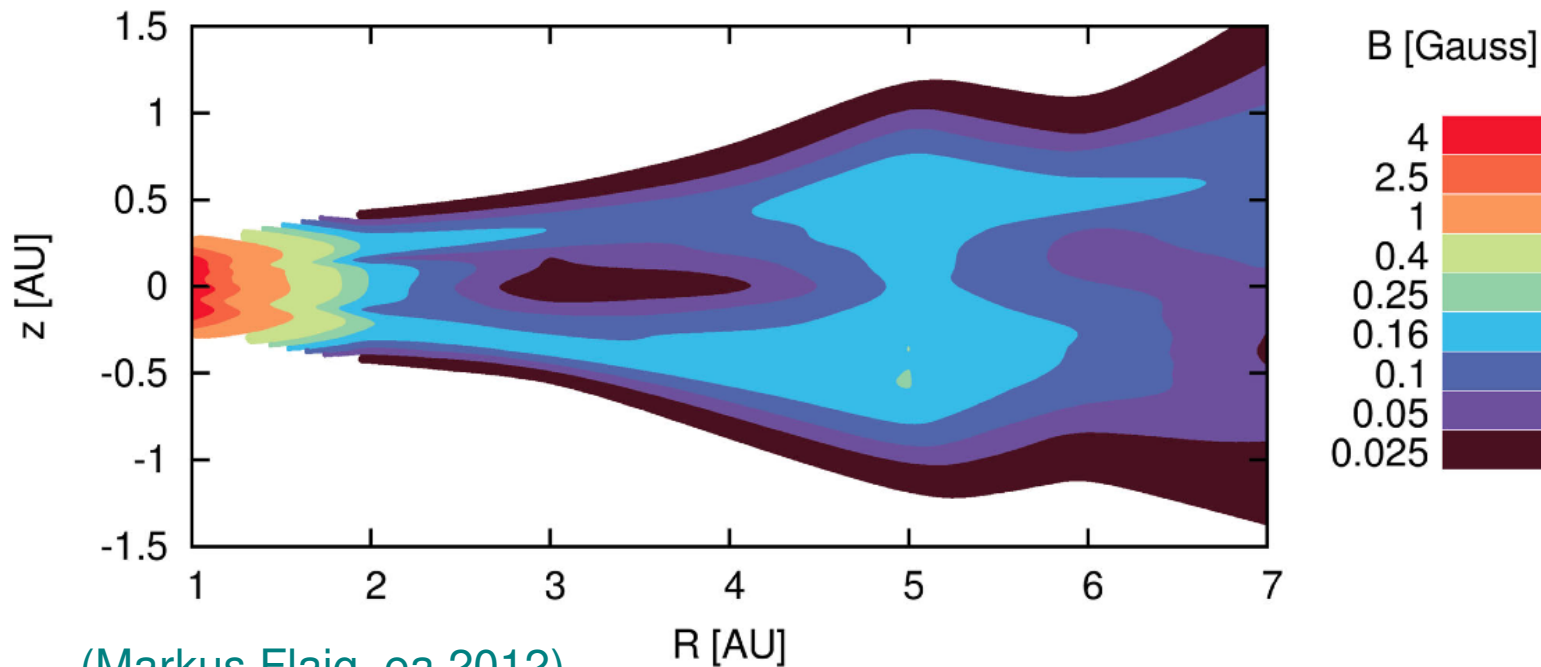
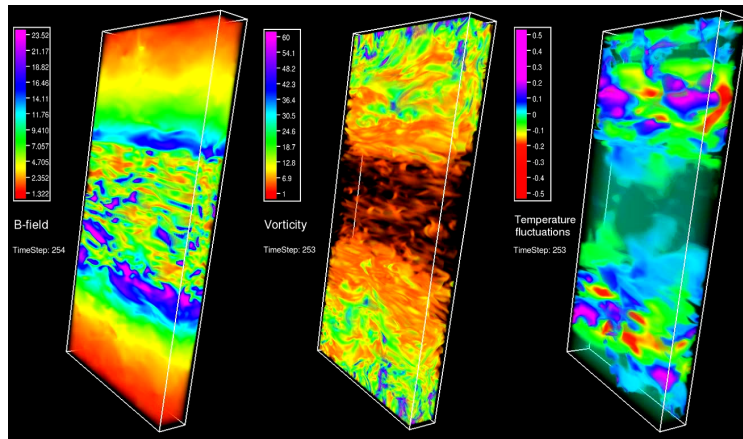


3D MHD Turbulence with **Radiation Transport and chemical network**
 in Accretion disks: **Stratified Local Shearing Box**, (Movie: 6 Orbits)

Outcome: (Talk: J.Simon, N. Turner

Poster: M. Flock, M. Hanasz)

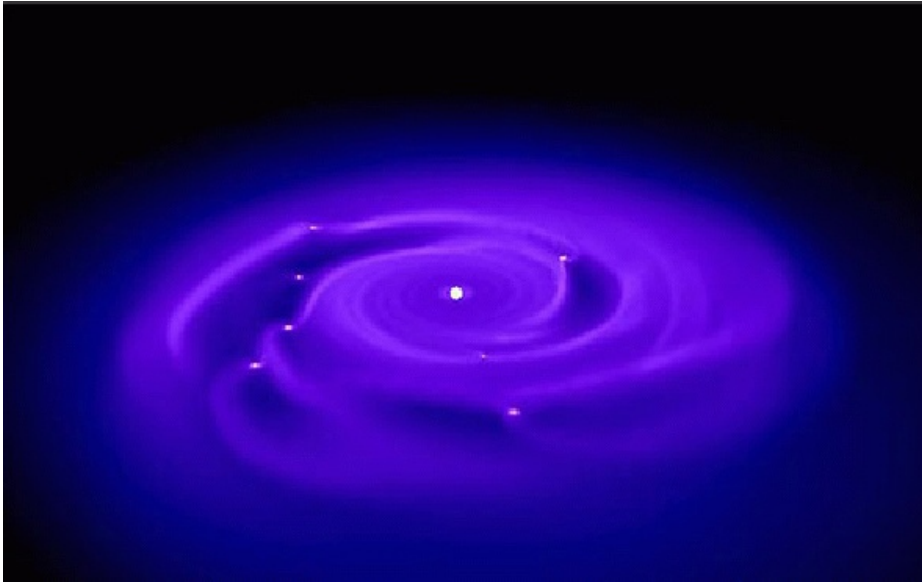
- Saturation level
- Vertical structure
- Surface temperature
- Transport efficiency (α)
- Deadzones



(Markus Flaig, ea 2012)



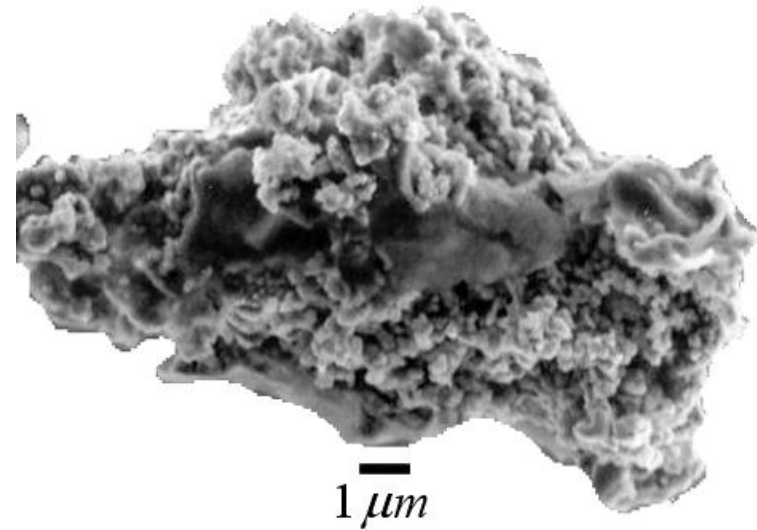
Gravitational-Instability (top-down)



(L. Mayer)

- Self-gravitating disk
- Density-Fluctuations grow
- Spiral arms \Rightarrow planets
- Fast formation (10^3 years)
- No cores**
- (Good for distant planets)

Sequential Accretion (bottom-up)

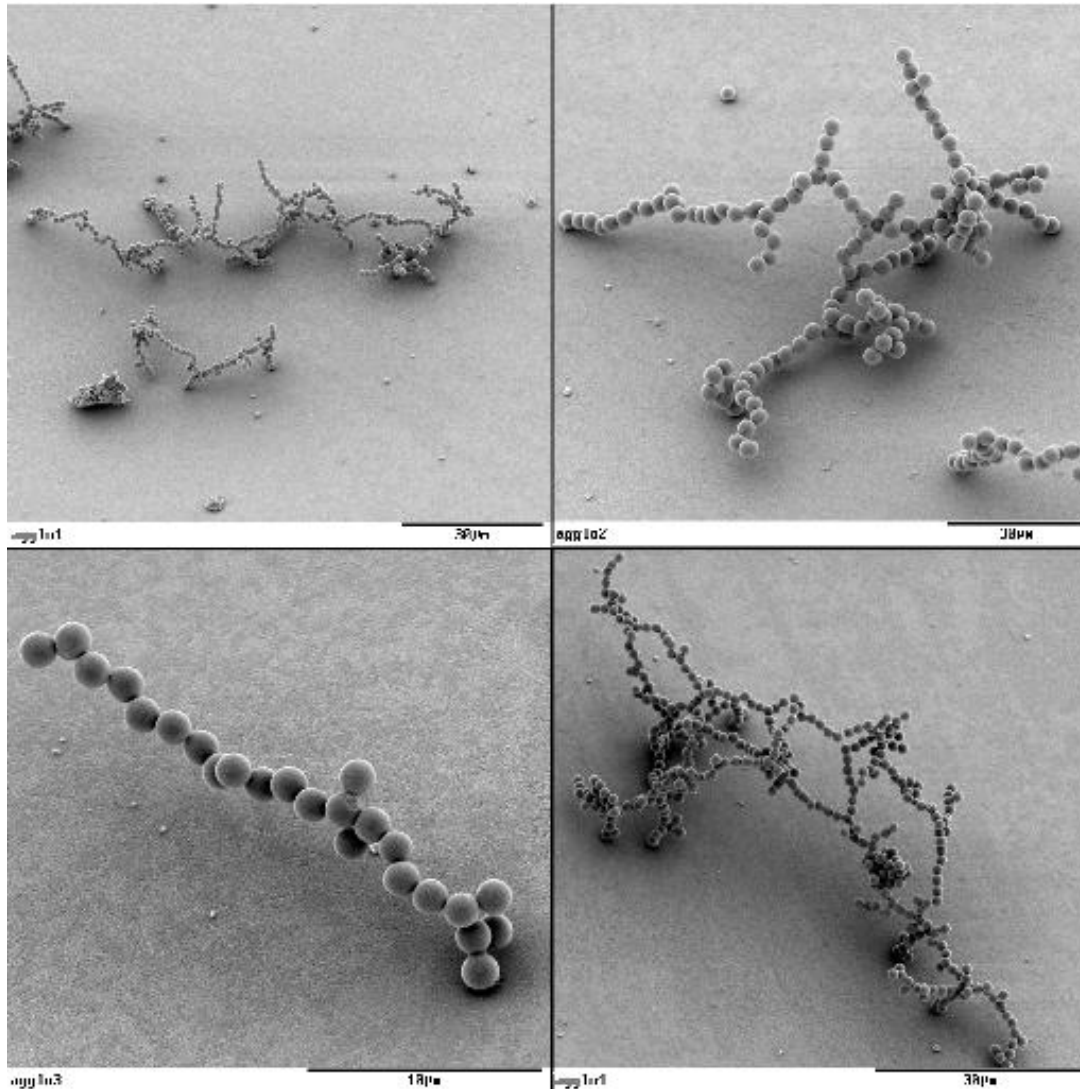


(NASA, U2)

- From small to large particles
- Slow formation (10^6 Years)
- Need: High sticking probability
- (Comets, asteroids, solid planets, cores of planets)**
- (Preferred for Solar System)



Laboratory-Experiments μm -sized particles

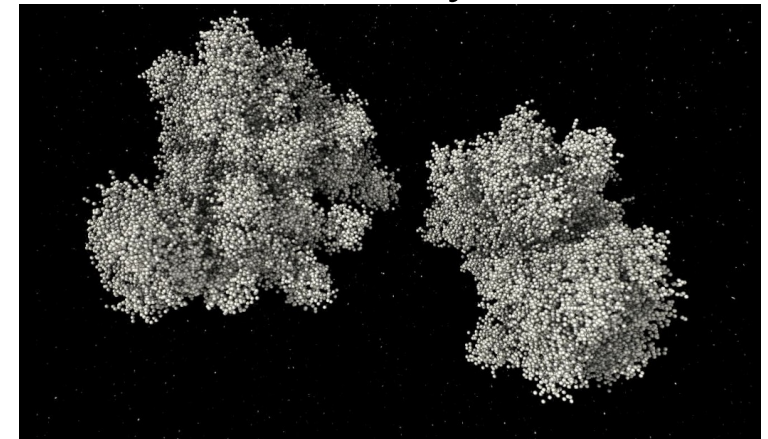


(J. Blum)

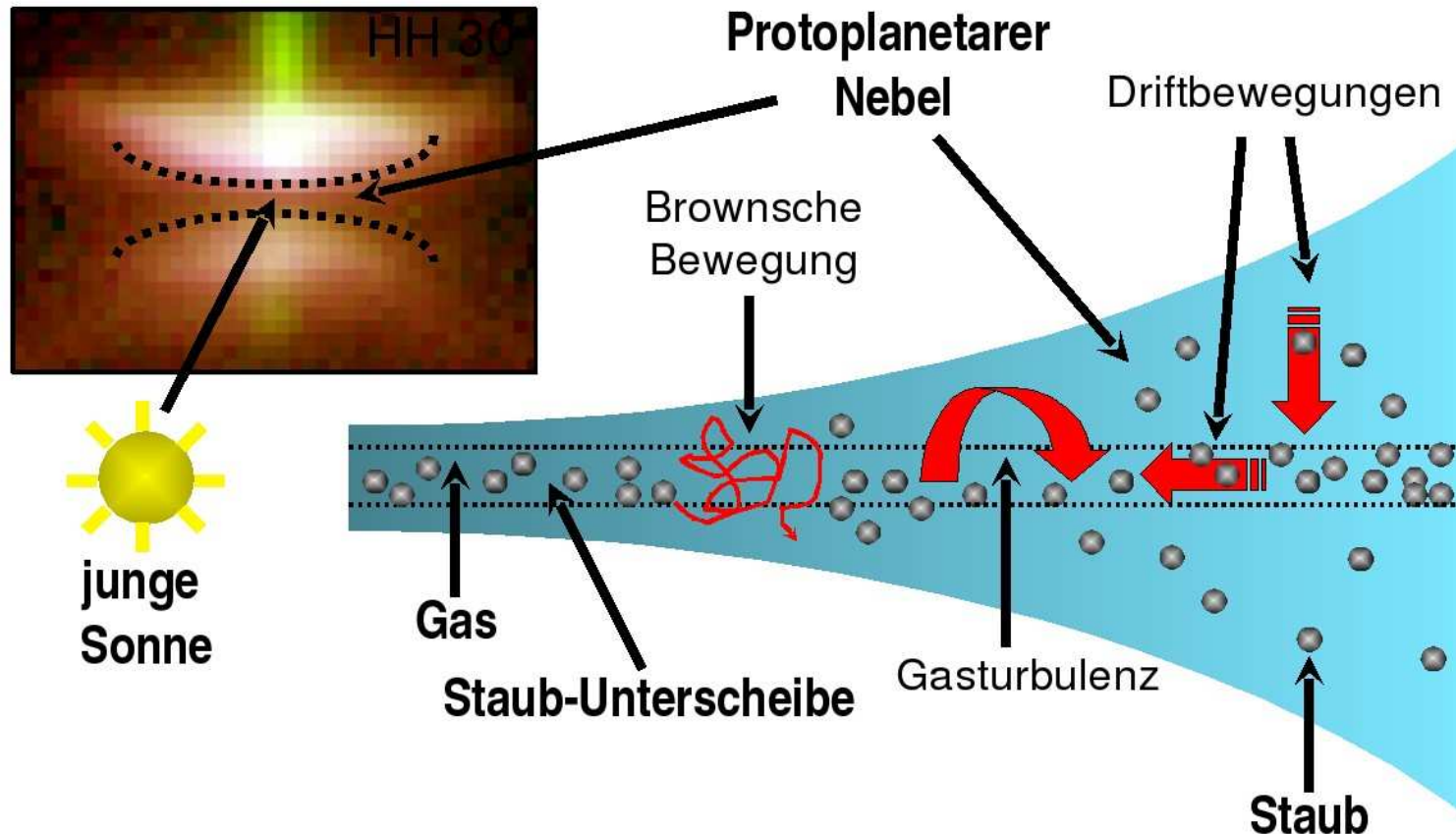
Sticking through:
Van der Waals forces
Fractal Growth
works up to cm-sizes

Numerical Simulations

Here: Molecular dynamics



(Alexander Seizinger, Tübingen)



(Jürgen Blum)

Particles have relative velocity with respect to the gas \Rightarrow frictional forces

Problem I: Fast radial drift towards star (for 1m Size: 1 AU / 100 Years)

