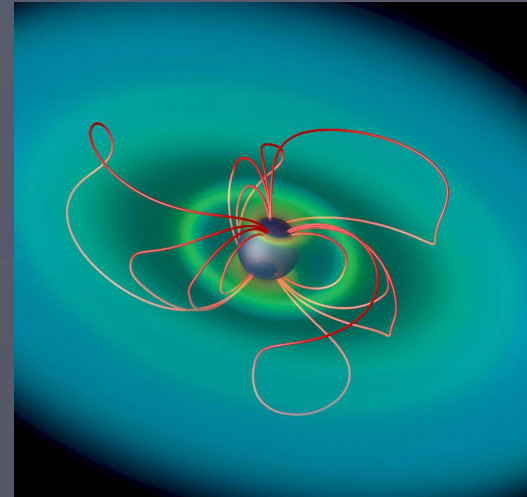


On some MHD aspects of Star-Disc-Jet systems

Jonathan Ferreira



Collaborators:

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P. Garcia, J. Bouvier, PO Petrucci, G Henri, G.
Pelletier

- I- Why should accretion discs host large scale B_z field ?
- II- Constraints from TTauri jets on accretion-ejection flows
- III- Naive views on star-disc interactions

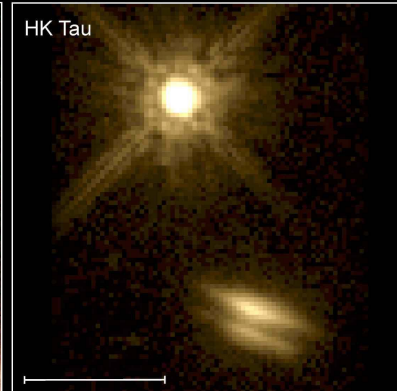
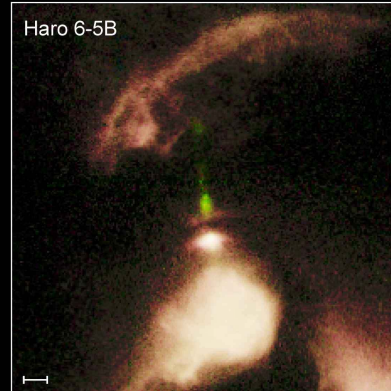
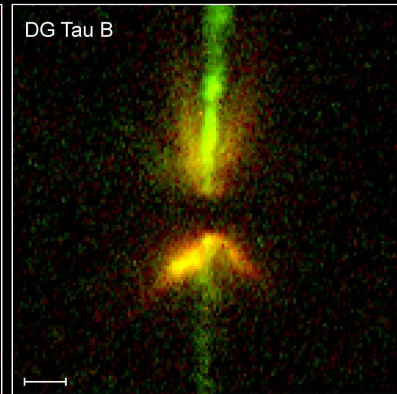
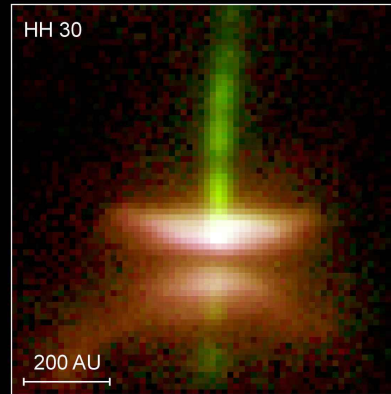
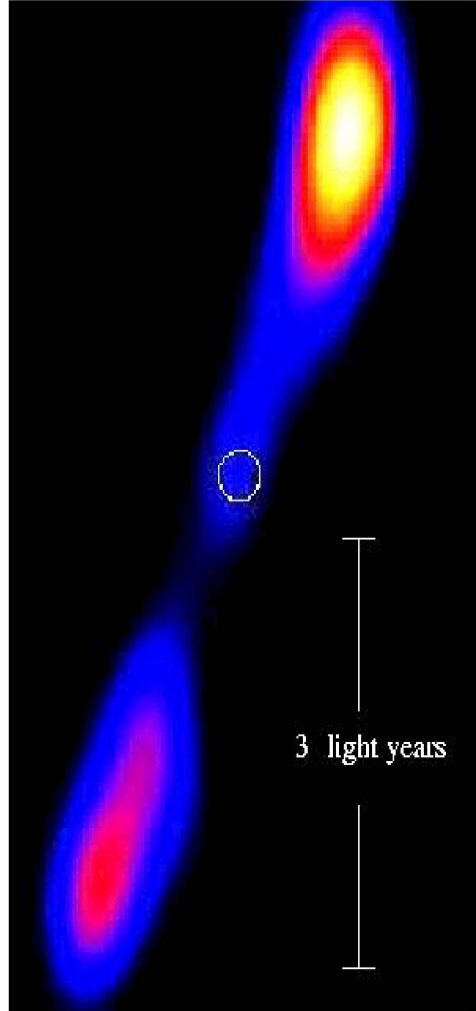
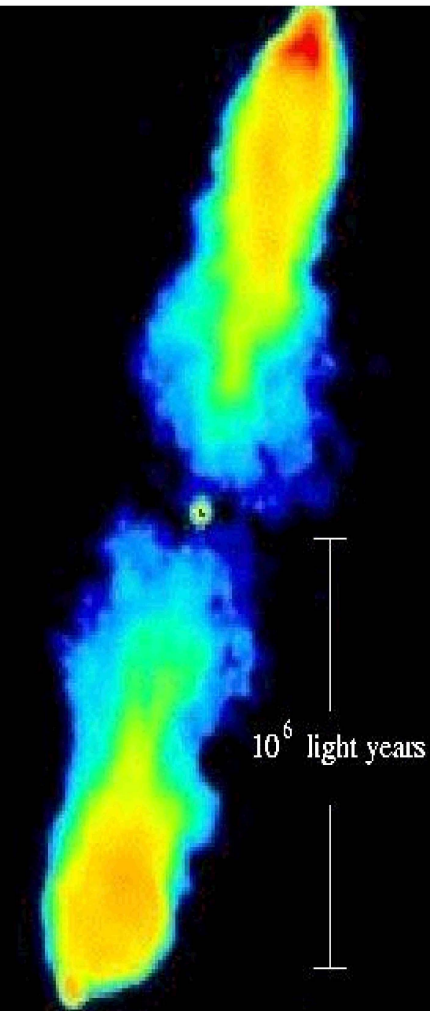
I-1 Astrophysical jets as disc winds

