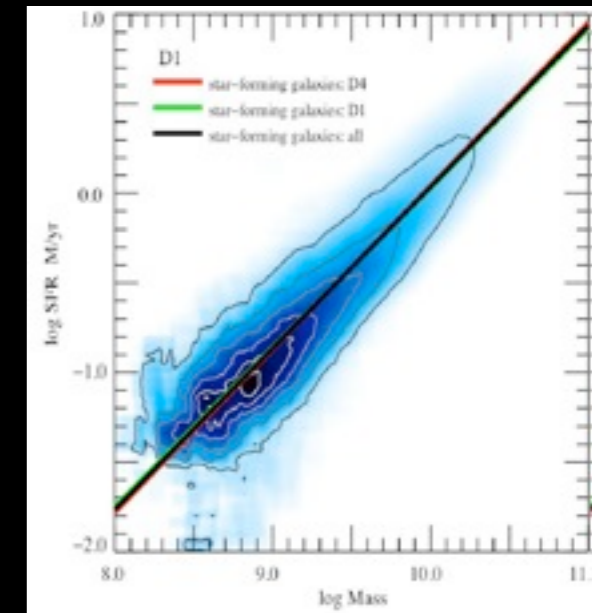
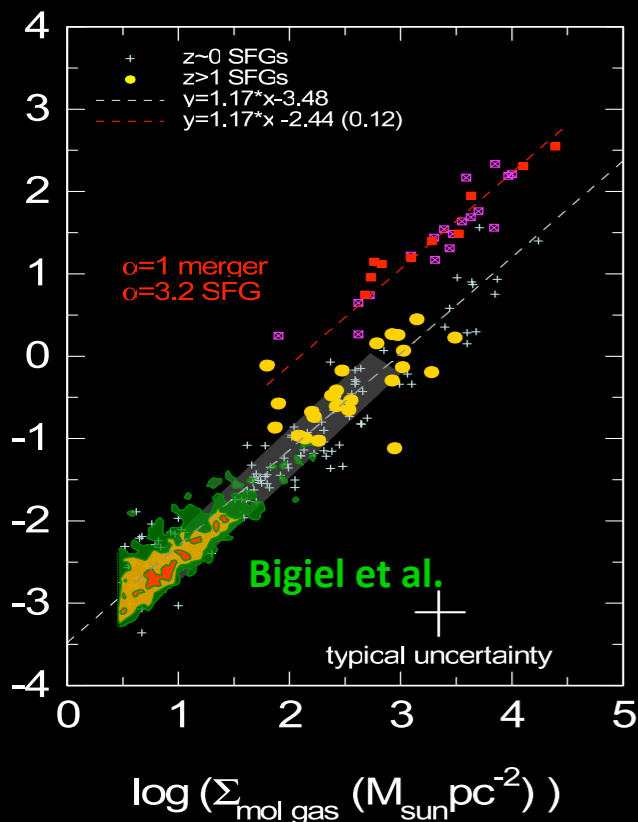
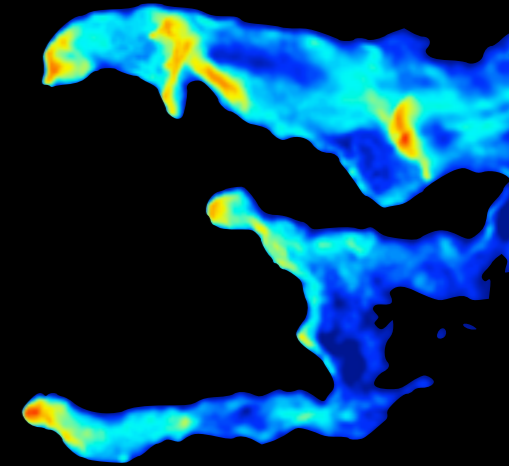