Problem II: Destructive Collisions

Note: Disk is hotter near central star, best condensation beyond iceline

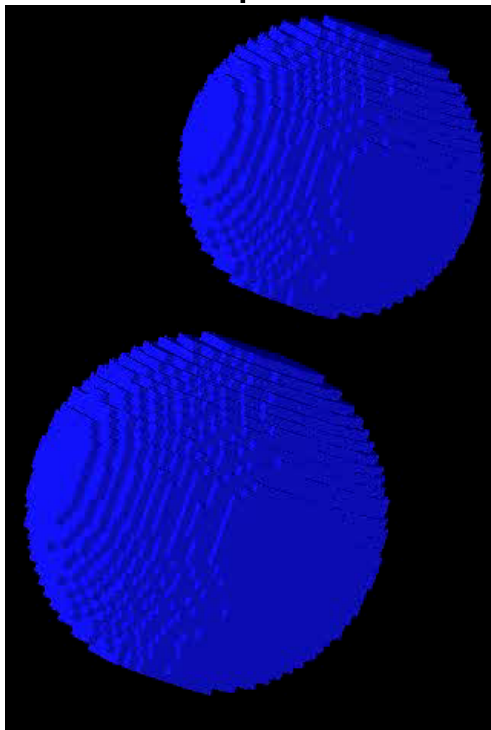


Check growth of planetesimals by collisions/accretion

SPH (Smoothed-Particle-Hydrodynamics) using 250,000-500,000 particles

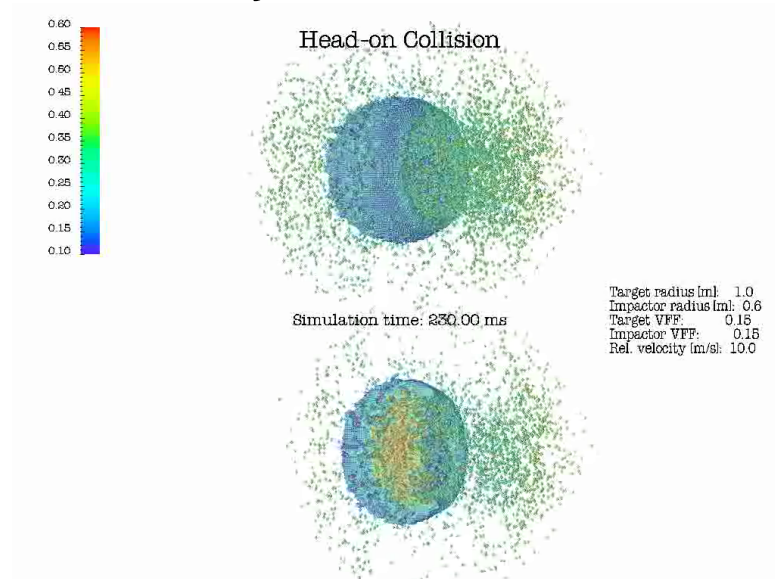
Here: Elastic-plastic strength model, formation and evolution of cracks

2 Basalt Spheres:



(Schäfer, Geretshauser, Speith, Meru; Tübingen)

Porous Objects:

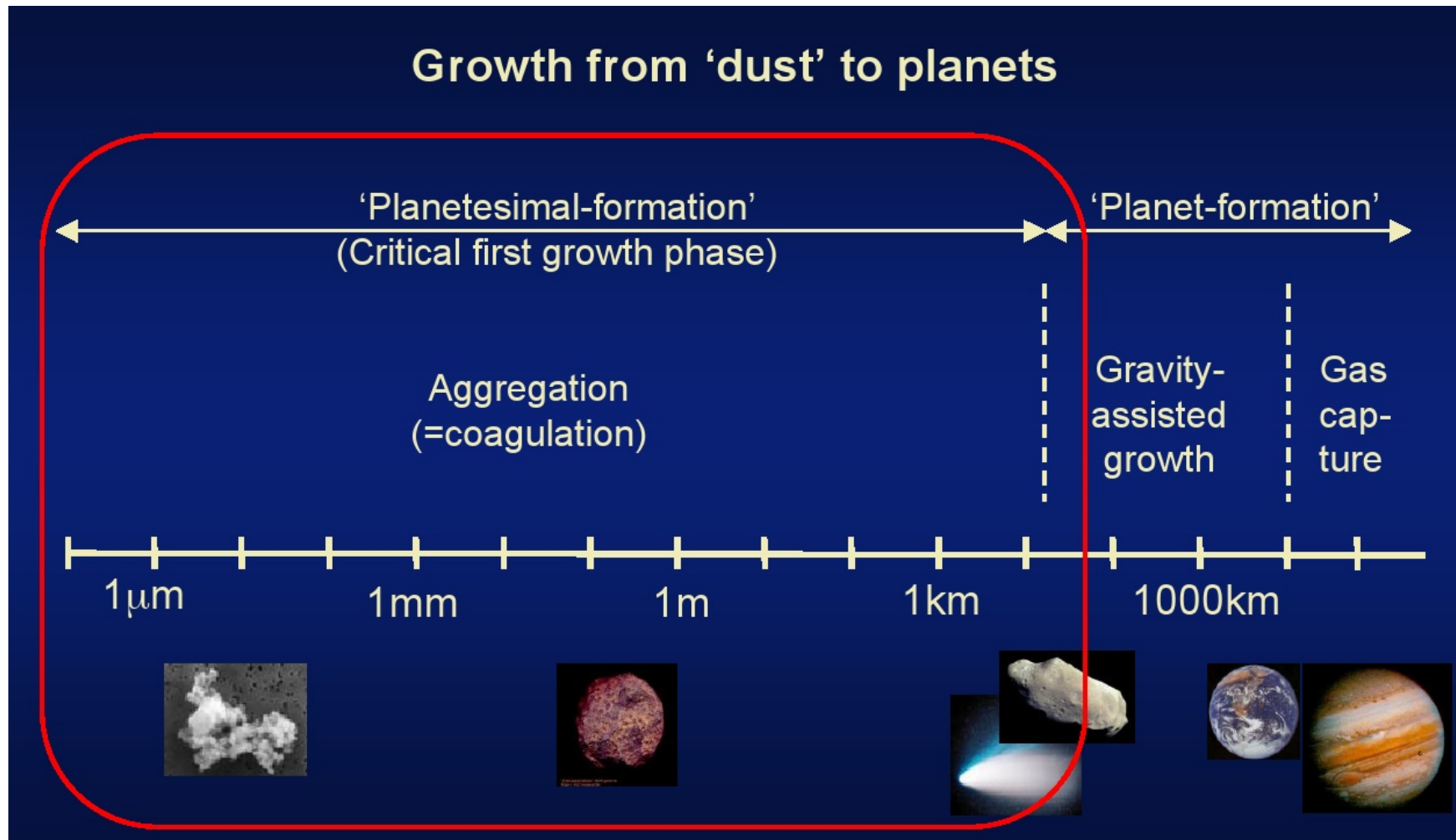


(cp. small objects in Solar System)

(Geretshauser ea., 2010, 2011a,b)

See talk by Roland Speith

to overcome growth barriers: Talks: J. Blum, F. Windmark



(C. Dullemond)

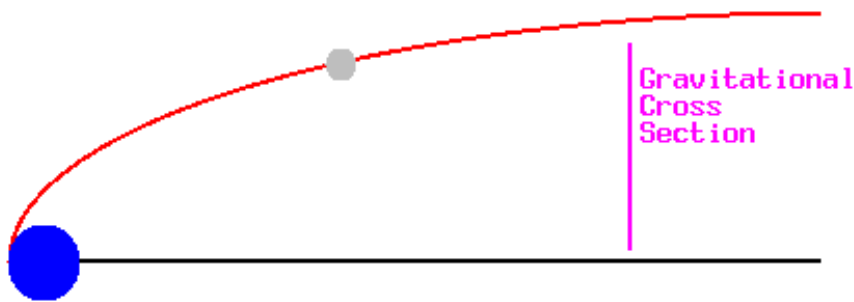
Dust \Rightarrow Planetesimals ($\mu\text{m} \Rightarrow 1\text{-}10\text{km}$, through Collisions)
 Mass rich planets: Gravitation & Gas Accretion



Important: Gravitational Focussing

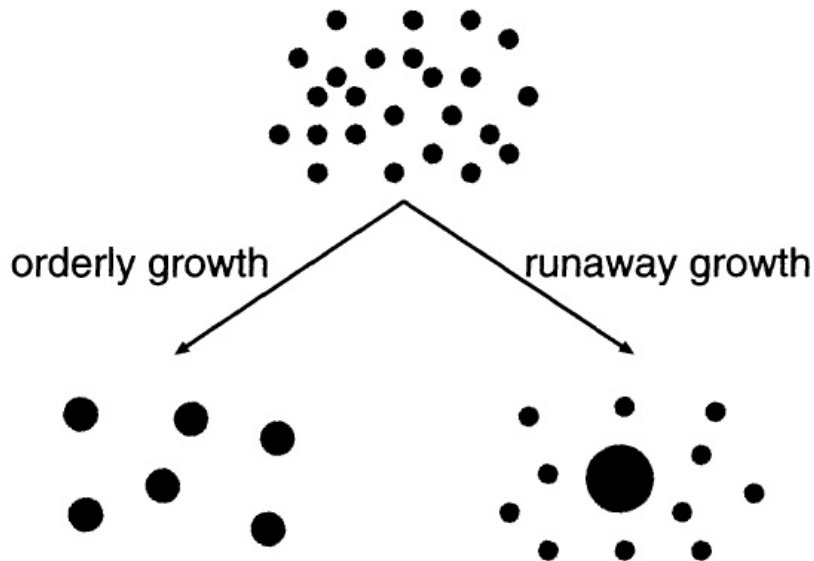
Two bodies grow through physical collisions

Mutual gravitational attraction increases the effective cross section



$$\sigma_{\text{grav}} = \sigma_{\text{geo}} \left[1 + \left(\frac{v_{\text{esc}}}{v_{\text{rel}}} \right)^2 \right]$$

⇒ Strongly enhanced collision probability



Two modes: (Kokubo, 2001)

a) **Ordered**

mass ratio of two
particles approaches unity

b) **Runaway**

Large particles grow faster $\propto M^{4/3}$

c) **Pebble accretion ?**: $\propto M^2$

(Lambrechts & Johansen, 2012;
Morbidelli & Nesvorny, 2012)



Core-instability model

First solid core (as terrestrial)
Then: Accretion of Gas Envelope
(runaway)

Core formation time $< 10^6$ years
Hydrostat. phase: $\approx 10^6$ years

Many details not known:

Opacity: κ_R

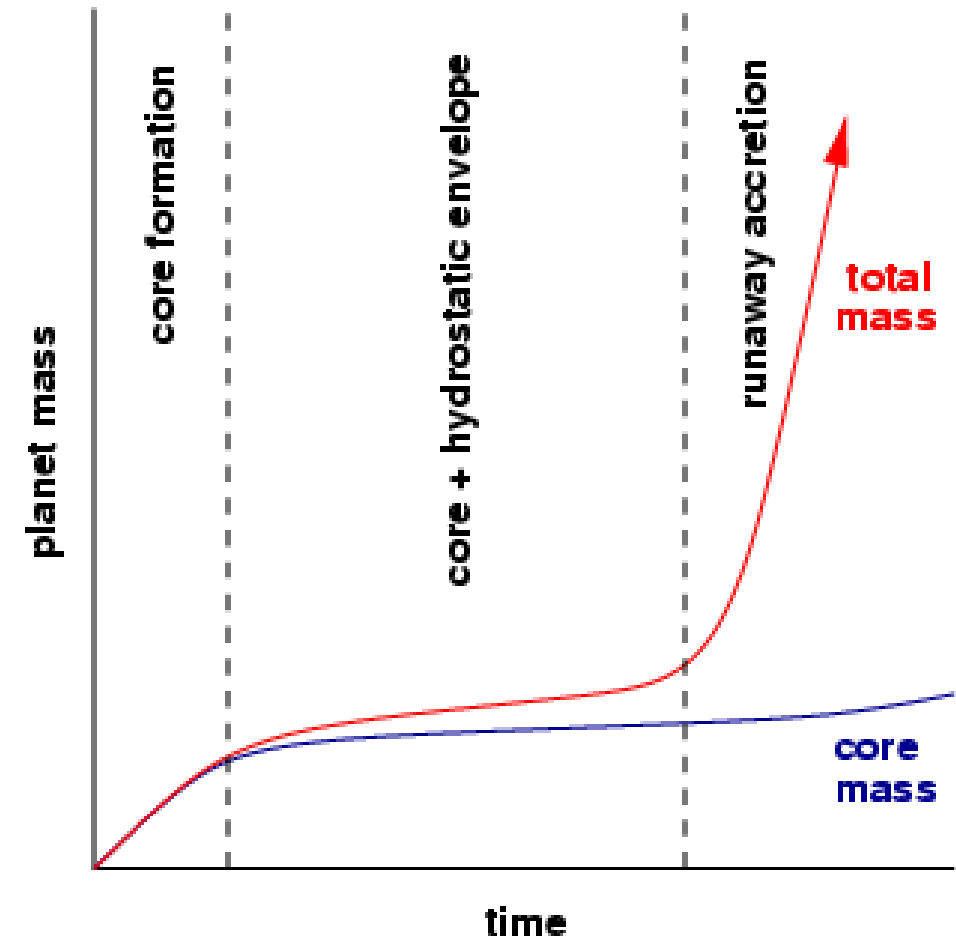
Chemical composition: μ

Accretion rate: \dot{M}_{env}

3D effects: CP-disk

Migration through disk

Schematic growth



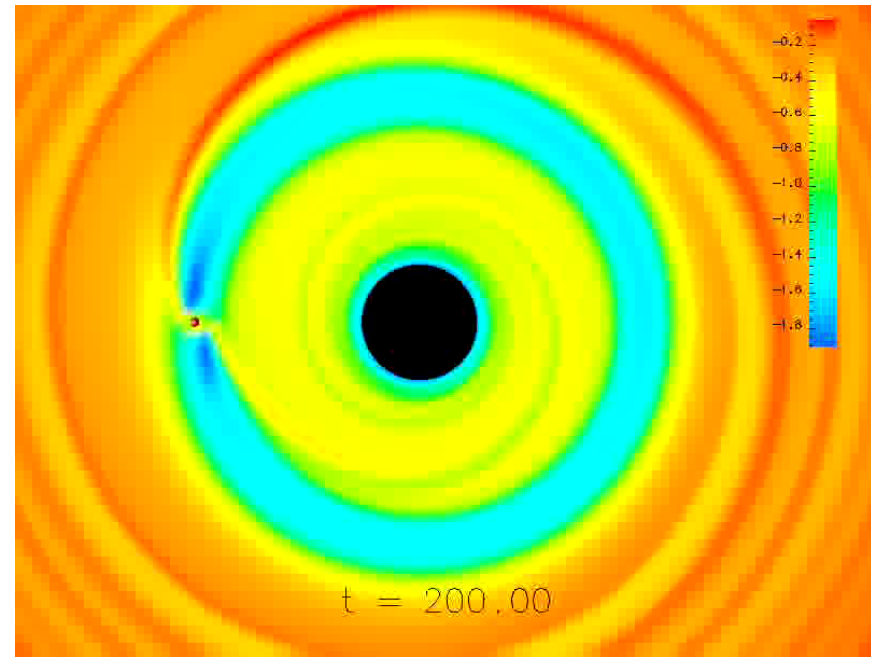
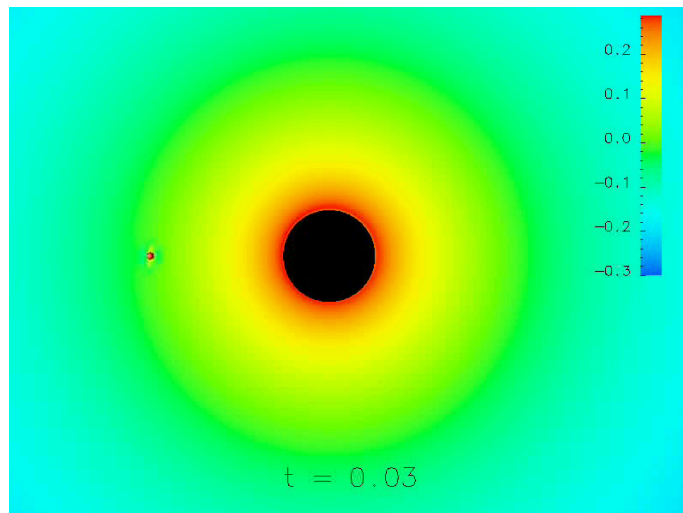
(Scholarpedia)



$M_p = 1 M_{\text{Jup}}$, $a_p = 5.2 \text{ AU}$, into disk around $1 M_{\text{sol}}$ star

Viscous hydrodynamical evolution

2D-Finite Volume Method



Spiral waves turn into shockwaves:

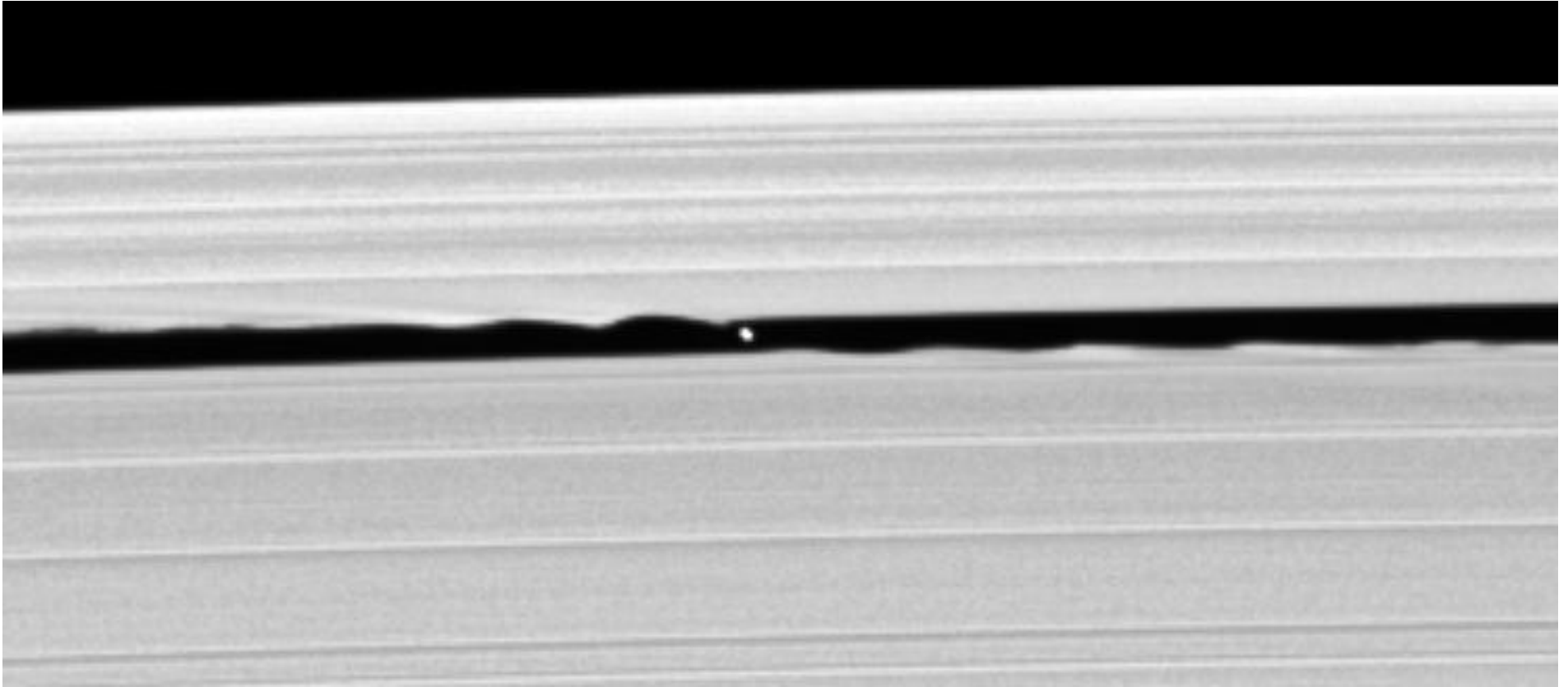
\Rightarrow dissipation & ang.mom. deposition \Rightarrow Gap

Gap formation **limits growth** to about $1 M_{\text{Jup}}$

Details (gap width & depth) depend on: Viscosity, pressure, planet mass



Moon (S/2005 S1) in Keeler Gap, Cassini (May 1, 2005)
(Gap-Width \approx 40 km, Planet-Diameter \approx 7 km)

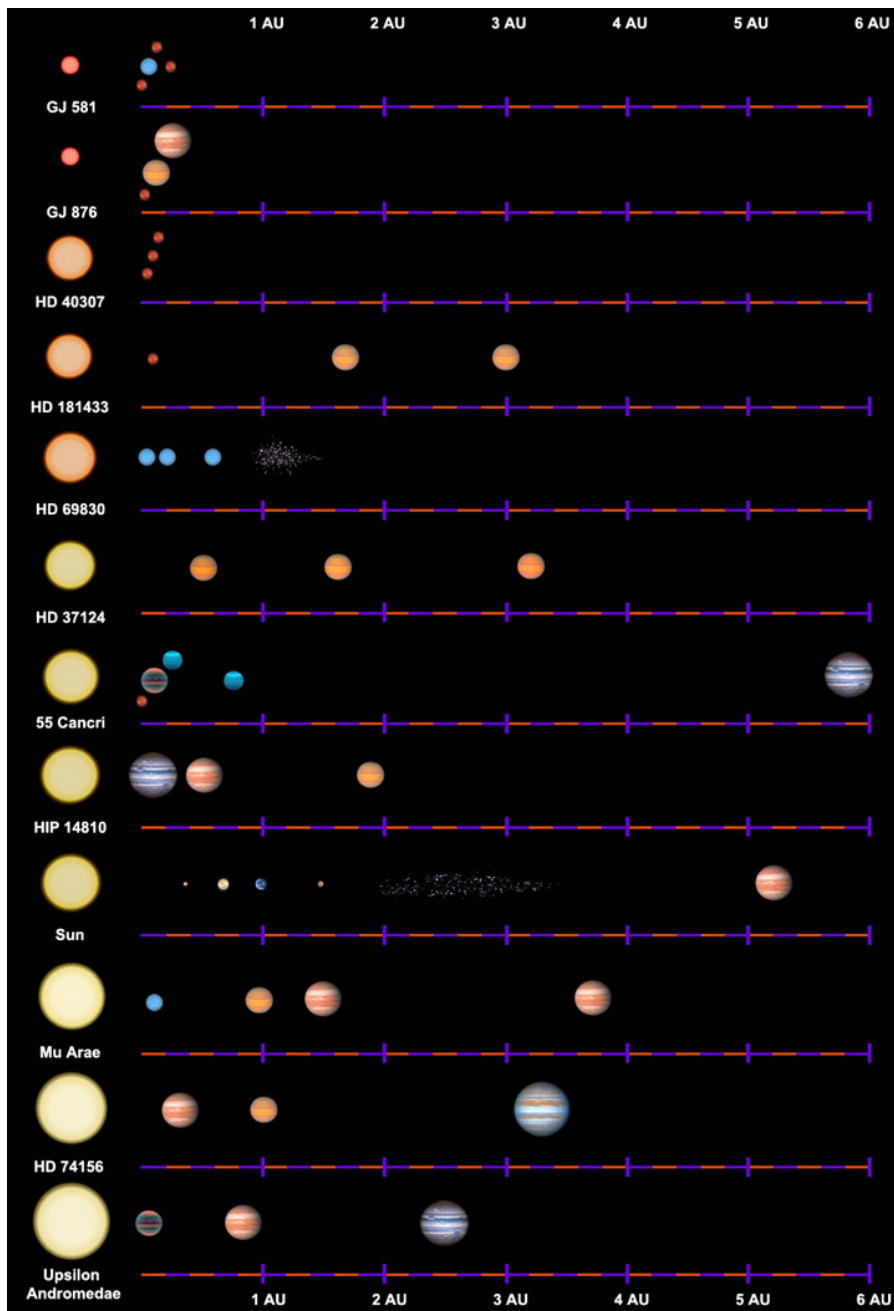


Here: Zero pressure, low viscosity \implies Very clean gap

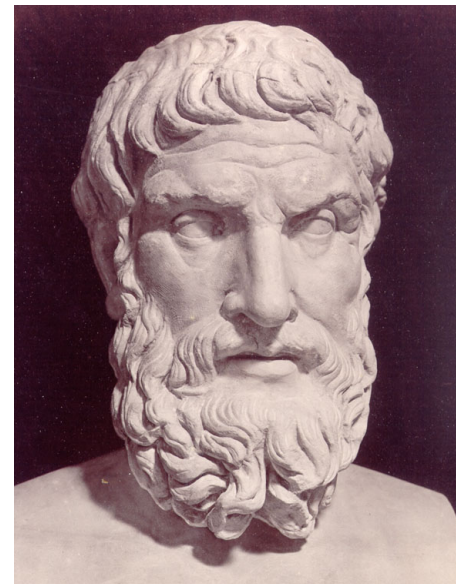


- Planets form in protoplanetary disk
(in one plane, circular orbits)
- Sequence of sticking collisions
- Inner planets: solid
- Outer planets: gaseous with cores
- Maximum Mass $\approx 1 M_{\text{Jup}}$ (gap formation)

What about the extrasolar planets ?



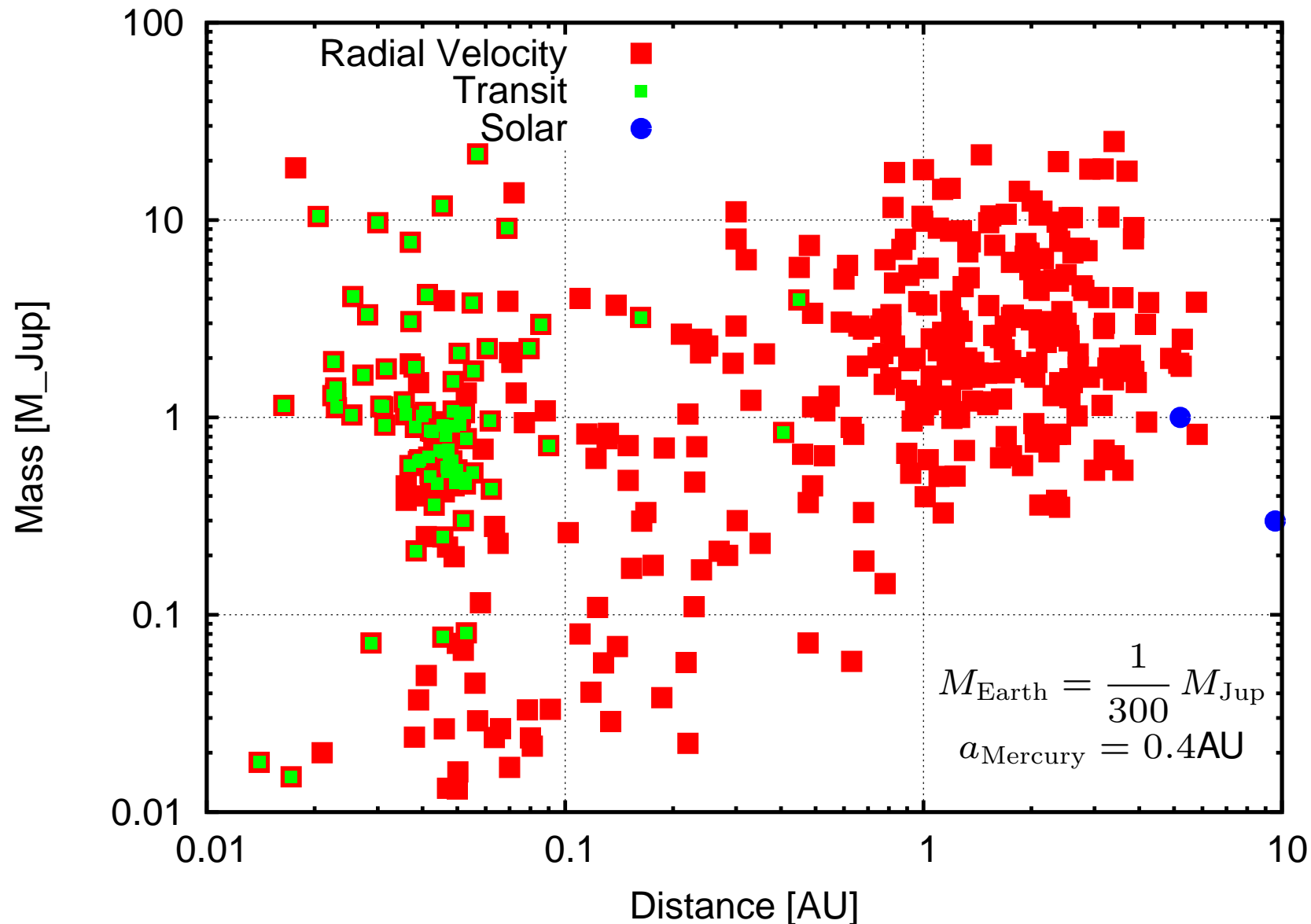
Epicurus (ca. 341-270 BC)
 “There is an **infinite number of worlds**, some similar to ours some very different.”





Small distances (hot Jupiters) & large masses

(Data: exoplanet.eu)



⇒ Need migration and mass growth!



- Not possible to form hot Jupiters in situ
 - disk too hot for material to condense
 - not enough material
- Difficult to form massive planets
 - gap formation

But planets grow in disks:

⇒ Have a closer look at planet-disk interaction

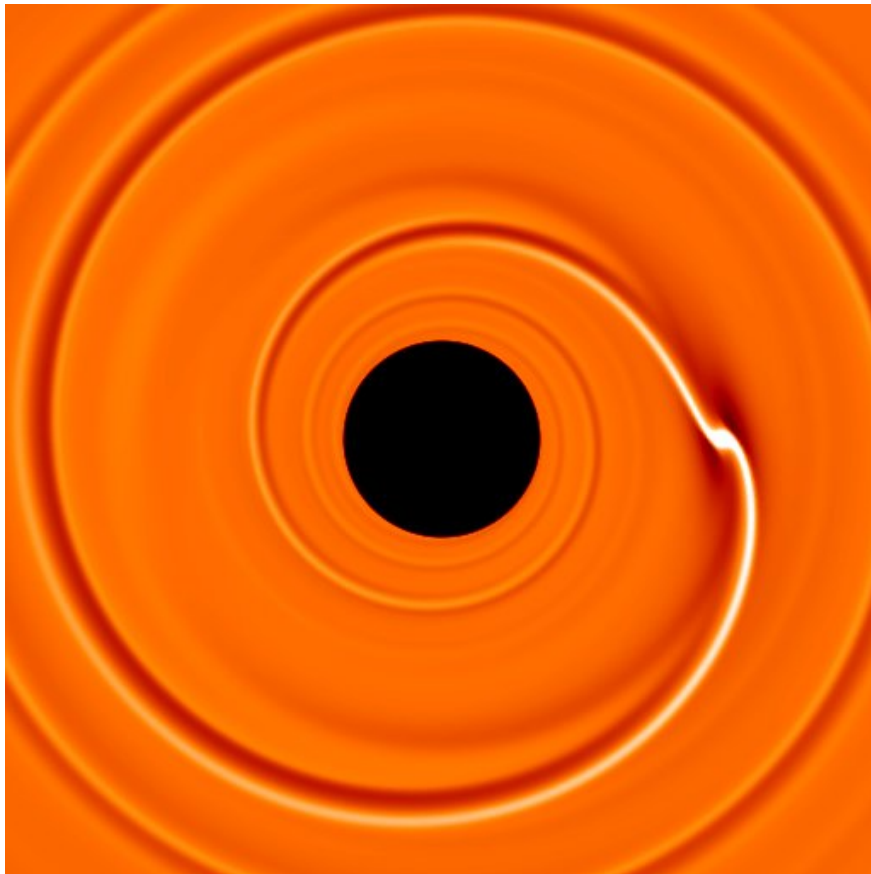
see Annual Review article: [Kley & Nelson, ARAA, 50 \(2012\)](#)



Young planets are embedded in gaseous disk

Creation of **spiral arms**:

- stationary in planet frame
- Linear analysis,
- 2D hydro-simulations



Inner Spiral

- pulls planet forward:
- positive torque

Outer Spiral

- pulls planet backward:
- negative torque

→ Net Torque

⇒ Migration

Most important:

Strength & Direction ?