1- A « universal » model

2- Jets not always present

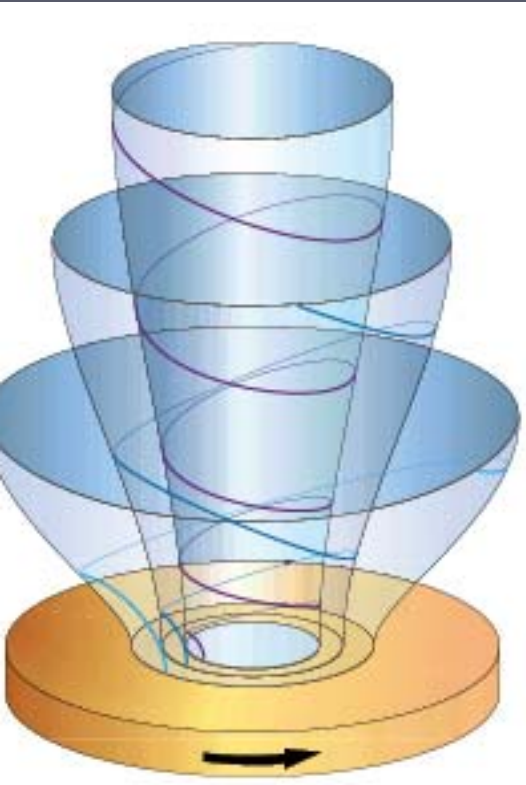
Quasar/radio galaxy

Microquasar 1E1740.7-2942



Disks around Young Stars
Hubble Space Telescope • WFPC2

I-2 Accretion-Ejection Systems



Axisymmetric jets are nested magnetic surfaces of constant magnetic flux: $a(r,z) = Cst$

Blandford & Payne 82

- Single fluid MHD description
- Non-relativistic equations (no light-cylinder)
- Steady-state
- Usually: polytropic energy equation
- Ideal MHD (no viscosity, no diffusivity)

A complex interplay between disc and jets

=> **M**agnetized **A**ccretion-**E**jection **S**tructure (MAES)

I-3 Some Ideas heard in conferences (and written too)

« Jets cannot be described by the Blandford & Payne picture »

Because:

- (1) There is never the correct bending at the disc surface
- (2) No large scale B_z can be maintained in the disc
- (3) In any case, accretion-ejection is by essence unstable

Jets may indeed be something else than what the accretion-ejection model depicts...but these ideas are wrong.

(1) « There is never the correct bending at the disc surface »

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Lubow *et al.* 94a

Heyvaerts, Priest & Bardou 96

Ogilvie & Livio 98

BP criterion for cold jets: $\theta > 30^\circ$ ($B_r > B_z$) at the disc surface.

This is never achieved in a **S**tandard **A**ccretion **D**isc (SAD). BP criterion requires a magn Reynolds number

$$\mathcal{R}_m = \frac{r u_r}{\nu_m} \geq \frac{r}{h}$$

In a SAD, « viscous » accretion gives

and in turbulent media $\nu_v \sim \nu_m \dots$

$$\mathcal{R}_e = \frac{r u_r}{\nu_v} \simeq 1$$

BUT the jet torque has been forgotten !

And it is dominant in a **J**et **E**mitting **D**isc (JED)...

(2) « No large scale B_z can be maintained in the disc »

Argument mainly based on previous calculations:

with $R_m \sim 1$, mass accretes whereas B_z stays behind... but because of imposed rigid BC!

Actually, a SAD can transport B_z such as to increase the disc magnetization towards the center :

Diffusion equation

$$v_m \frac{\partial B_z}{\partial r} \simeq u_r B_z$$

leads to

$$B_z \propto r^{-R_m}$$

Since

$$P_{\text{tot}} = \frac{\dot{M}_a \Omega_k^2 h}{6\pi v_v} \propto r^{-3/2-\delta}$$

where

$$h(r) \propto r^\delta$$

one gets

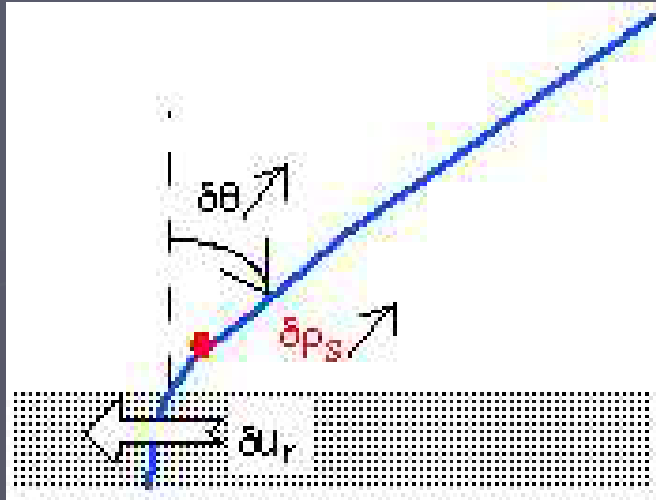
$$\mu = \frac{B_z^2}{\mu_o P_{\text{tot}}} \propto r^{-\varepsilon}$$

with $\varepsilon \sim 1$ for typical values for δ

Note: the only real issue is ionization and B/plasma coupling

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

(3) « In any case, accretion-ejection is by essence unstable »



Lubow *et al.* 94b
Ogilvie & Livio 01
Cao & Spruit 02

Calculations assumed a hydrostatic density profile.

This is a description far too crude : MHS quasi-equilibrium

- Disc rotation produces **larger** \mathbf{B}_ϕ if B_r becomes larger
- The magnetic field produces a strong vertical compression
- The SM point lies **below** the Sonic point (wrong mass flux)

=> as θ increases **less** mass ejected, not more!

Accretion-Ejection are stable (against this particular effect)

I-4 Semi-analytical studies of Accretion-Ejection systems

Wardle & Konigl 93, Ferreira & Pelletier 93,95
Ferreira 97, Vlahakis et al 00, Casse & Ferreira 00,0
Ferro-Fontan & Gomez 03, Campbell 99->05

QuickTime™ and a
TIFF (Uncompressed) decompressor
are needed to see this picture.