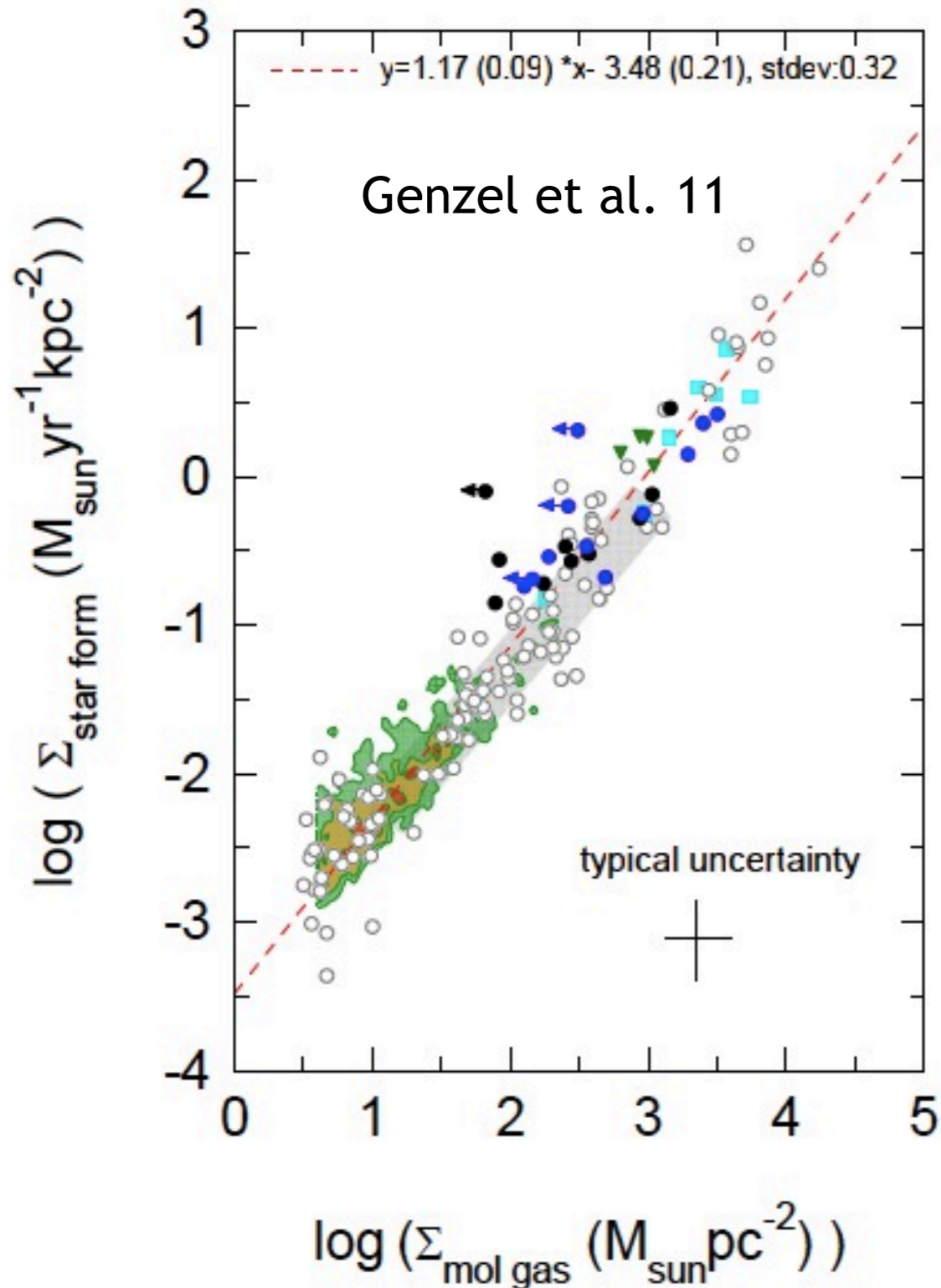
# Self-regulated star formation

Andreas Burkert



C. Dobbs, E. Ntormousi, K. Fierlinger,  
J. Ngoumou, J. Pringle, S. Walch

# Evidence for self-regulation



$$SFR = \frac{M_{H_2}}{\tau_{sf}} \text{ with } \tau_{sf} \approx 1 - 2 \cdot 10^9 \text{ yrs}$$

- $\tau_{sf}$  is almost independent of redshift
- Gas depletion timescale **50 times** greater than local free-fall timescale.