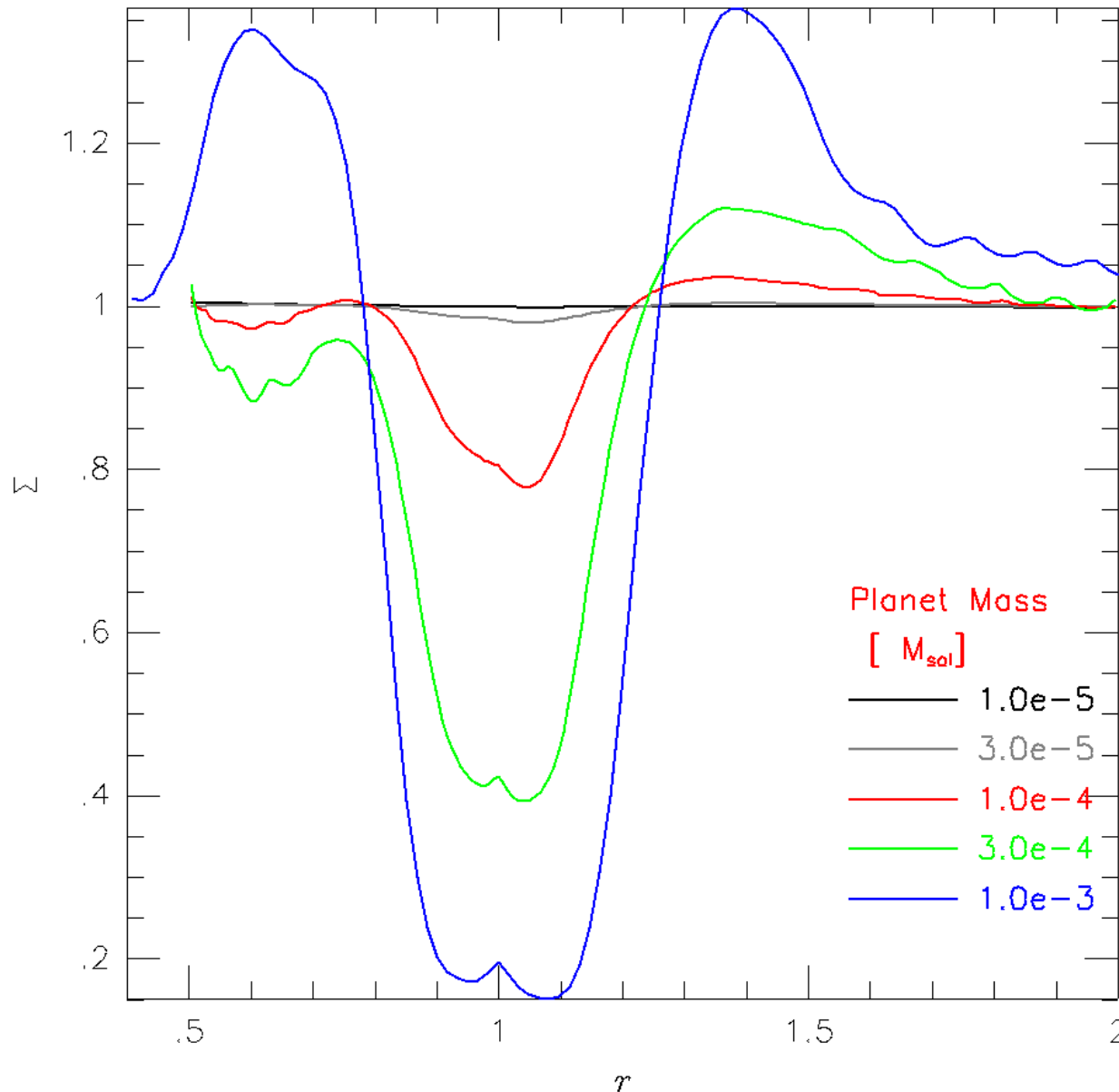
Typically: Outer spiral wins

⇒ Inward Migration

Torque scales with:

inv. Temp. $(H/r)^{-2}$, M_p^2 , M_d

(Masset, 2001)



$$M_p = 0.01 M_{\text{Jup}}$$

$$M_p = 0.03 M_{\text{Jup}}$$

$$M_p = 0.1 M_{\text{Jup}}$$

$$M_p = 0.3 M_{\text{Jup}}$$

$$M_p = 1.0 M_{\text{Jup}}$$

Depth depends on:

- M_p
- Viscosity
- Temperature

Torques reduced:

Migration slows

Type I \Rightarrow Type II

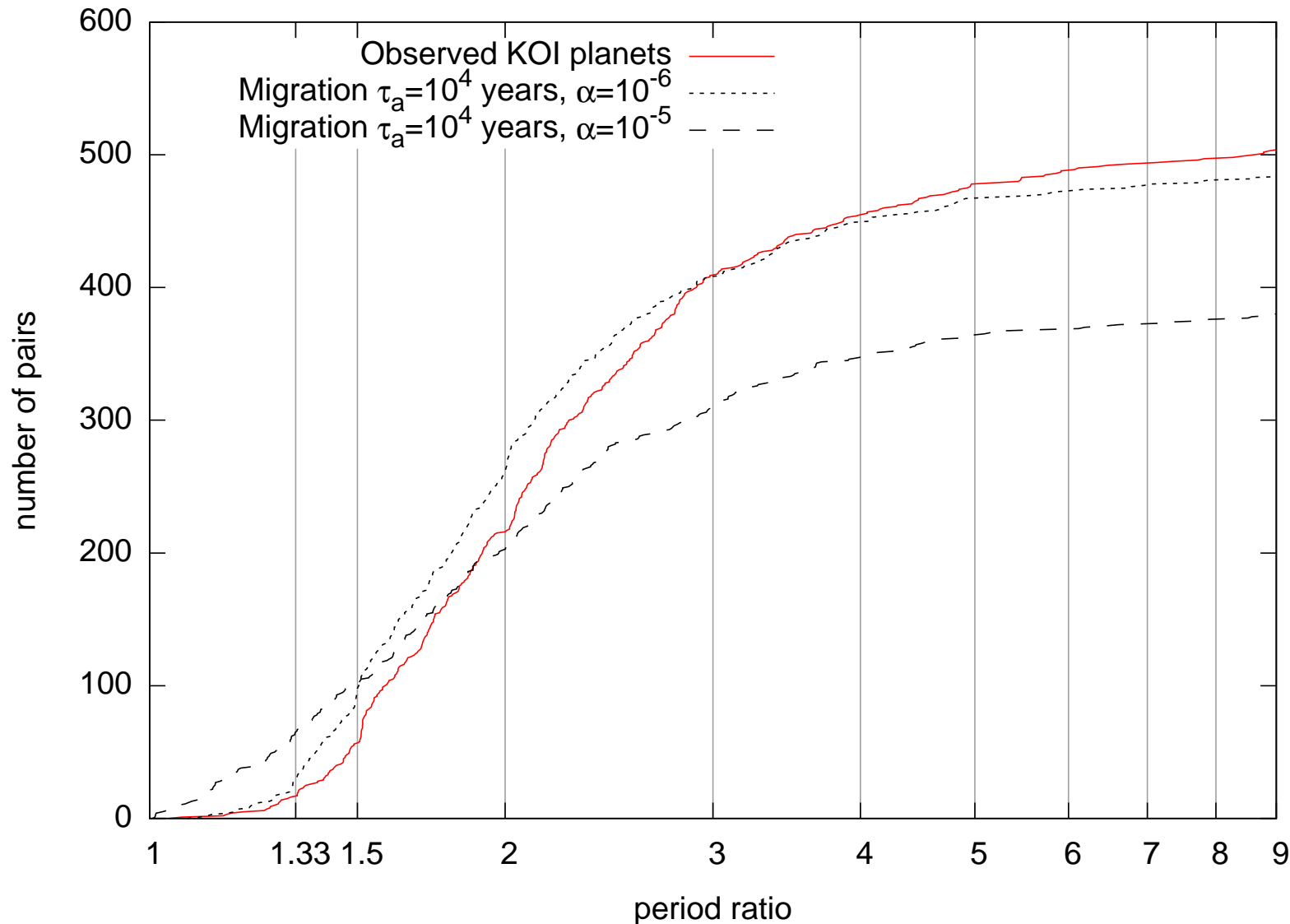
linear \Rightarrow non-linear



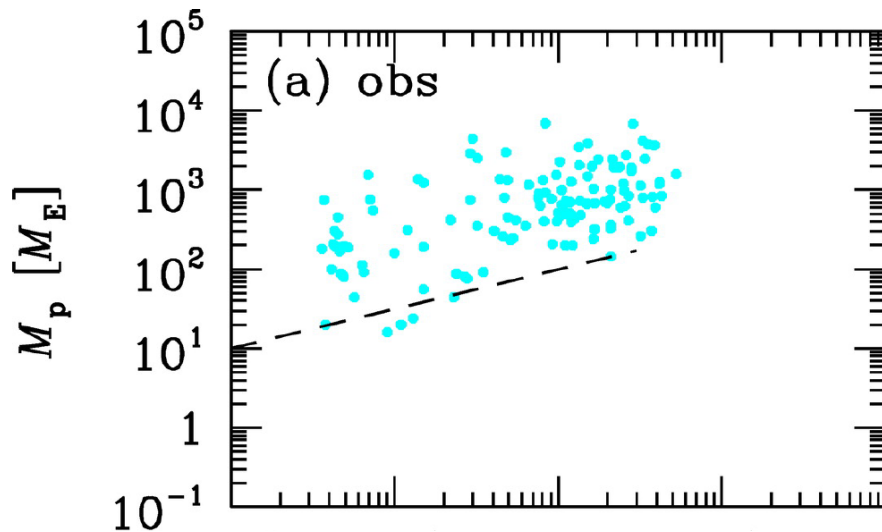
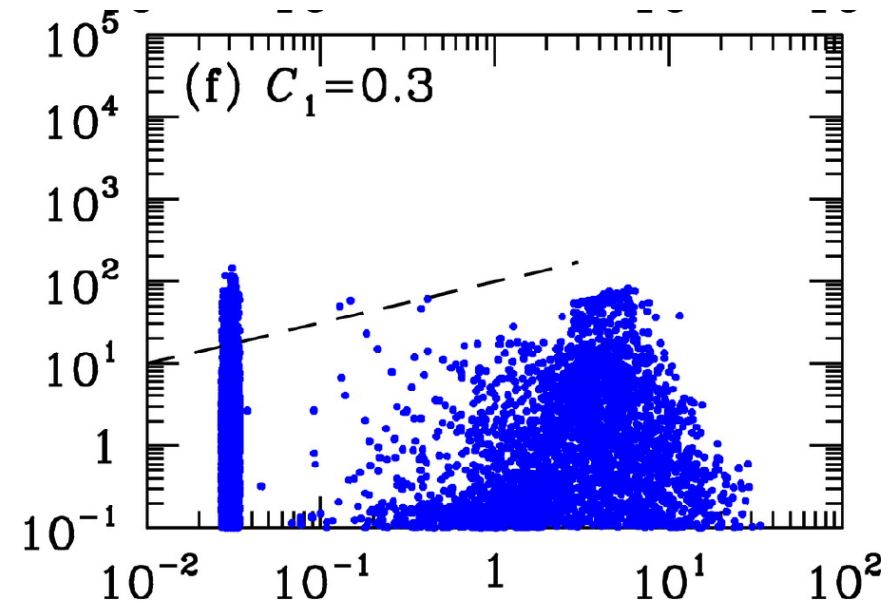
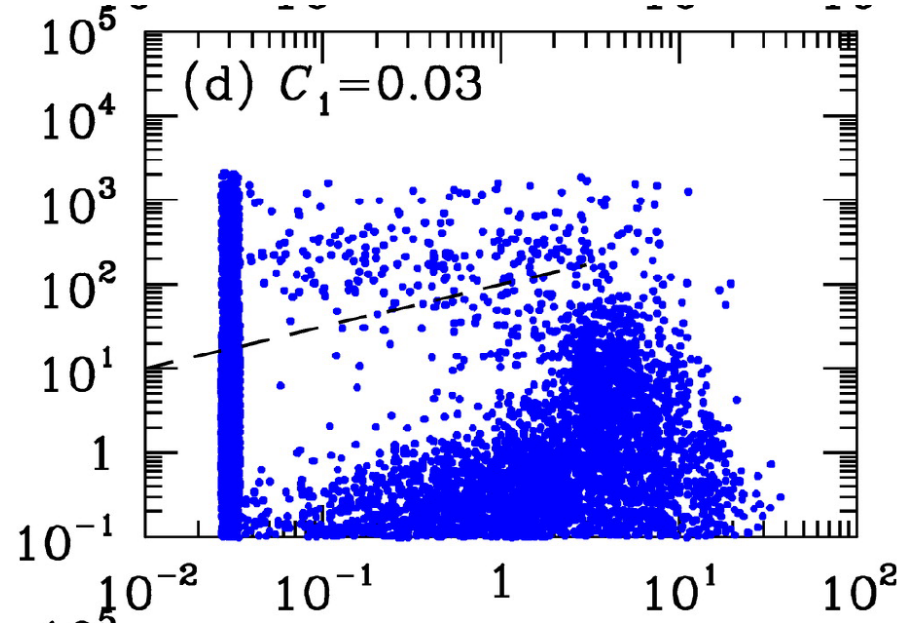
- RV-Obs.: ≈ 50 multi-planet extrasolar planetary systems
 $\approx 1/4$ contain planets in a low-order **mean-motion resonance** (MMR)
mostly in a 2:1 configuration (eg. GJ 876, HD 128311, HD 82943,)
recently 3:2 (HD 45364) and 3:1 (HD 60532)
In **Solar System**: 3:2 between Neptune and Pluto (plutinos)
- Resonant capture through convergent migration process
dissipative forces due to disk-planet interaction
- Existence of resonant systems
 \implies **Clear evidence for planetary migration**
- Hot Jupiters (Neptunes) & Kepler systems
 \implies **Clear evidence for planetary migration**



Migration ($\tau = 10^4$ yrs) with stochastic forcing (H. Rein, 2012)



— observed, - - - stronger stochastic forcing, ... weak forcing



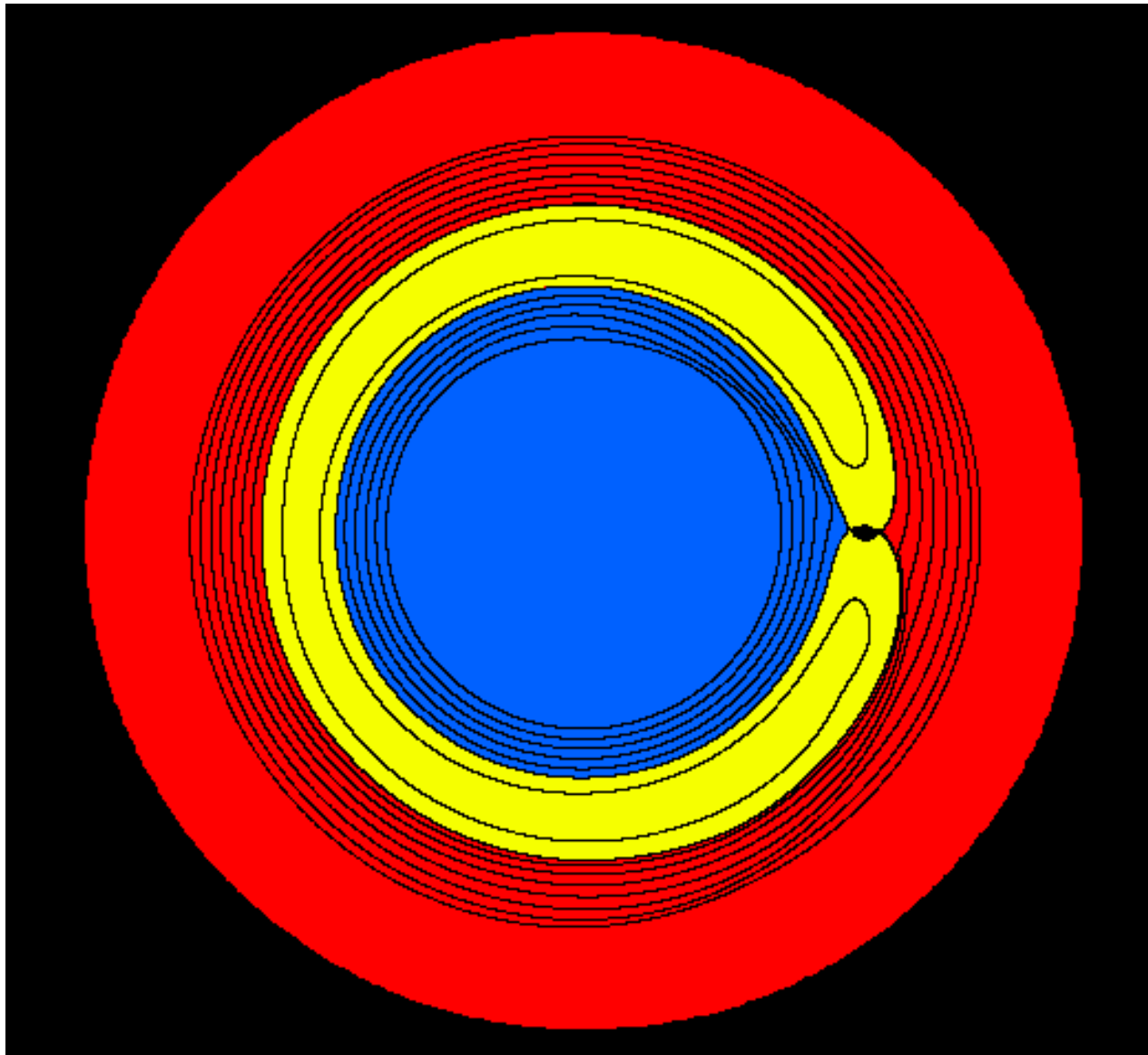
Migration too efficient!