These studies address the mass load issue and the physical conditions within the disc required to drive super-fast jets.

=> The disc radial + vertical structures must be solved for.

Casse & Ferreira 00

Major results:

1- magnetic B_z field close to equipartition ($\mu \sim 1$)

2- high level of turbulence ($\alpha_m \sim 1$) required for stationarity

I-5 Numerical studies of Accretion-Ejection systems

Many numerical simulations where the disc is a boundary condition and mass loss imposed (eg. Ouyed & Pudritz 97,99,03, Ustyugova 99, Krasnopolsky et al 99, 03, Anderson et al 05...)

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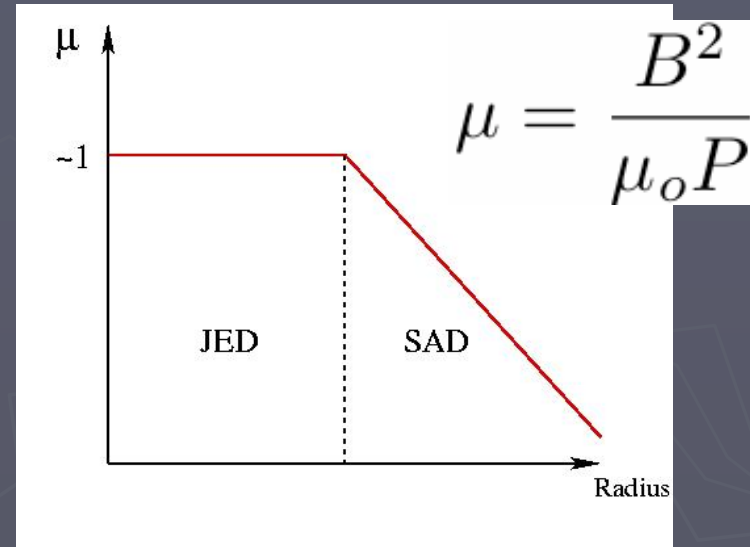
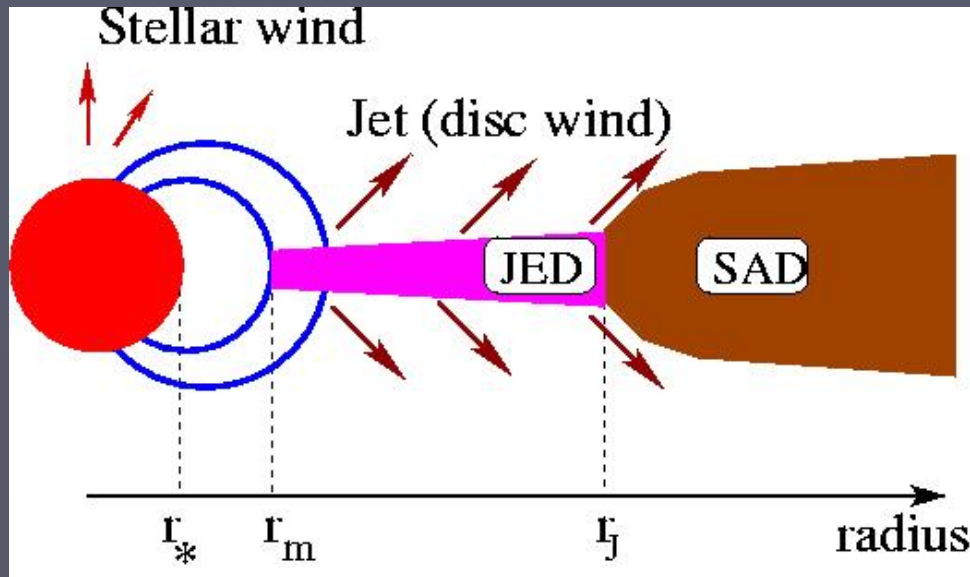
Casse & Keppens 02, 04

QuickTime™ and a
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Zanni et al 07

=> Main results from self-similar calculations are confirmed by MHD simulations where the disc is **also** computed.

I-6 A picture for the innermost disc regions



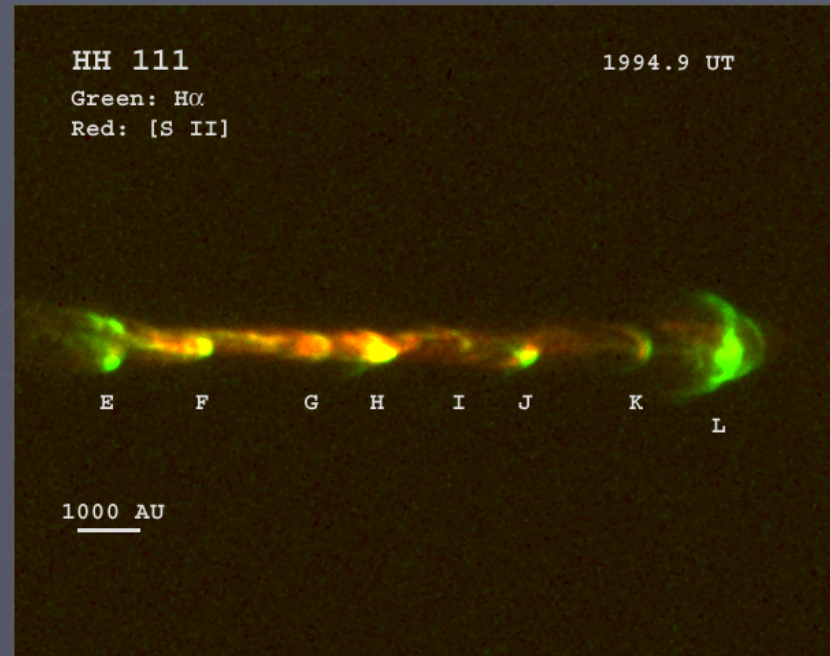
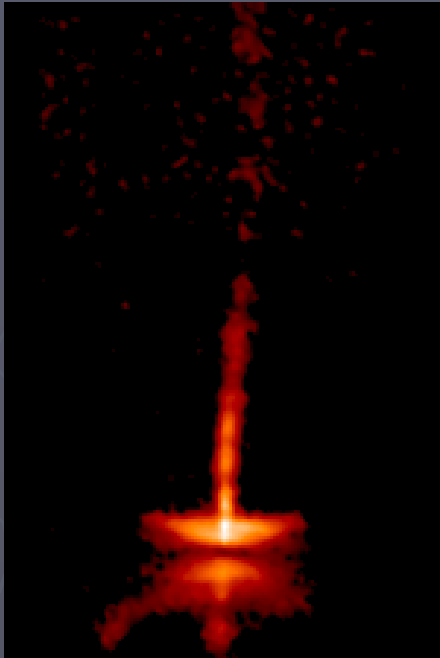
Jet Emitting Disc with \sim equipartition large scale B_z field

$$B_z \simeq 0.2 \left(\frac{M}{M_\odot} \right)^{1/4} \left(\frac{\dot{M}_a}{10^{-7} M_\odot / \text{yr}} \right)^{1/2} \left(\frac{r_o}{1 \text{ AU}} \right)^{-5/4 + \xi/2} \text{ Gauss}$$

Ferreira & Pelletier 95

\Rightarrow Existence of such B_z recently confirmed by spectro-polarimetric observations around FU Or (\sim kG @ 0.05 AU) Donati et al 05, Nature 10

II-1 Constraints from Ttauri Jets

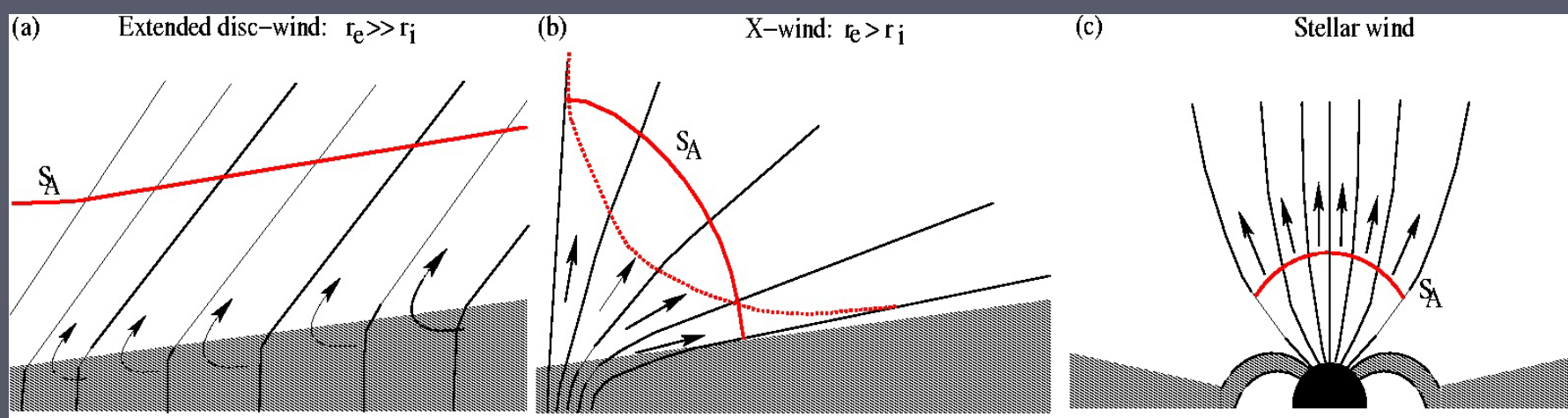


Images: degree of collimation, evolution (HH objects)

Spectroscopy: jet kinematics (line profiles, PV diagrams) and physical conditions such as density and temperature (line ratios).

=> **Strong constraints on all MHD models**

II-2 The 3 basic steady-state jet models



Blandford & Payne 82

Ferreira & Pelletier 93,95,97

Wardle & Königl 93

Casse & Ferreira 00, 04

Shu et al 94, 95

Fendt 9, 00

Shang et al 98, 02

Weber & Davis 68

Hartmann & McGregor 80

DeCampli 82

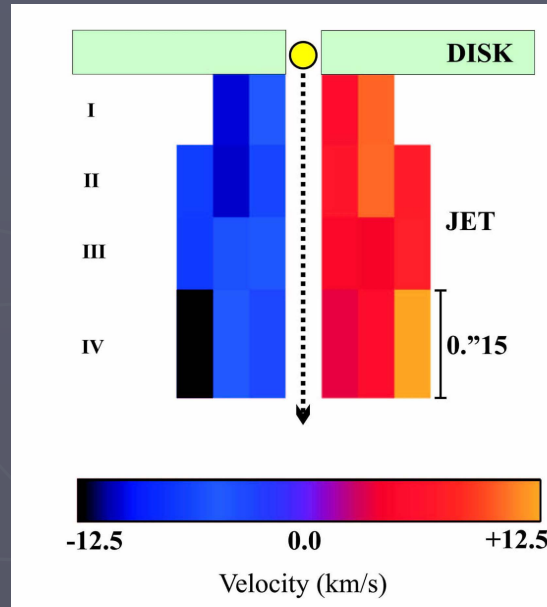
Sauty & Tsinganos 94, 02

- share the same physics: rotating body + large scale B_z
 - Governed by the same set of MHD equations
 - Apart the « extended disc wind » model, mass flux is imposed
- => Can observations discriminate between these models ?

II-3 First: are jets indeed rotating?