Only strong reduction of Type I gives reasonable results

(Ida & Lin; Mordasini, Alibert & Benz)

⇒ Need improvements:

- stochastic migration
- inviscid, self-grav. disks
- radiative disks (corotation effects)



3 Regions

Outer disk (spiral)

Inner disk (spiral)

⇒ Lindblad torques

Horseshoe (coorbital)

⇒ Corotation Torques
(Horseshoe drag)

Scaling with:

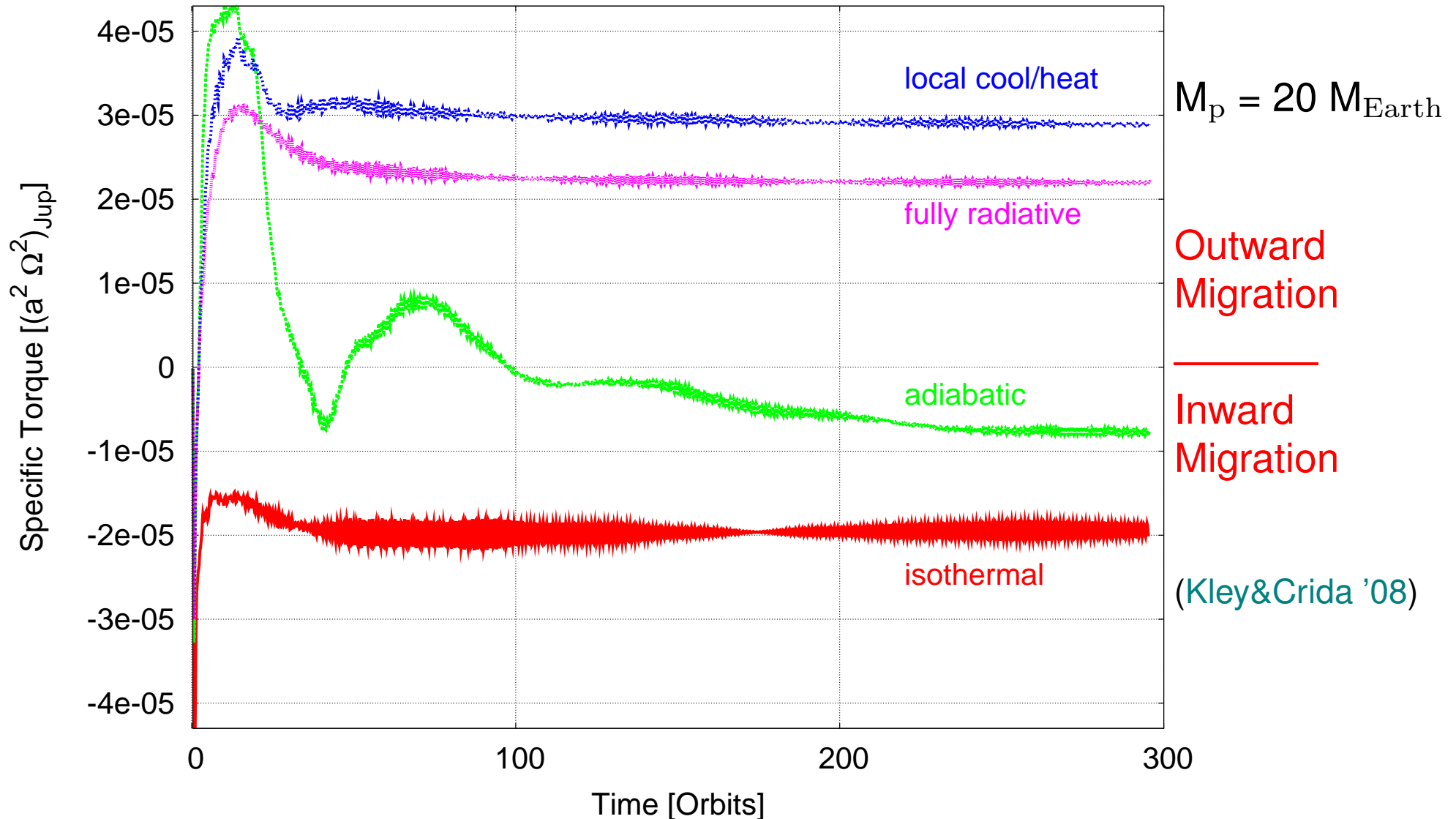
- Vortensity gradient
- Entropy gradient

(Talk by: C. Baruteau)

(F. Masset)



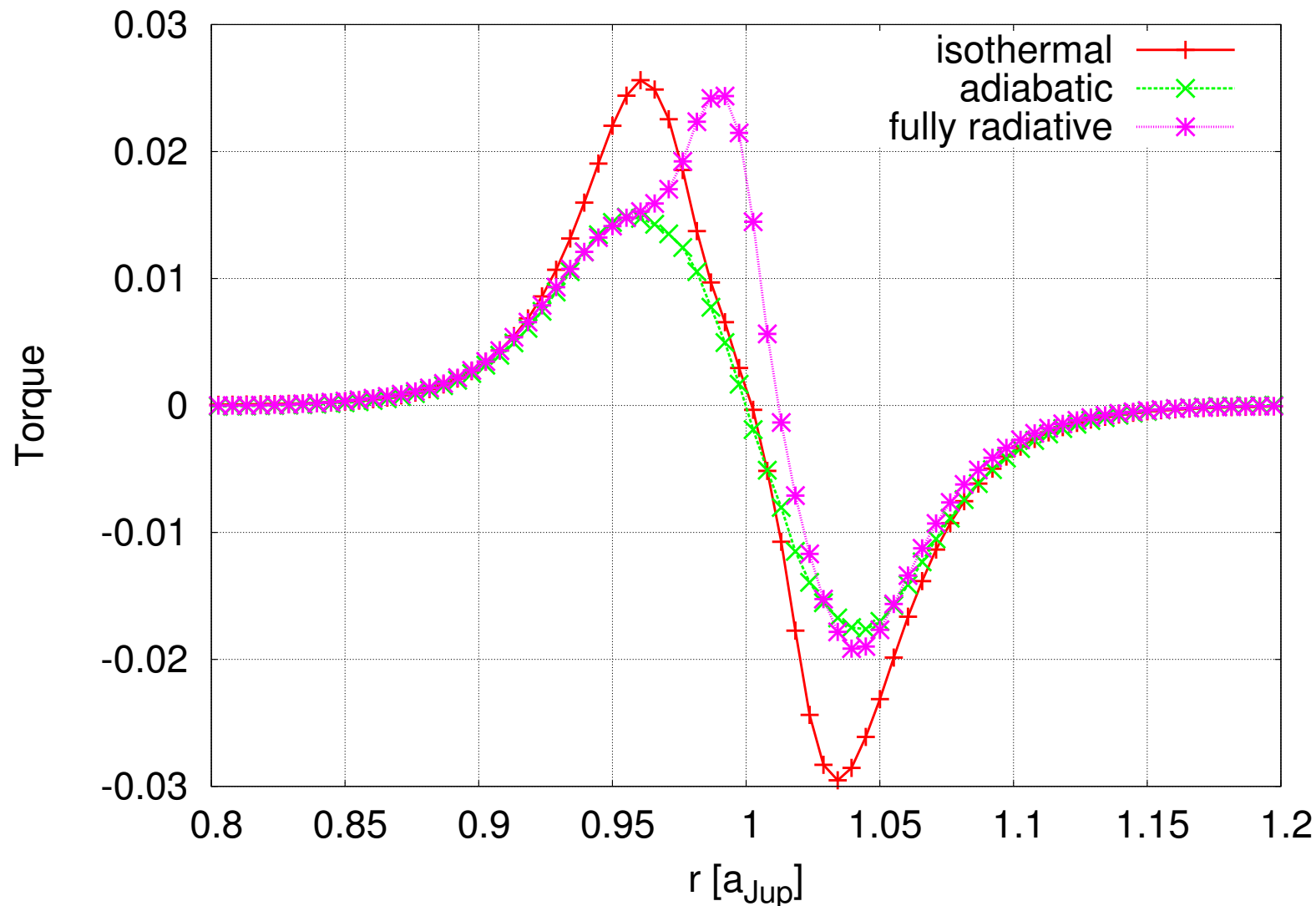
$$\frac{\partial \Sigma c_v T}{\partial t} + \nabla \cdot (\Sigma c_v T \mathbf{u}) = -p \nabla \cdot \mathbf{u} + D - Q - 2H \nabla \cdot \vec{F}$$

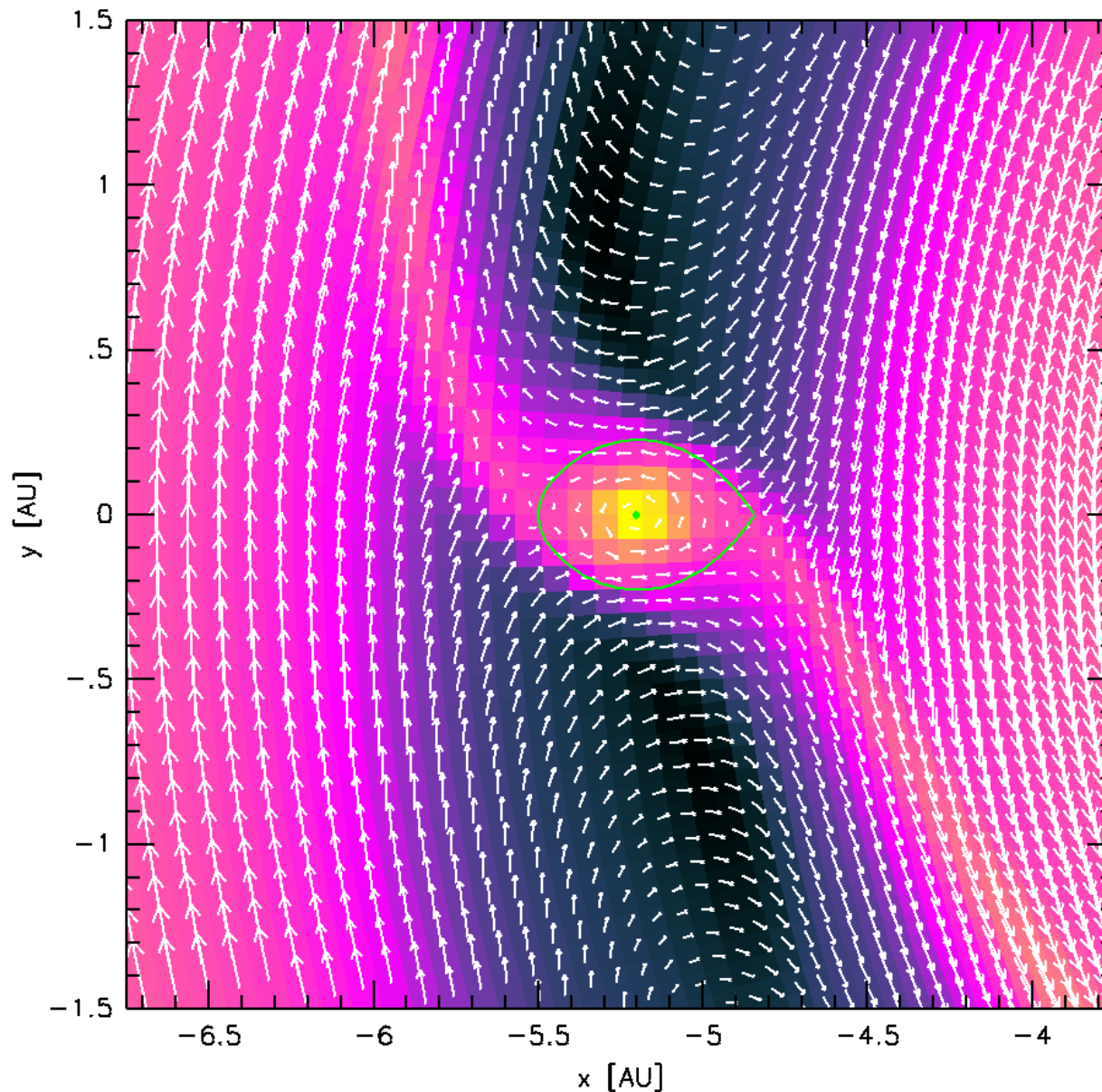




3D-simulations, radiative diffusion, 20 M_{Earth} planet (Kley, Bitsch & Klahr 2009)

$d\Gamma/dm$, with $\Gamma_{\text{tot}} = 2\pi \int (d\Gamma/dm) \Sigma dr$ Radiative: \Rightarrow additional positive contrib.





Surface Density
at 200 Orbits

Green Dot: Planet
Green Line:
Roche-Lobe

$m_p = 1 M_{\text{Jup}}$
 $a_p = 5.2 \text{ AE}$

Flow-Field

→ Mass growth
up to a few M_{Jup}
→ prograde rotation

(WK, 2000)



Planet-disk interaction: Torques on Planet

Isothermal Migration is inward & rapid (lose planets)

But: $\Gamma_{\text{tot}} = \Gamma_{\text{L}} + \Gamma_{\text{HS,ent}} + \Gamma_{\text{HS,vort}}$

Outward in radiative disks

Mass limit due to gap opening

Driven by:

Vortensity gradient

maintained by: viscosity

Entropy gradient

maintained by: rad. diffusion (or cooling)

- cooling time \approx libration time

Need viscosity

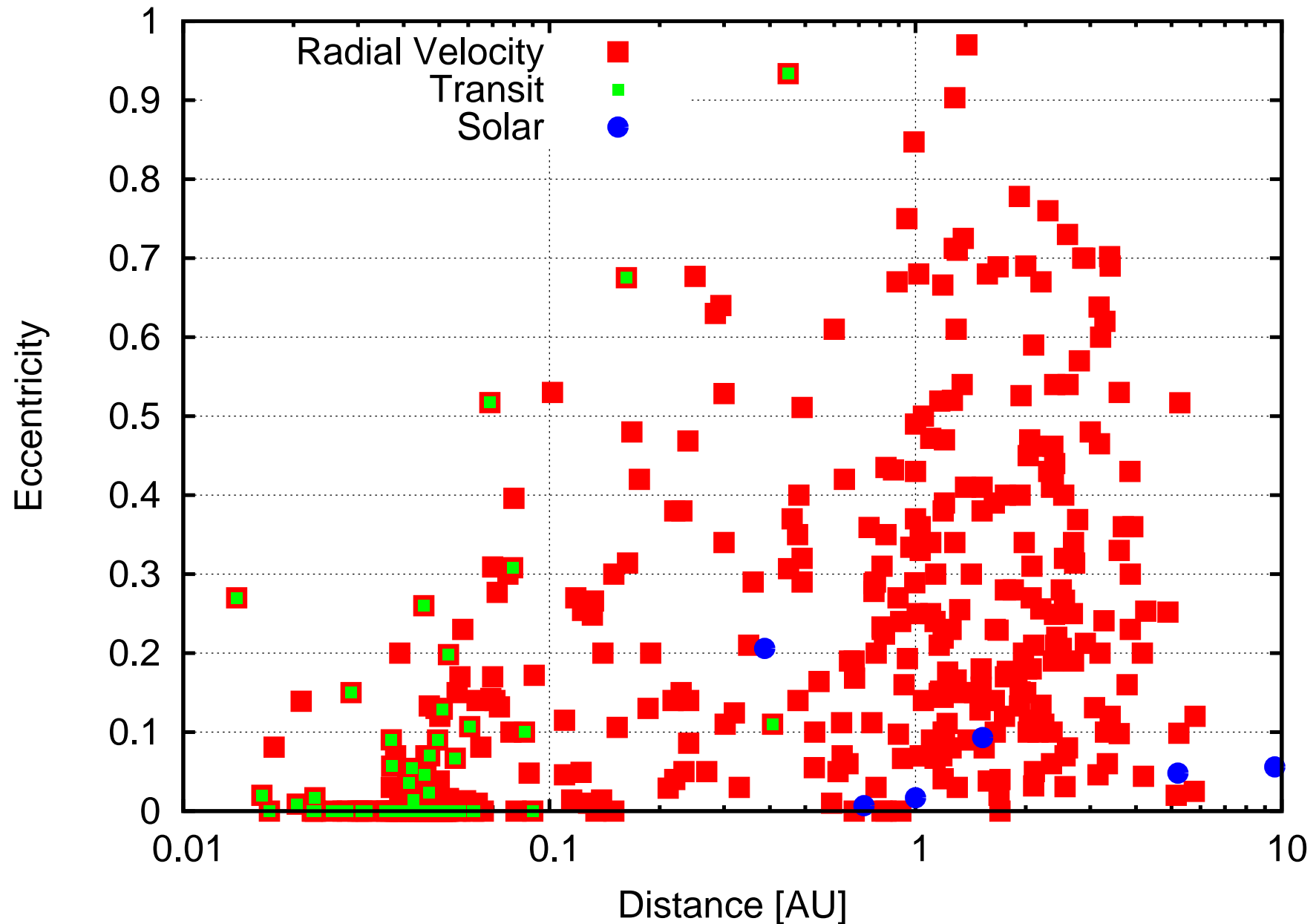
approximate torque formulae: [Paardekooper et al.](#); [Masset&Casoli](#)

More details: [Talks by C. Baruteau](#) and [B. Bitsch](#)



Large eccentricities (similar to binary stars)

(Data: exoplanet.eu)