DG Tau
Observations
HST/STIS

Bacciotti et al 00

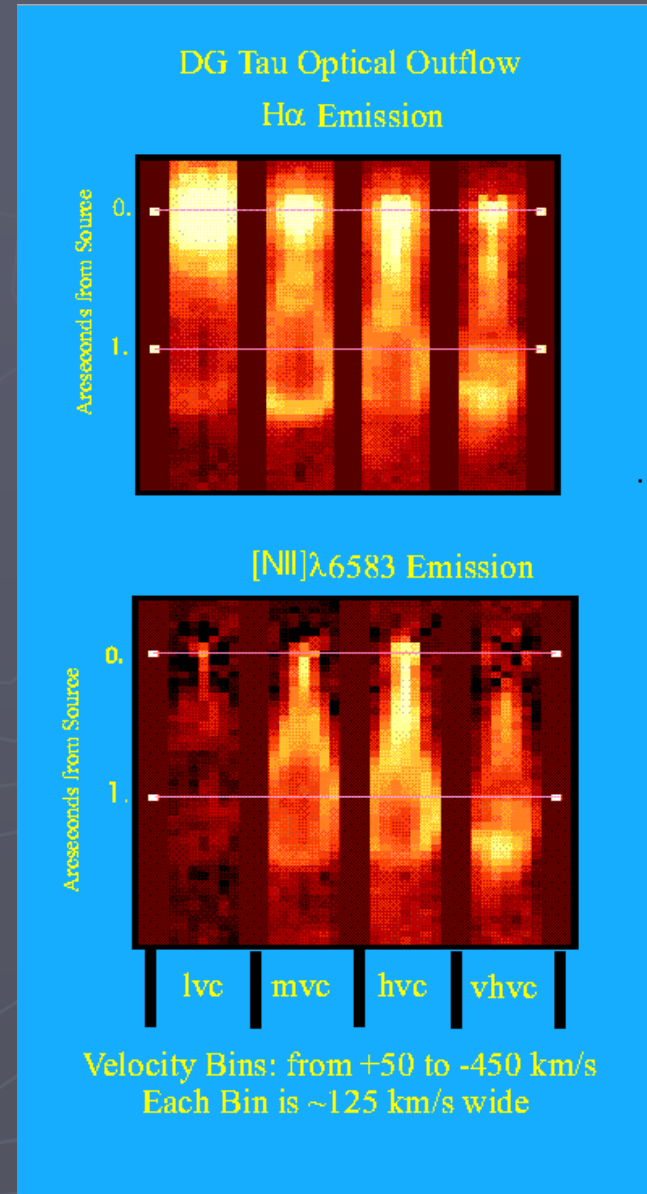


1. The collimation degree increases with the jet speed (higher closer to the axis)

2. Unresolved acceleration scale <20 au

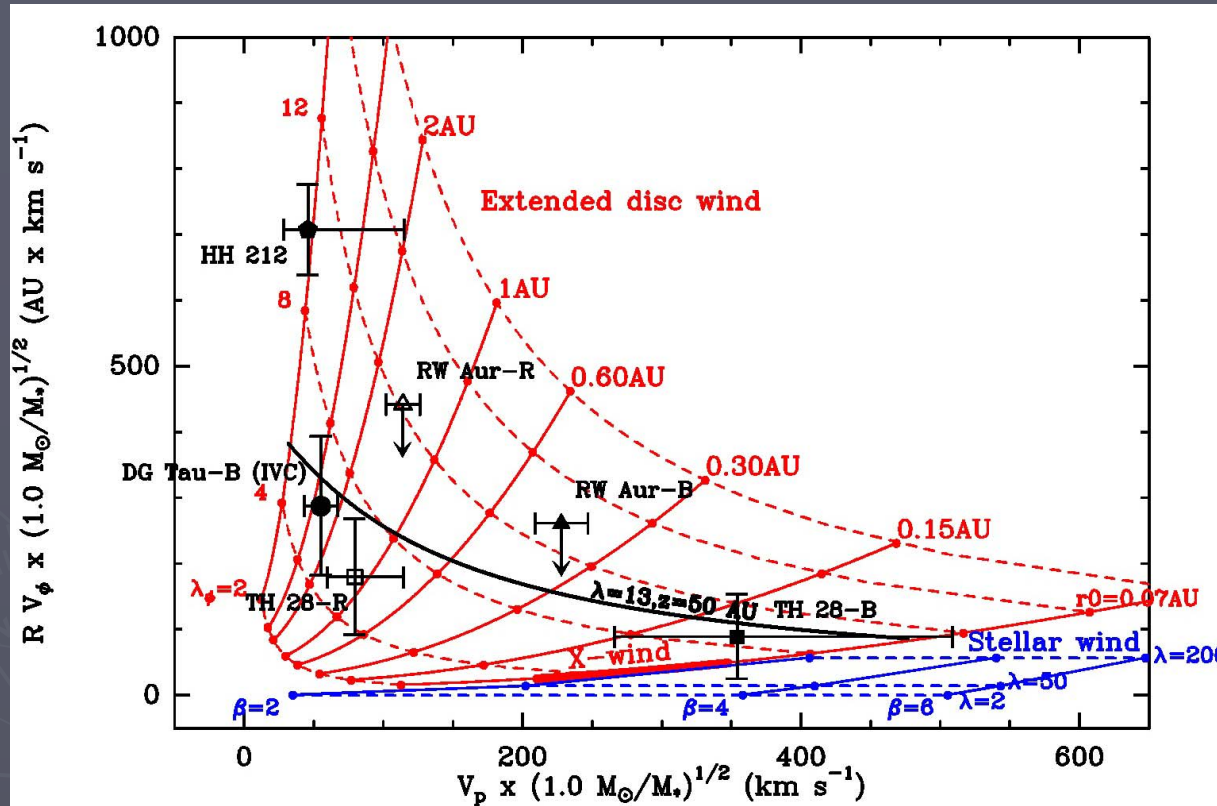
=> MHD: launching radius at 2-3 au

(Anderson et al 03, Pesenti et al 04)



II-4 Putting all constraints together

Ferreira *et al.* 06b



Observations:

- jet radius r
- velocity V_{ϕ}
- velocity V_p

Models:

- anchoring radius r_0
- magnetic lever arm

$$\lambda \simeq \frac{r_A^2}{r_0^2}$$

$\Rightarrow \lambda \sim 10$ needed

- Jet velocity gradients incompatible with current X-wind models
- Observed mass fluxes incompatible with stellar winds only.
- **IF** velocity shifts are rotation: launching radius from 0.2 to 3 AU

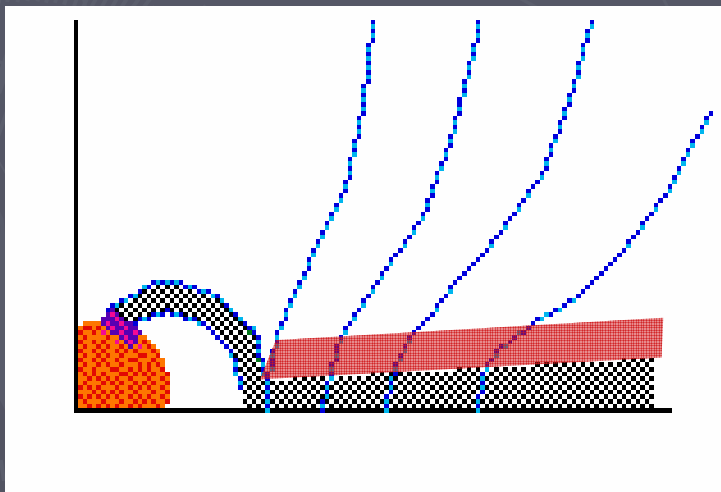
II-5 The need for a warm « corona »