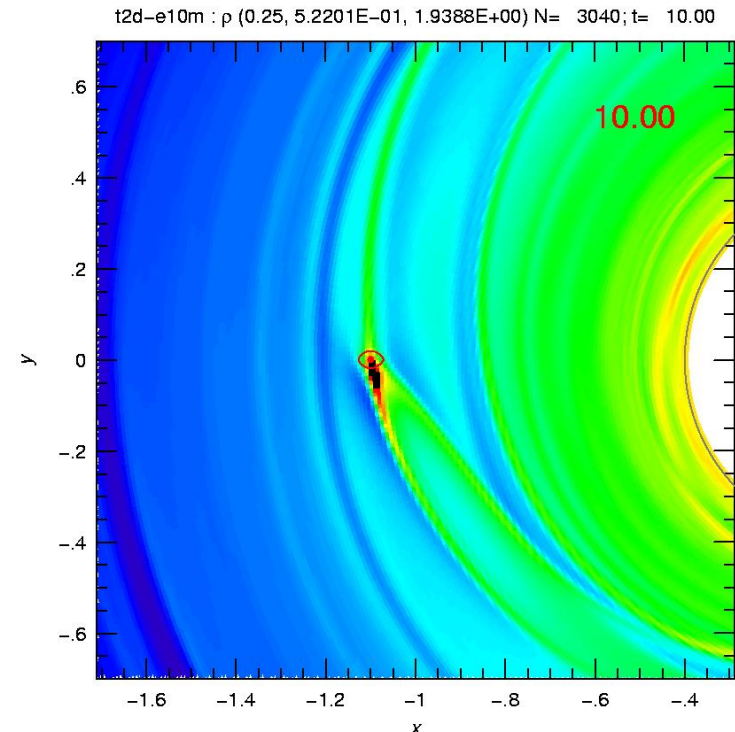
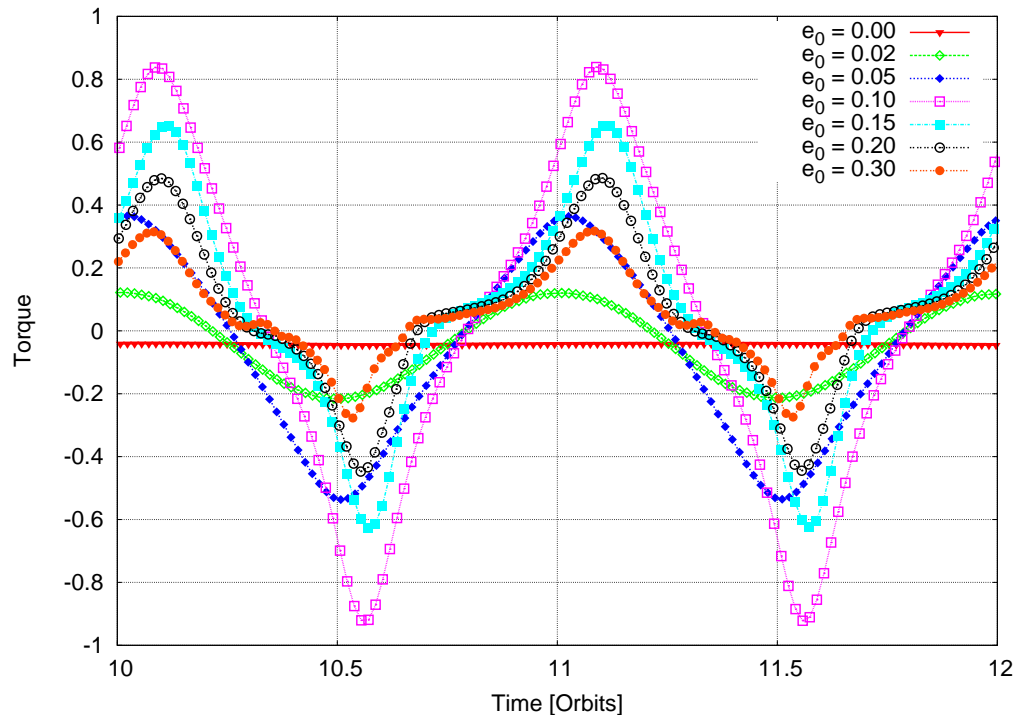


Torque on planet due to disk

Power: Energy loss of planet

$$\Gamma_{\text{disk}} = \int_{\text{disk}} (\vec{r}_{\text{P}} \times \vec{F}) \Big|_z df$$

$$P_{\text{disk}} = \int_{\text{disk}} \dot{\vec{r}}_{\text{P}} \cdot \vec{F} df$$



$$L_{\text{p}} = m_{\text{p}} \sqrt{GM_* a} \sqrt{1 - e^2}$$

$$\frac{\dot{L}_{\text{p}}}{L_{\text{p}}} = \frac{1}{2} \frac{\dot{a}}{a} - \frac{e^2}{1 - e^2} \frac{\dot{e}}{e} = \frac{\Gamma_{\text{disk}}}{L_{\text{p}}}$$

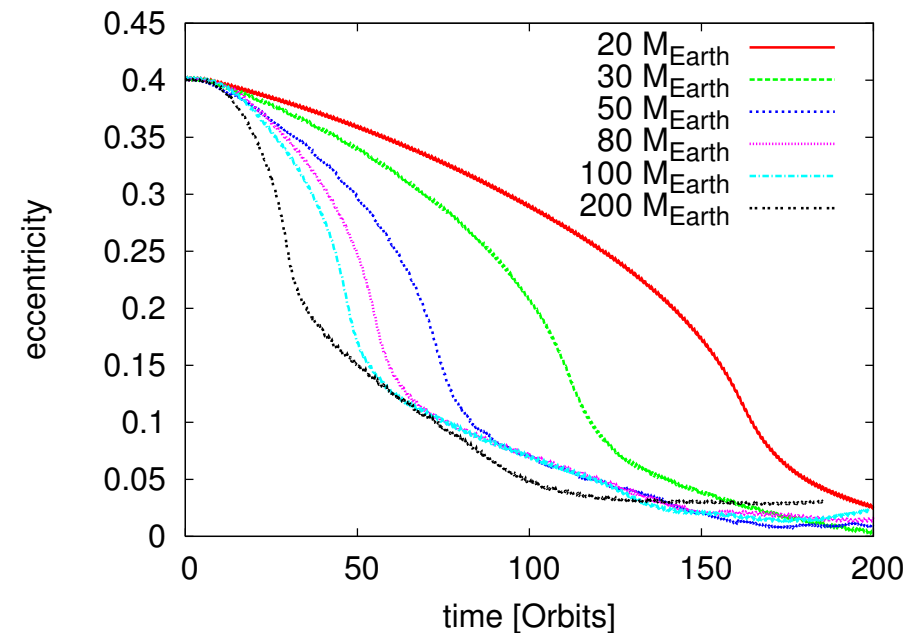
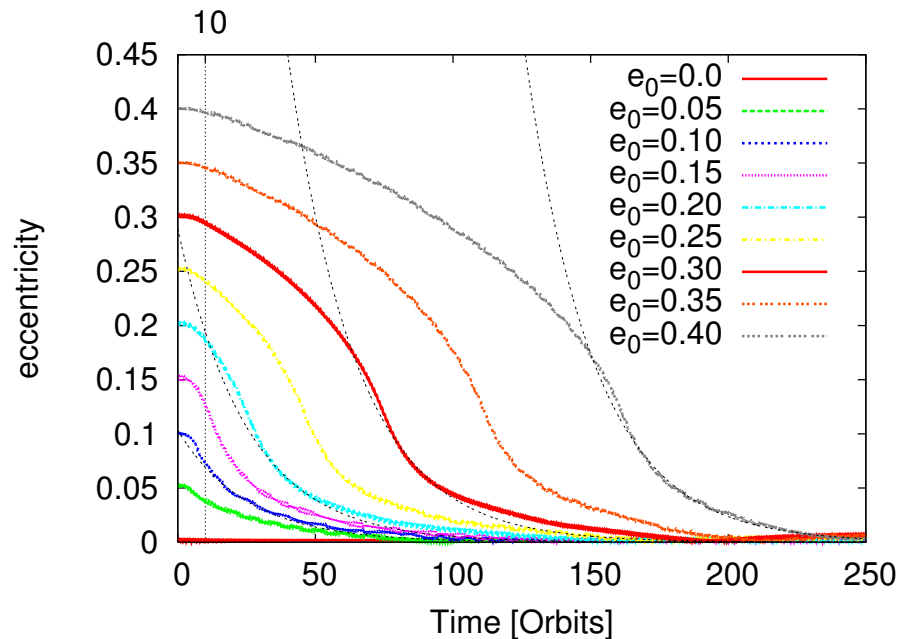
$$E_{\text{p}} = -\frac{1}{2} \frac{GM_* m_{\text{p}}}{a}$$

$$\frac{\dot{E}_{\text{p}}}{E_{\text{p}}} = \frac{\dot{a}}{a} = \frac{P_{\text{disk}}}{E_{\text{p}}}$$



Fix planet mass $M_p = 20M_{\text{Earth}}$
- Vary initial Eccentricity

Vary Planet Mass 10 – $200M_{\text{Earth}}$
- Same $e_0 = 0.40$



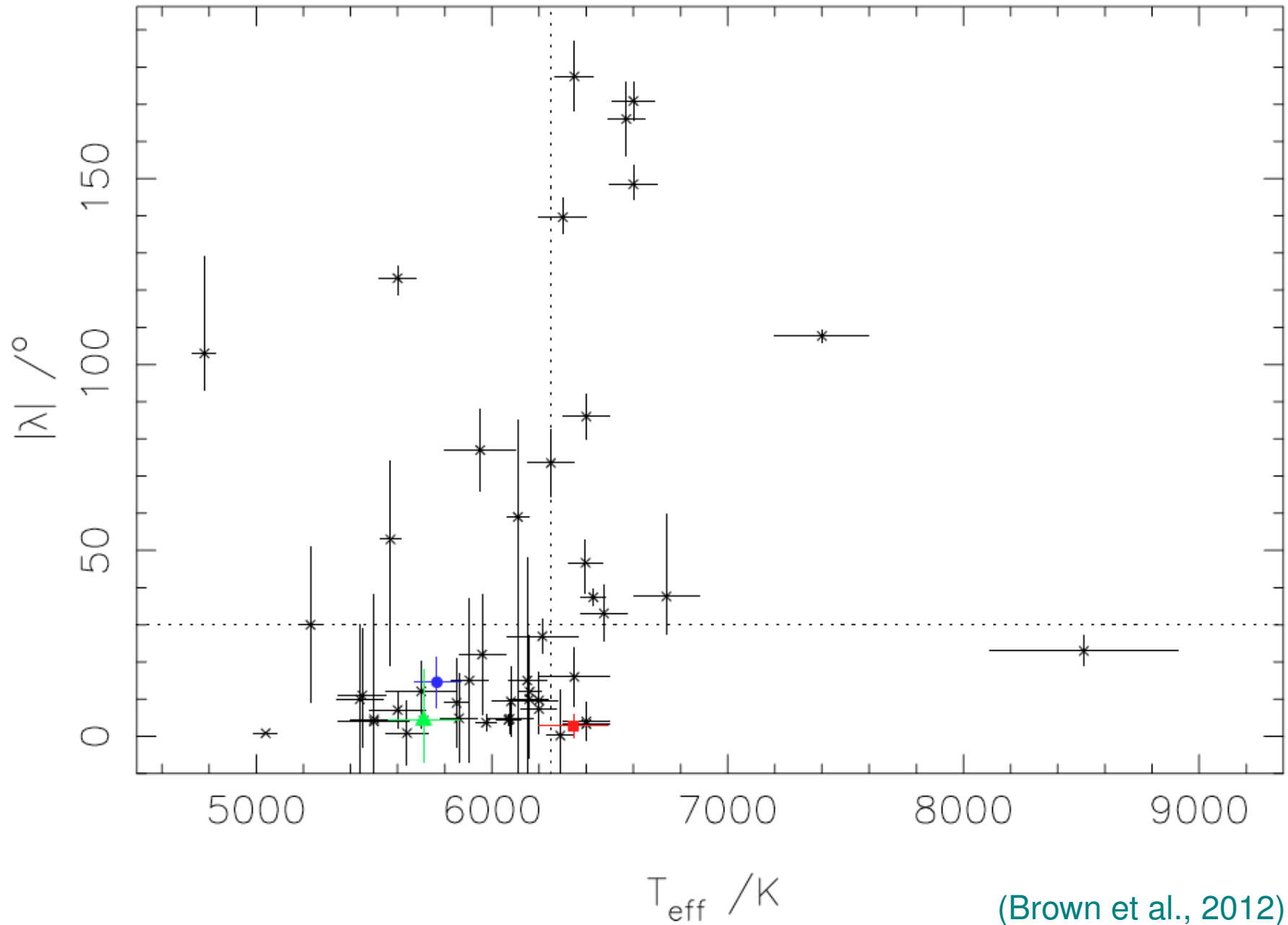
(Bitsch&Kley '10)

- e -damping for all planet masses.

Small e : exponential damping, large e : $\dot{e} \propto e^{-2}$

- Migration outward upto $e \approx 0.02 - 0.03$

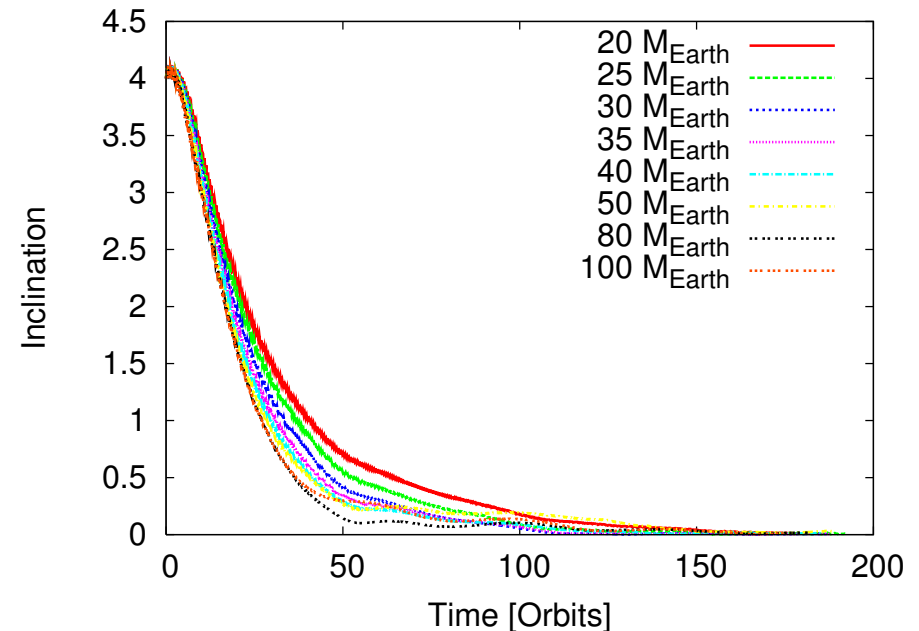
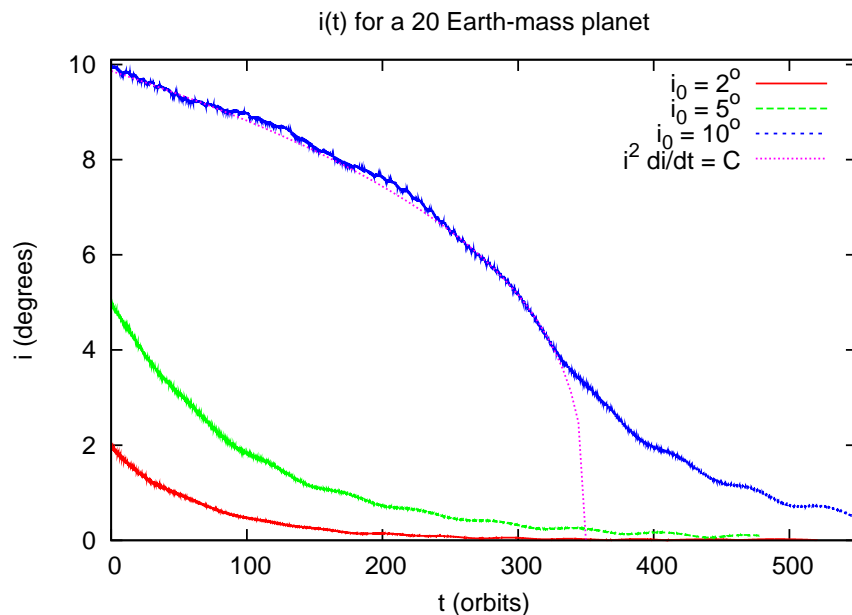
\implies Need multiple objects ! (and Scattering)





Fix planet mass $M_p = 20M_{\text{Earth}}$
- Vary initial Inclination

Vary Planet Mass 20 – $100M_{\text{Earth}}$
- Same $i_0 = 4\text{deg}$



(Cresswell et al 2007; Bitsch & Kley 2011)

- i -damping for all planet masses.