Thin discs => enthalpy negligible in jets: « cold jet » models

In that case, $\lambda \sim 50-100$ (Ferreira 97, Casse & Ferreira 00a)

- Mass fluxes too low $\frac{2\dot{M}_{jet}}{\dot{M}_{acc}} \simeq 1/\lambda < 10\%$ observed
- Velocities too high

Higher mass fluxes with $\lambda \sim 10$ can only be done if some energy is deposited at the disc surface (Casse & Ferreira 00b).



Warm jets from thin discs can arise if

1. **Stellar UV/X illumination** (accretion shock)
2. **Local dissipation of accretion energy** (coronal heating, Galeev 79, Heyvaerts & Priest 89)

II-6 Can X-rays do the job?

$$\tau_X = 0.1$$

QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

$$\tau_X = 1$$

SM

$$\tau_X = 10$$

QuickTime™ and a
TIFF (Uncompressed) decompressor
are needed to see this picture.

$$10^{-6} M_{\odot}/\text{yr}$$

$$10^{-7} M_{\odot}/\text{yr}$$

$$10^{-8} M_{\odot}/\text{yr}$$

For relevant T Tauri star parameters (Macc, X-ray flux and spectrum, Imanishi et al 03, Wolk et al 05, Guedel 05) X-rays are unable to **heat** the base of the jet to the level required to enhance the mass flux.

→ *heating must be of turbulent origin*

Garcia et al, to be subm.

II-7 Black Hole XrBs: a similar situation?

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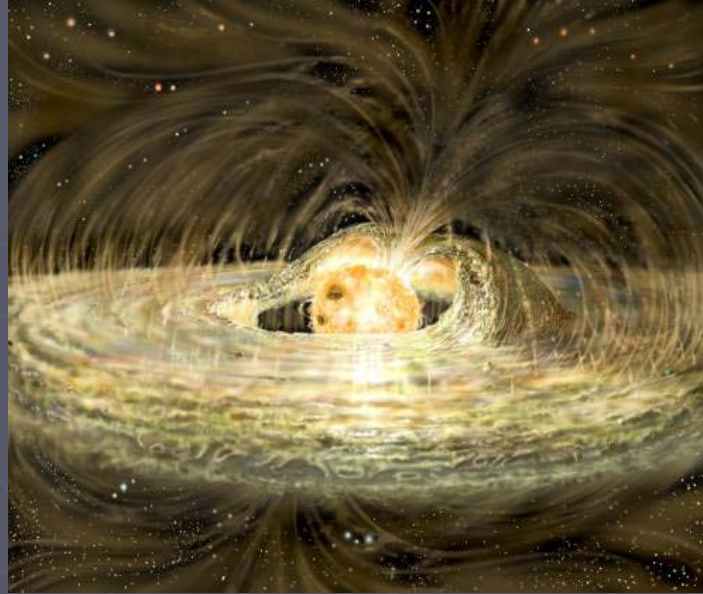
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Ferreira et al 06a, A&A

Here, a corona is needed in order to account for the hard X-ray spectrum when the disc emission is lacking *and jets observed.*

=> Consistent with JEDs

III-1 The star-disc magnetic interaction



See review in PPV
Alencar et al

T Tauri stars are slow rotators despite contraction+accretion (Bertout et al 89, Bouvier et al 97, Rebull et al 02)

=> Accretion must help to remove stellar angular momentum

Evidences of a magnetospheric interaction (Edwards *et al.* 94, 98, Calvet 04, Muzerolle et al 01, Bouvier *et al.* 99, Günther et al 99, Johns-Krull et al 99,01, Feigelson & Montmerle 99)

=> What kind of magnetic configuration?

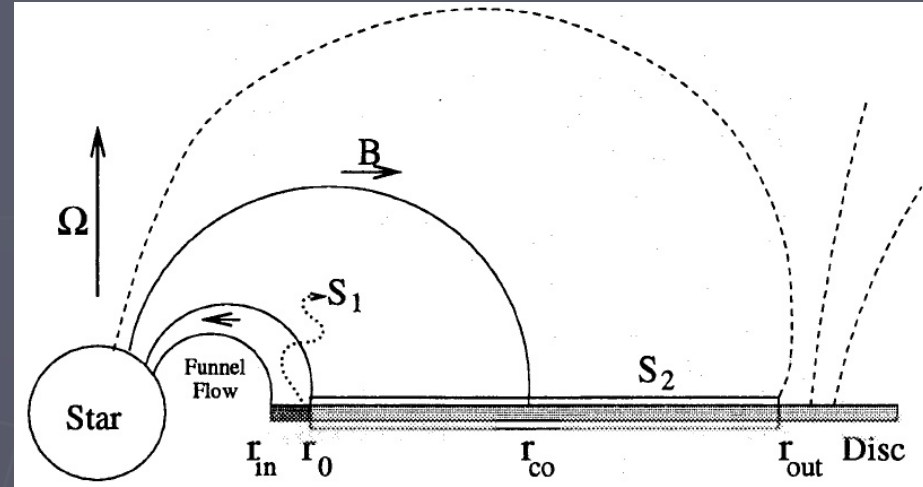
III-2 Two contradictory requirements

- (1) The disc-locking paradigm
- (2) Accretion must still proceed

Gosh & Lamb 79

Cameron & Campbell 93, Li 96

Matt & Pudritz 04



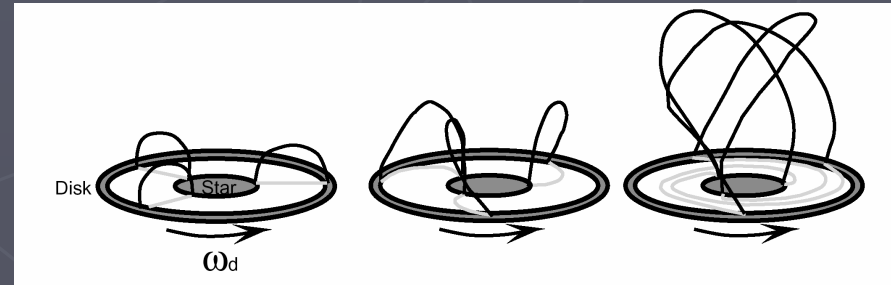
$-r_{in} > r_{co}$: star is spun down \Rightarrow accretion is prevented

$-r_{in} < r_{co}$: star is spun up \Rightarrow accretion is allowed

\Rightarrow Efficiency of disc viscosity??

Second process: reconnection

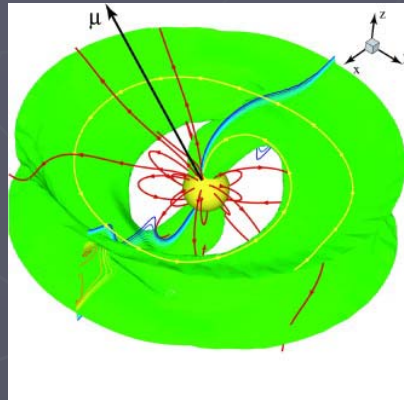
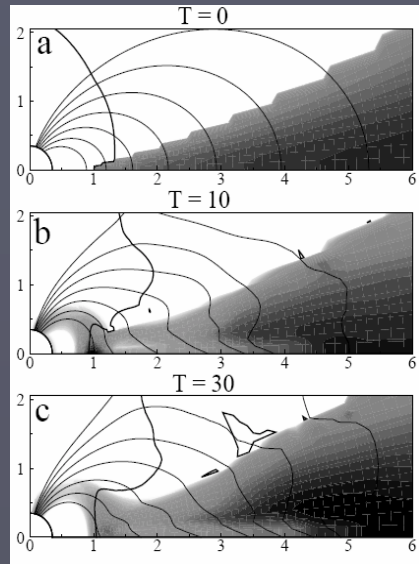
\Rightarrow Efficiency of disc magnetic diffusivity??



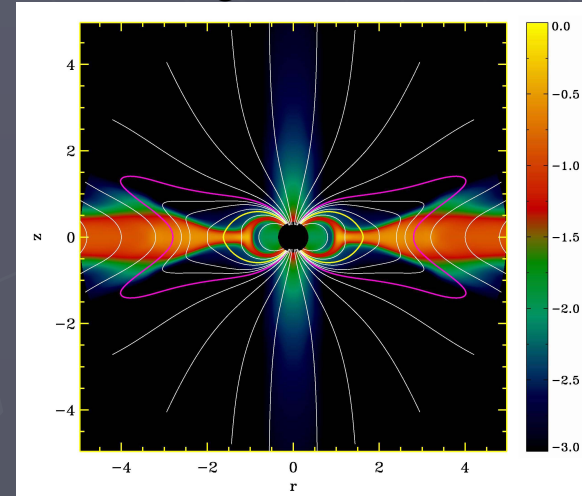
Lovelace et al 95, 99, Bardou & Heyvaerts 96
Uzdensky et al 02, Matt & Pudritz 05

III-3 Heavy numerical simulations

Hayashi et al 96, Miller & Stone 97, Hirose et al 97, Goodson et al 97,99, Kueker et al 03, Romanova et al 02,03,04,05, Long et al 05, 06 von Rekowski & Brandenburg 04,05



QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.



Zanni et al, in pre

Funnel flows are a natural feature in a non force-free magnetosphere, **BUT stellar spin up.**

Main difficulties (and differences):

- Treatment of disc physics (mass, v_v and v_m)
- Boundary conditions at the star (!!)
- Numerical resolution

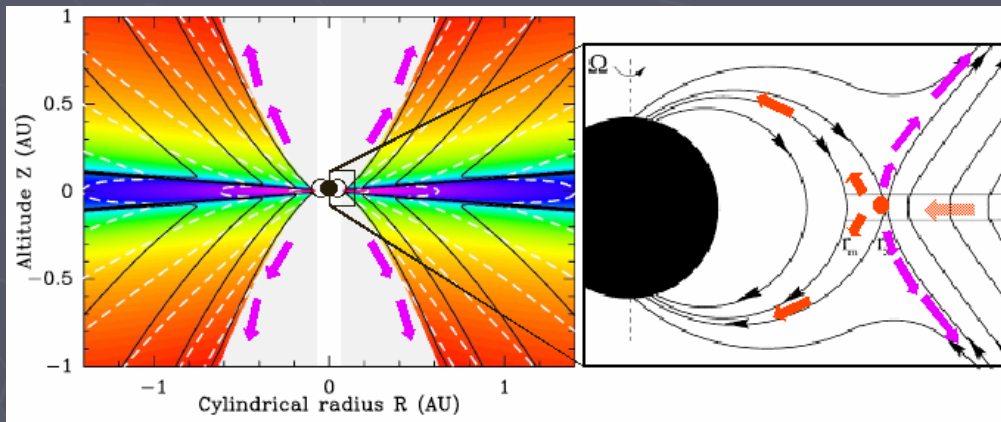
QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.

Concluding remarks

The « stellar angular momentum problem » requires a wind as a sink:

⇒ Accretion-powered stellard winds ?

⇒ Reconnection X-winds ?