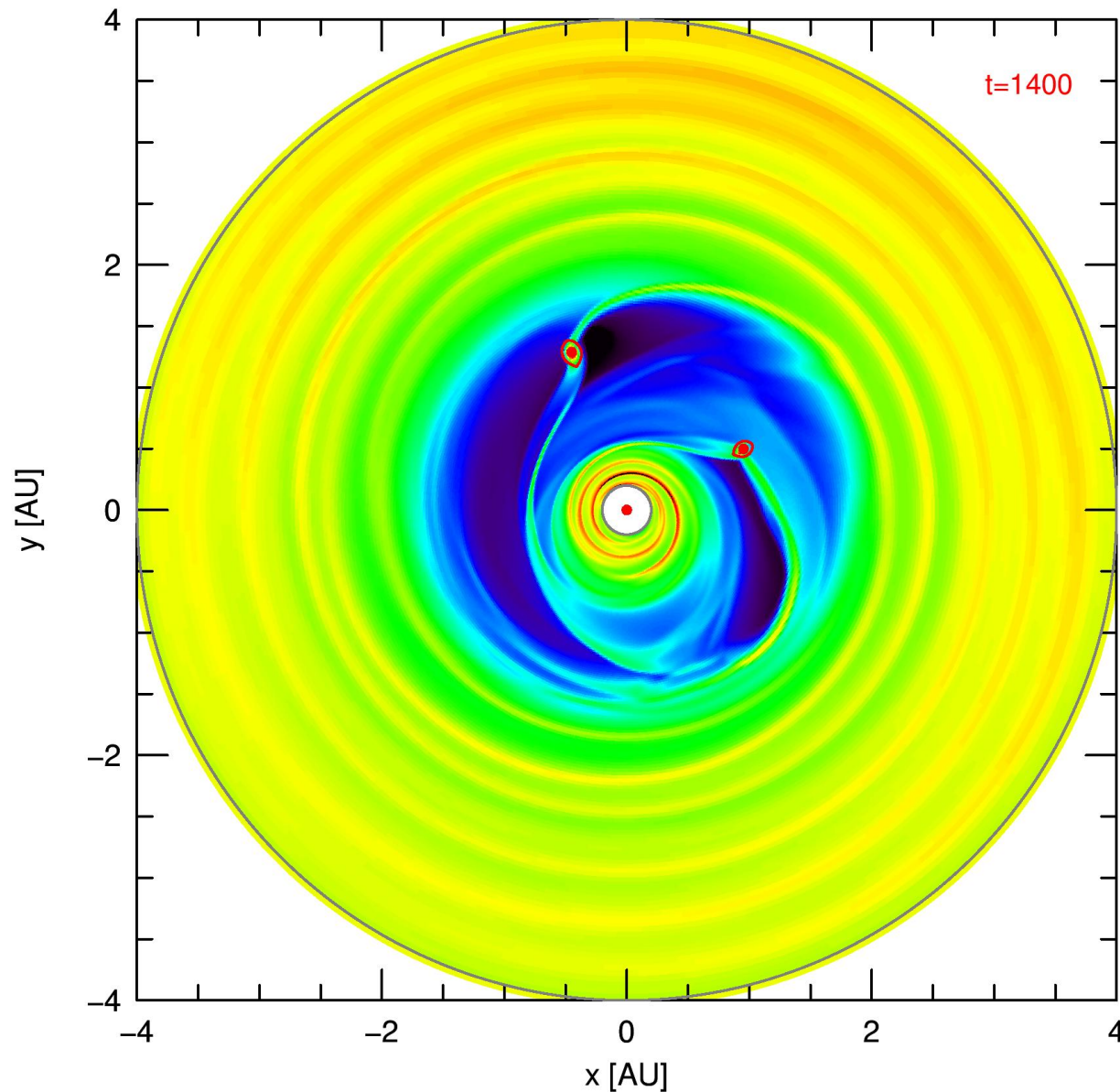
Small i : exponential damping, large i : $\dot{i} \propto i^{-2}$

- Migration still outward upto $i \approx 4^\circ$

⇒ Need multiple objects ! (Scattering)



2 massive Planets in disk



Two planets:
joint, large gap

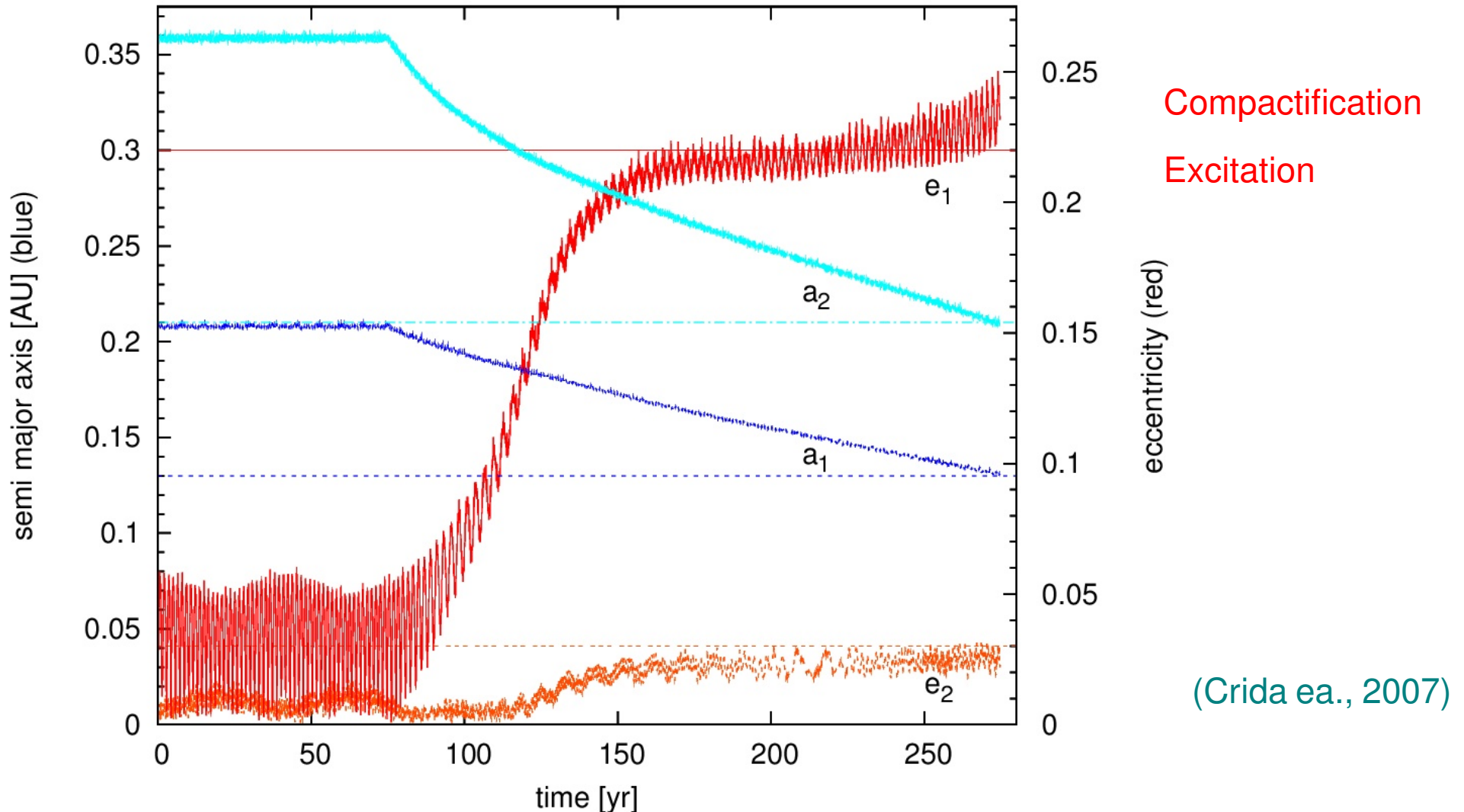
Outer planet :
Pushed inward

Inner planet :
Pushed outward

Separation reduction:
Resonant capture

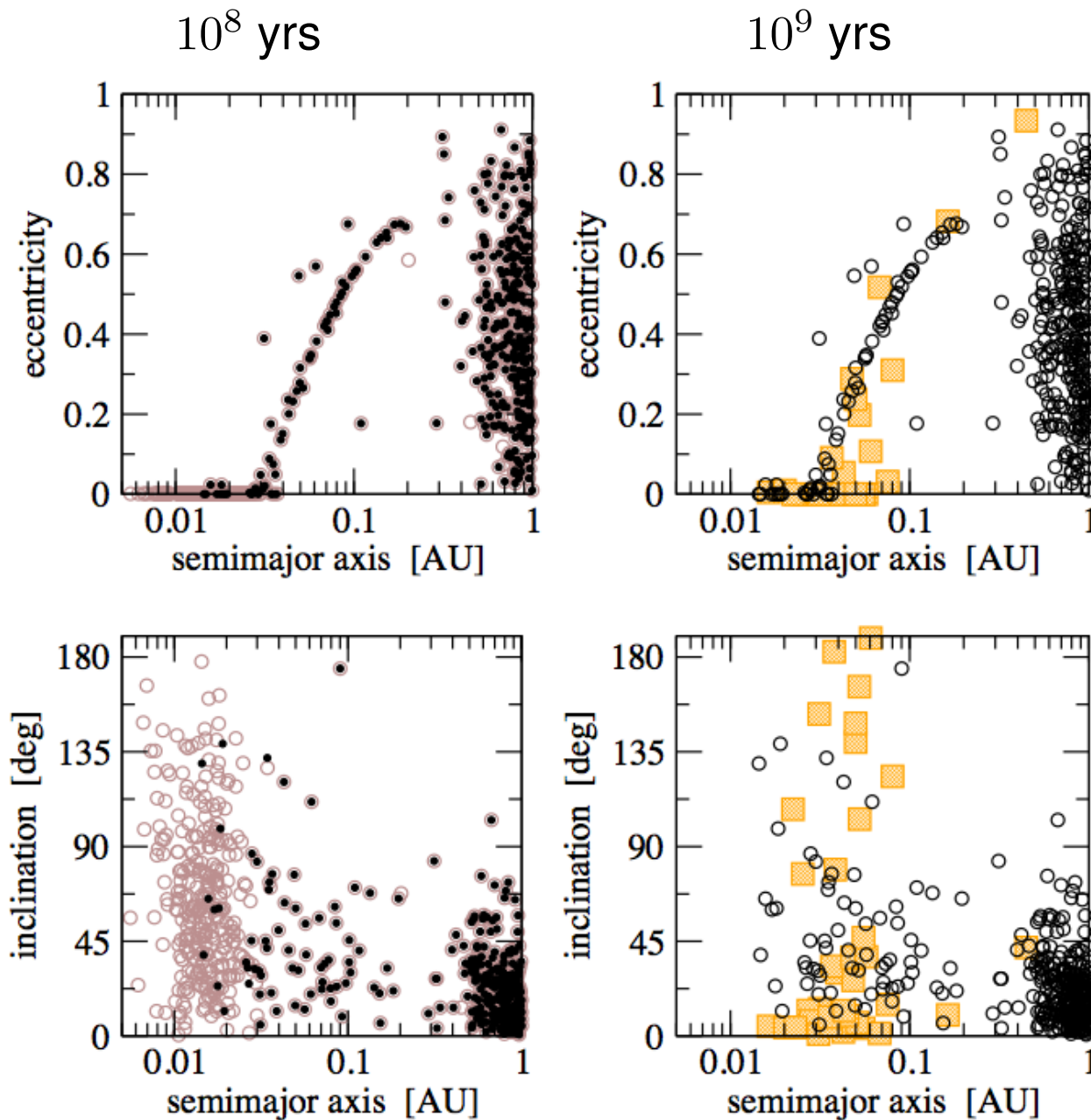


Here: **System-parameter of GJ 876** (2 planets in 2:1 resonance, 60:30 days)



System ends in: **apsidal corotation**, with **correct eccentricities**

Less disk damping: \Rightarrow much higher $e \Rightarrow$ Instability



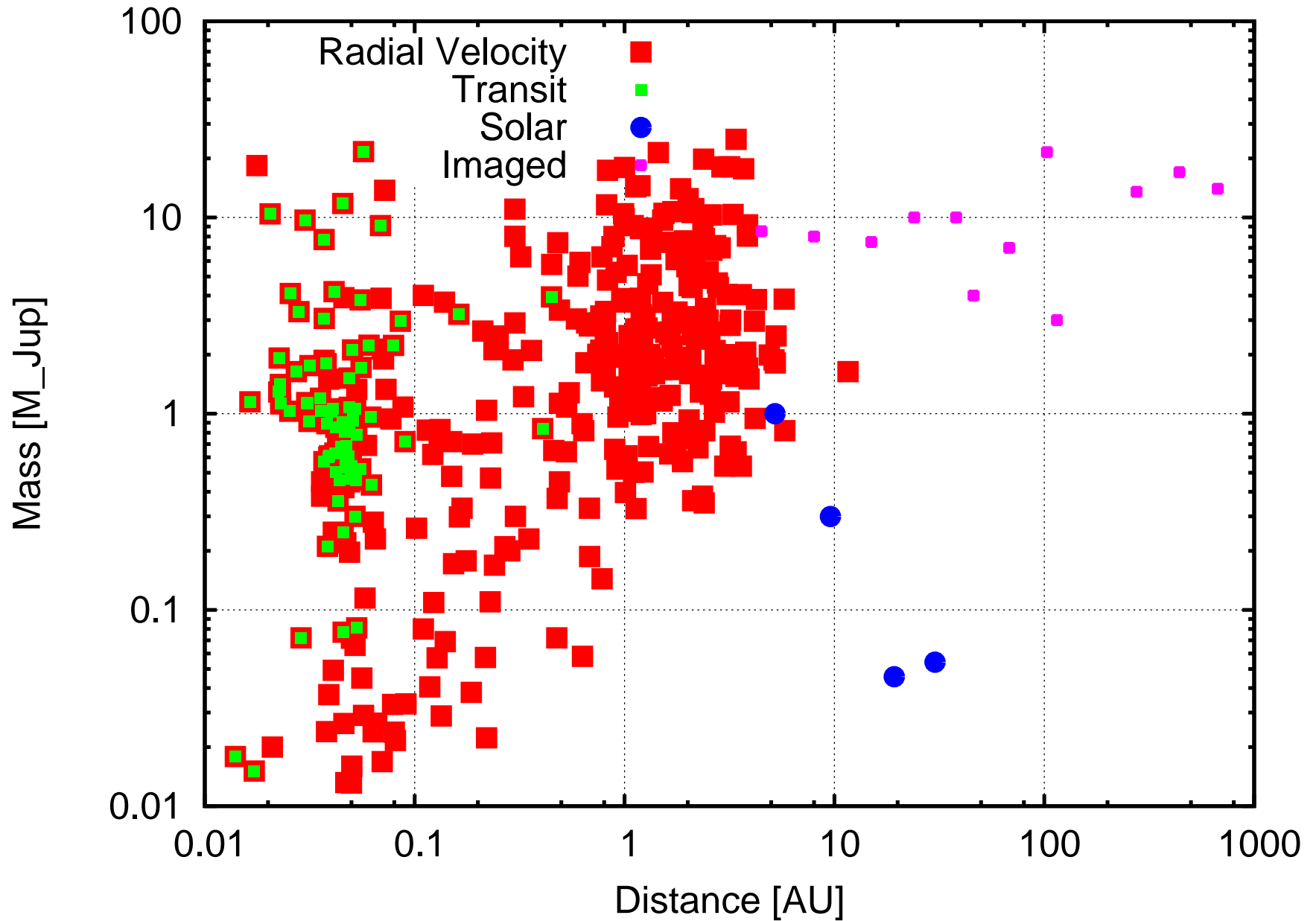
3 or 4 body simulations:
 Start with:
 - system close to resonance
 Include:
 - star-planet interaction

(Beauge & Nesvorny 2012)



Large distances (several hundred AU)

(Data: exoplanet.eu)





Consider local density perturbation in disk

Analytical

Stability-Criterion (Toomre)

$$Q \equiv \frac{c_s \kappa_0}{\pi G \Sigma_0} > 1$$

with

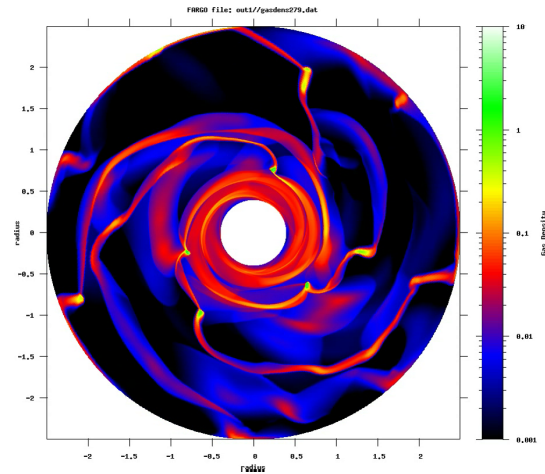
- c_s = sound velocity
- κ_0 = Epicyclic-Frequency (Ω_K)
- Σ_0 = surface density

- Pressure & Rotation stabilize
- Density destabilizes

Numerical

Evolution of an **isothermal** Disk

- Finite-Difference Hydrodynamics
- Viscous Disk



(Tobias Müller, Tübingen)