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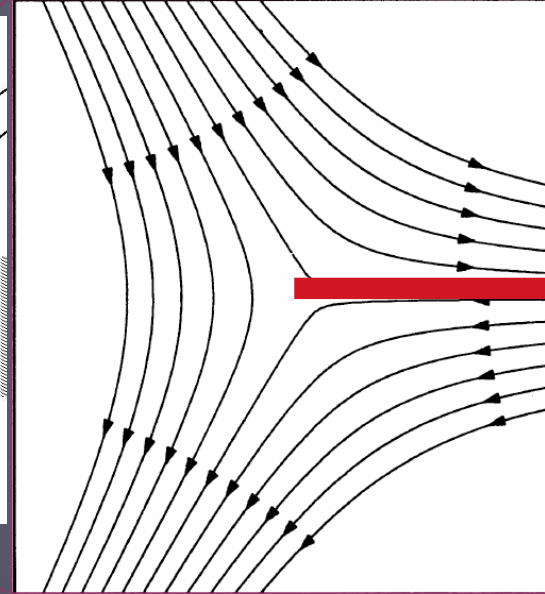
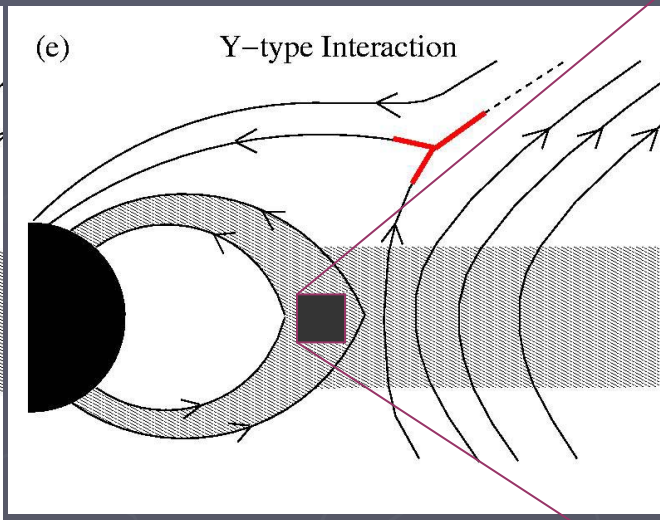
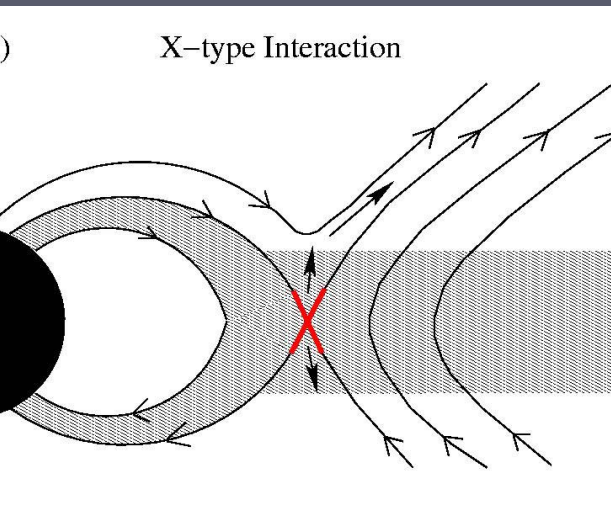
Matt & Pudritz 05

Ferreira, Pelletier & Appl 00

The answer probably relies on **MHD experiments**:

- Quality is greatly and rapidly improving over the years
- BUT the outcome is strongly dependent on the disc microphysics and turbulence (transport coefficients, corona)

Two (over-)simple configurations



Ferreira, Pelletier, Appl 00

Shu et al 94a

Camenzind 90

Matt & Pudritz 05

Ostriker & Shu 95

Both configurations give birth to a magnetic neutral line at the equator : chondrules (Gounelle et al 06)

-Reconnection X-winds: due to oppositely directed fields

-X-wind: unspecified origin

II-5 Testing models against observations

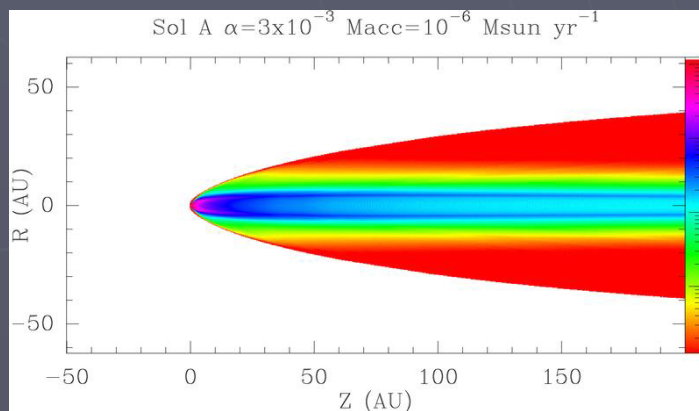
MHD model
 V, ρ, B

+

Thermal and Ionization Structures
 T, n_H, x_e

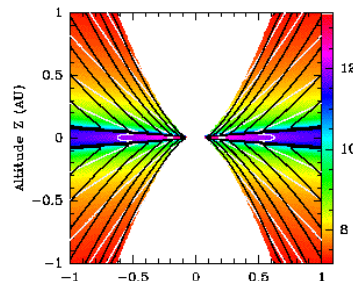
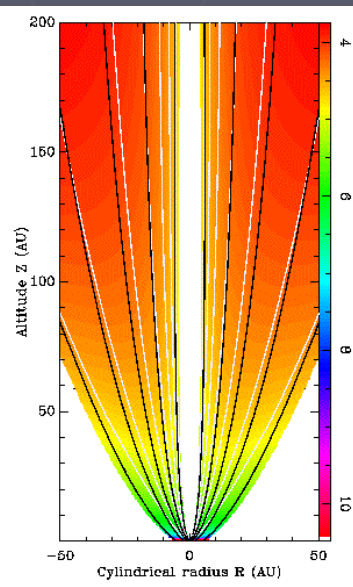
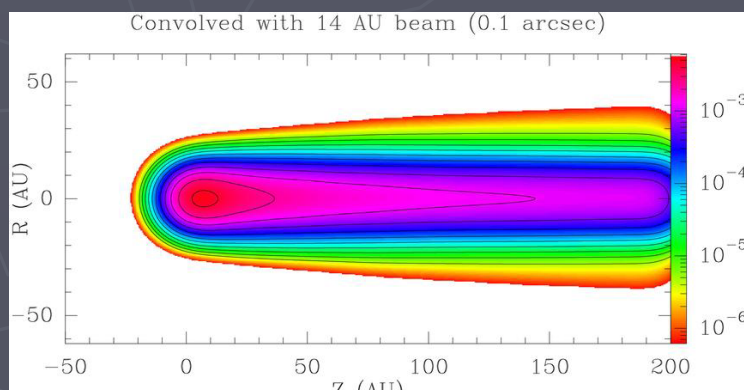
Projection

Emissivity map



Convolution

Synthetic image



QuickTime™ and a
TIFF (LZW) decompressor
are needed to see this picture.