Disk heats up upon compression, need fast cooling

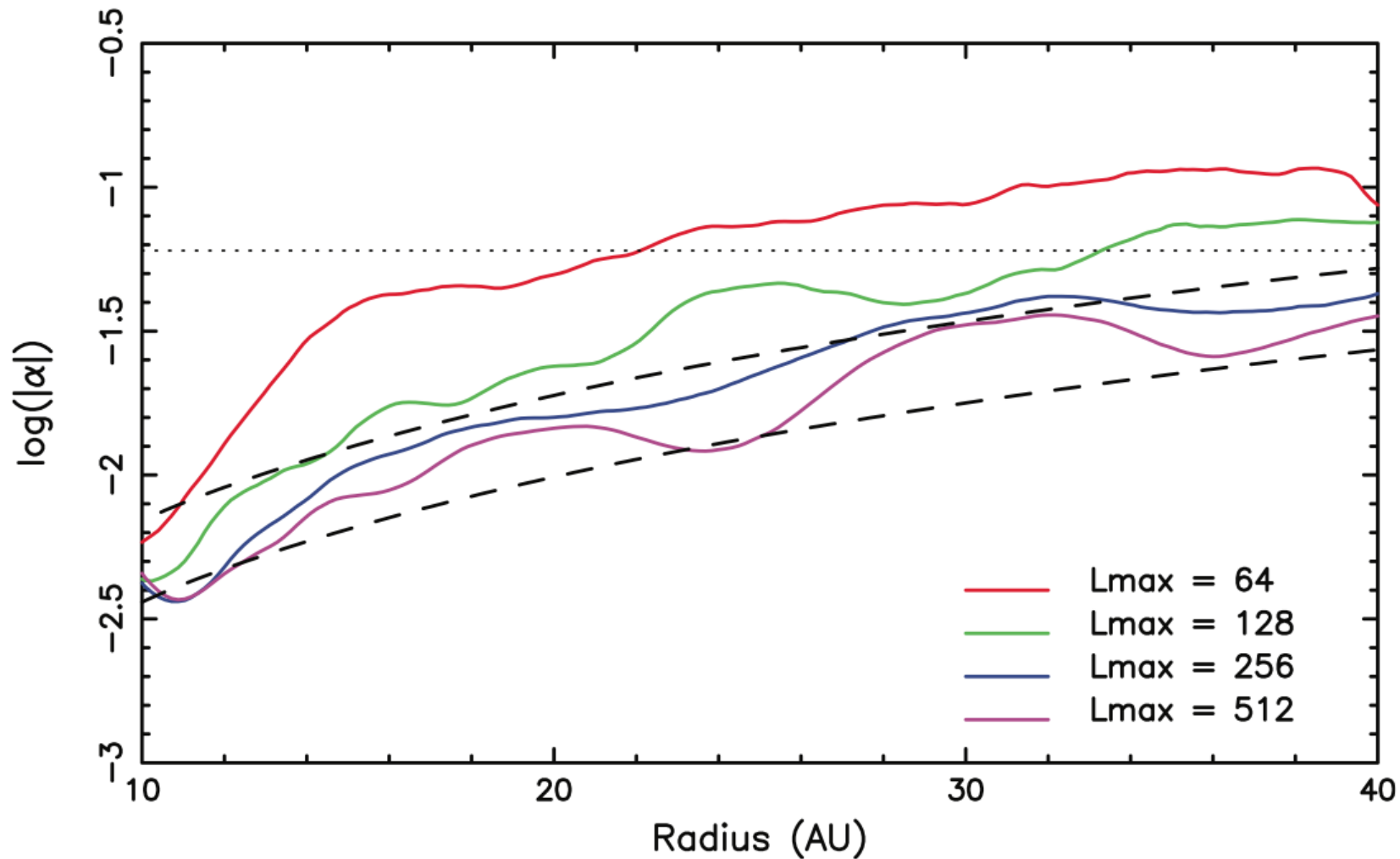
Require: **coolingtime** \approx **orbital period** \Rightarrow fast formation

Only in **large distances** from star (larger 40-50 AU, no cores)

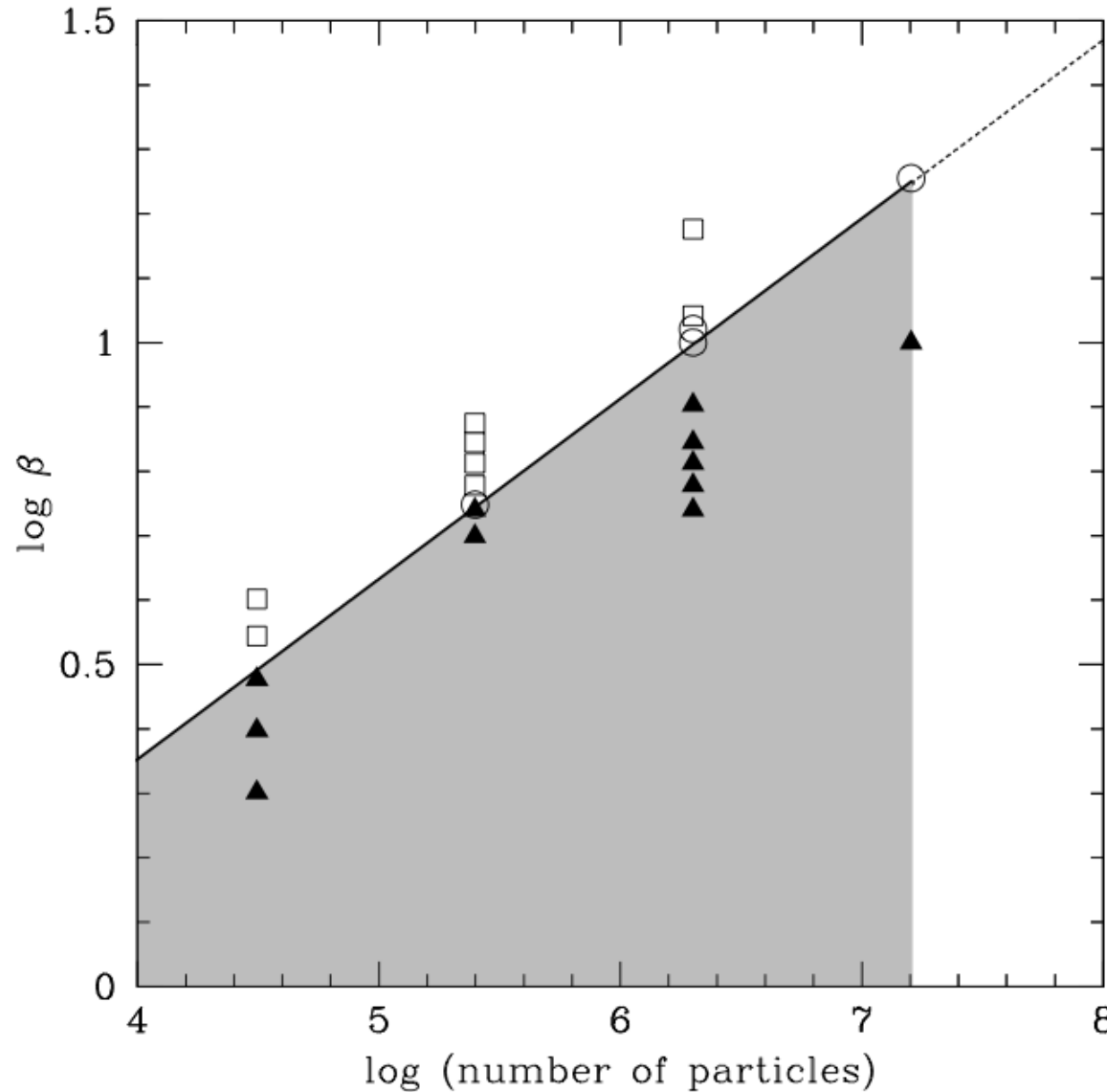


3D grid simulations - Gravitational stresses: $T_{r\varphi} \propto \int_{disk} g_r g_\varphi df$

- lower stresses for higher resolution



(Scott ea. 2012)



3D SPH simulations

$$\tau_{cool} = e_{th} / (de_{th}/dt)$$

$$\tau_{cool} = \beta \Omega^{-1}$$

Test Instability:

□ stable

▲ fragmenting

More fragmentation
in HighRes disks

(Meru & Bate 2010)



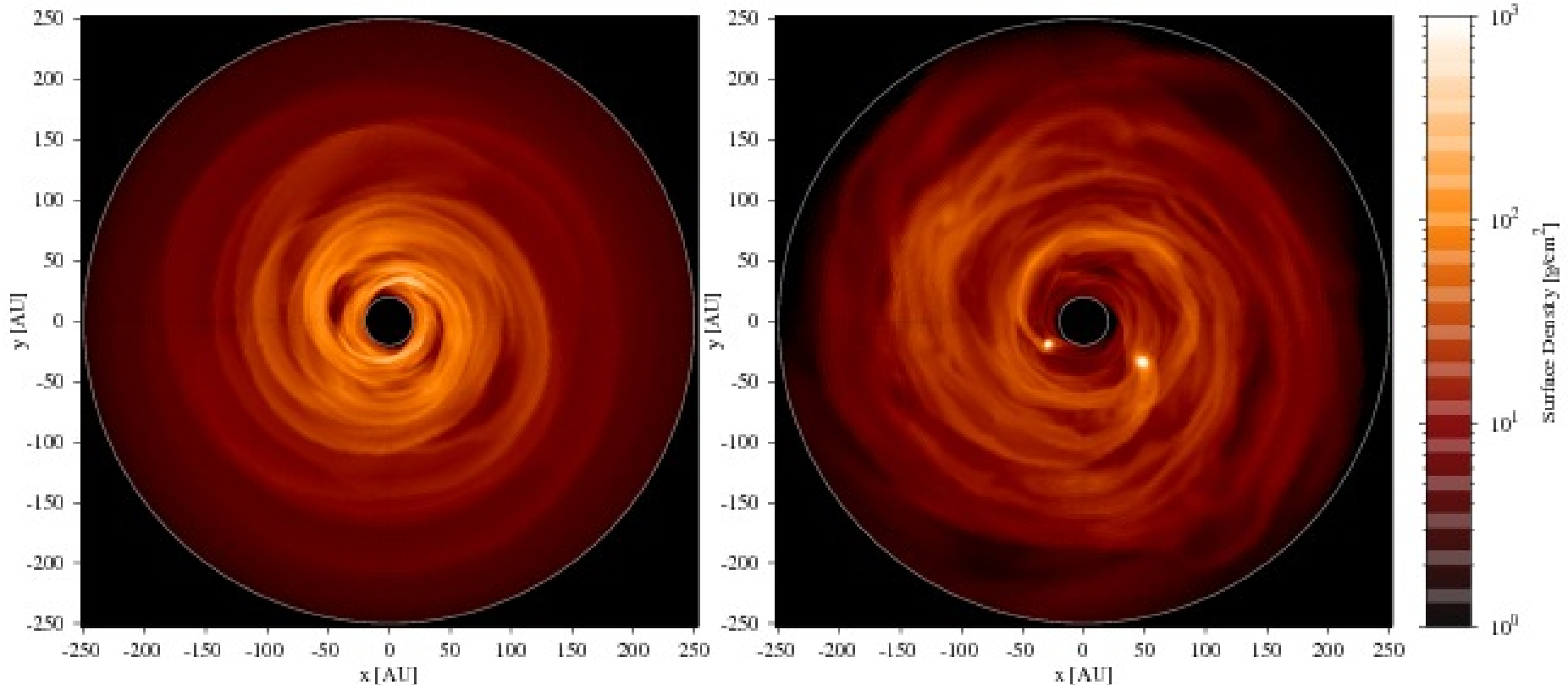
2D Grid simulations (FARGO) (smoothing takes vertical height into account)

$$\Psi \propto \frac{1}{(s^2 + \epsilon^2)}$$

$$\epsilon = 0.6H$$

$$\epsilon = 0.006H$$

(Müller, Kley & Meru 2012)

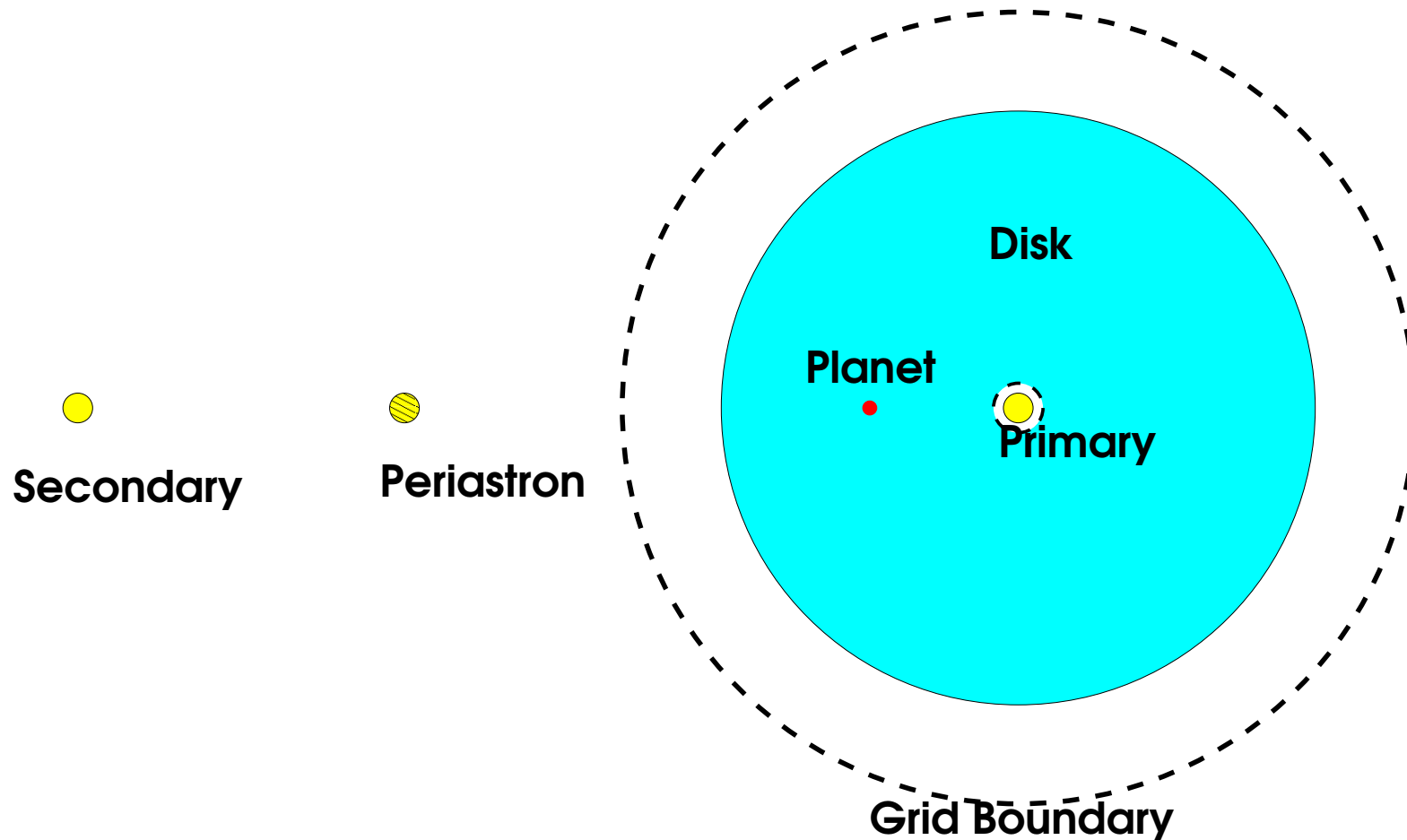


Realistic smoothing: $\epsilon \approx H \Rightarrow$ Less fragmentation



-
- Numerics
 - Cooling efficiency
 - Irradiation from star
 - Fate of fragments
 - tidal disruption
 - core formation
 - relation to FU Ori

Talk by: S. Nayakshin



Binary (γ Cephei)

$M_1 = 1.4 M_{\odot}$, $M_2 = 0.4 M_{\odot}$, $a_b = 20$ AU, $e_b = 0.4$

Grid: [0.5, 8.0] AU

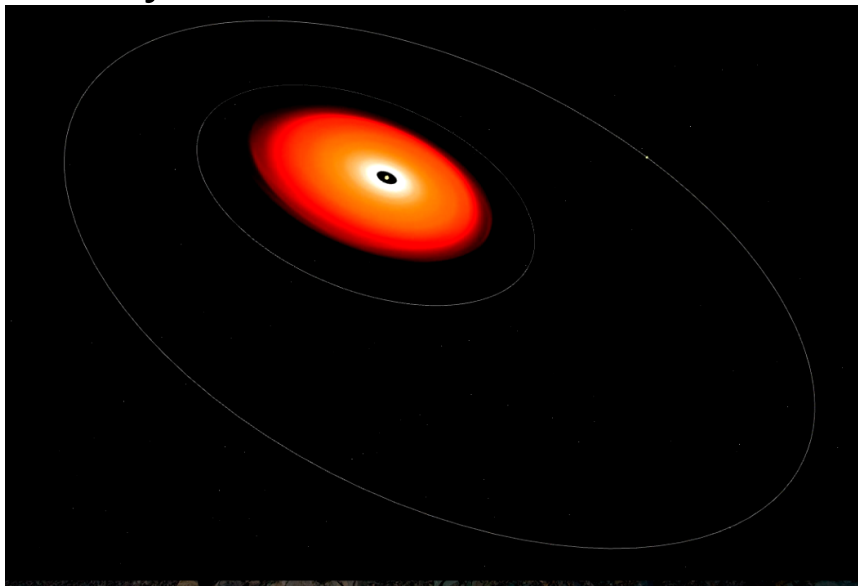


2D viscous accretion disk

locally isothermal

 $r_{\min} = 0.5, r_{\max} = 8.0$ AUGrid: 256×576

Density structure:



(Tobias Müller)

Strong dynamical interaction

 \Rightarrow Planet Formation more difficult
(A. Nelson, 2001)

both: sequ. acc. & GI

- Disk heating
- Enhanced collision velocities



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- Planets form in disks
 - flat systems
 - Planet-disk interaction moves planets
 - Inward for isothermal disks
 - possibly outward/slowed in **radiative disks**
 - Eccentricity & Inclination damped by disk
 - Eccentric & inclined planets through scattering
 - Distant planets through gravitational instability
 - Planet formation in binaries more difficult



Thank you for your attention !

(A. Crida)