$$\tau_{ff} \ll \tau_{sf} < \tau_{\text{Hubble}}$$



continuous replenishment

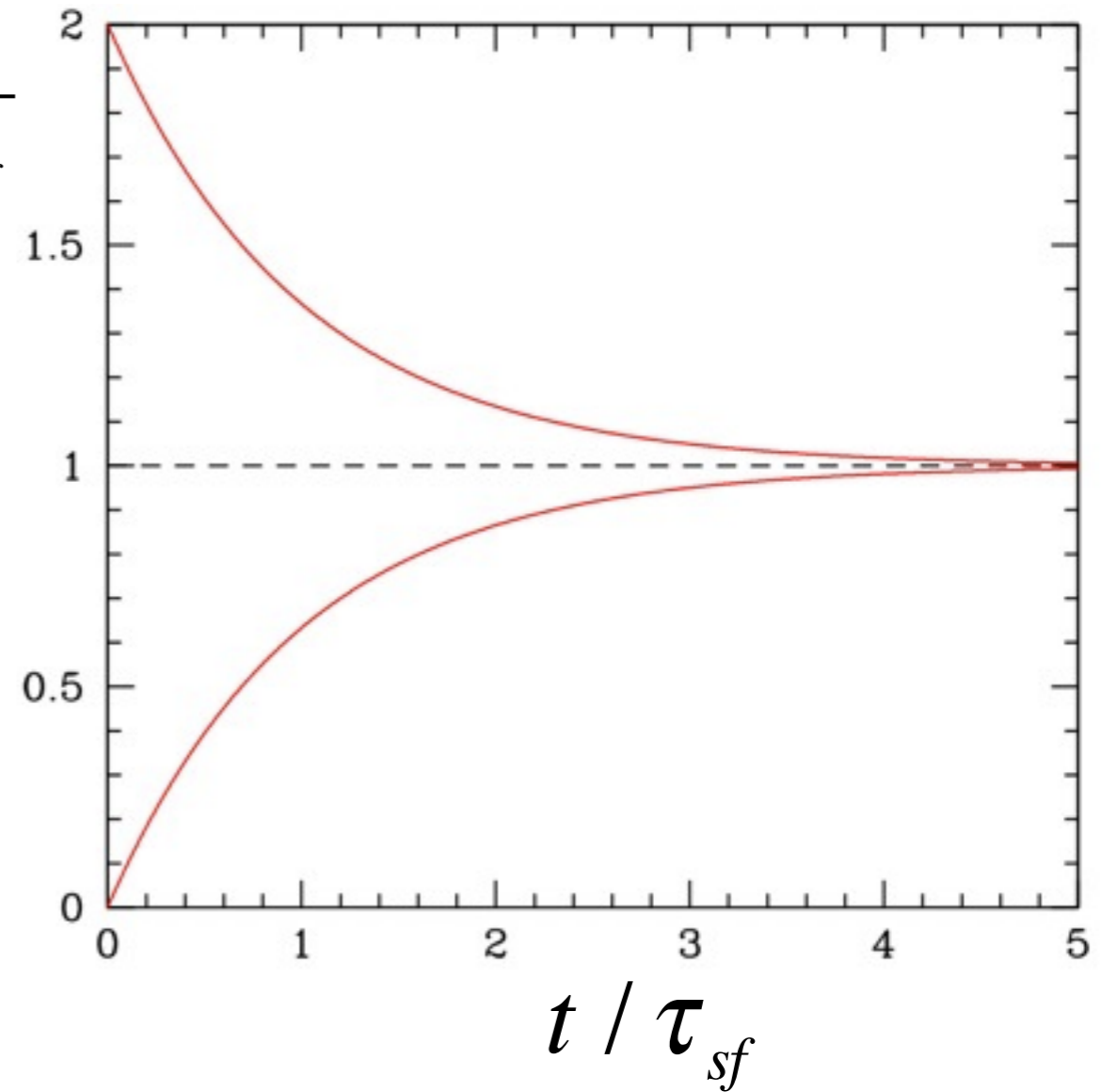
Bouché et al. 07, McKee & Ostriker 08, Genzel et al. 10,11, Daddi et al. 10

# What determines SFR?

(Bouche et al. 10; R. Davé et al. 11a,b)

$$\frac{dM_g}{dt} = \left( \frac{dM_g}{dt} \right)_{acc} - \frac{M_g}{\tau_{sf}} (1 - R + \alpha_{wind})$$

$$\frac{SFR}{\dot{M}_{acc,eff}}$$



$$\dot{M}_{acc,eff}$$

$$SFR = \frac{M_g}{\tau_{sf}} = \frac{1}{1 - R + \alpha_{wind}} \left( \frac{dM_g}{dt} \right)_{acc} \left( 1 - \exp\left( -\frac{t}{\tau_{sf}} \right) \right)$$



$$SFR = \dot{M}_{acc,eff}$$

# What determines SFR?

(Bouche et al. 10; R. Davé et al. 11a,b)

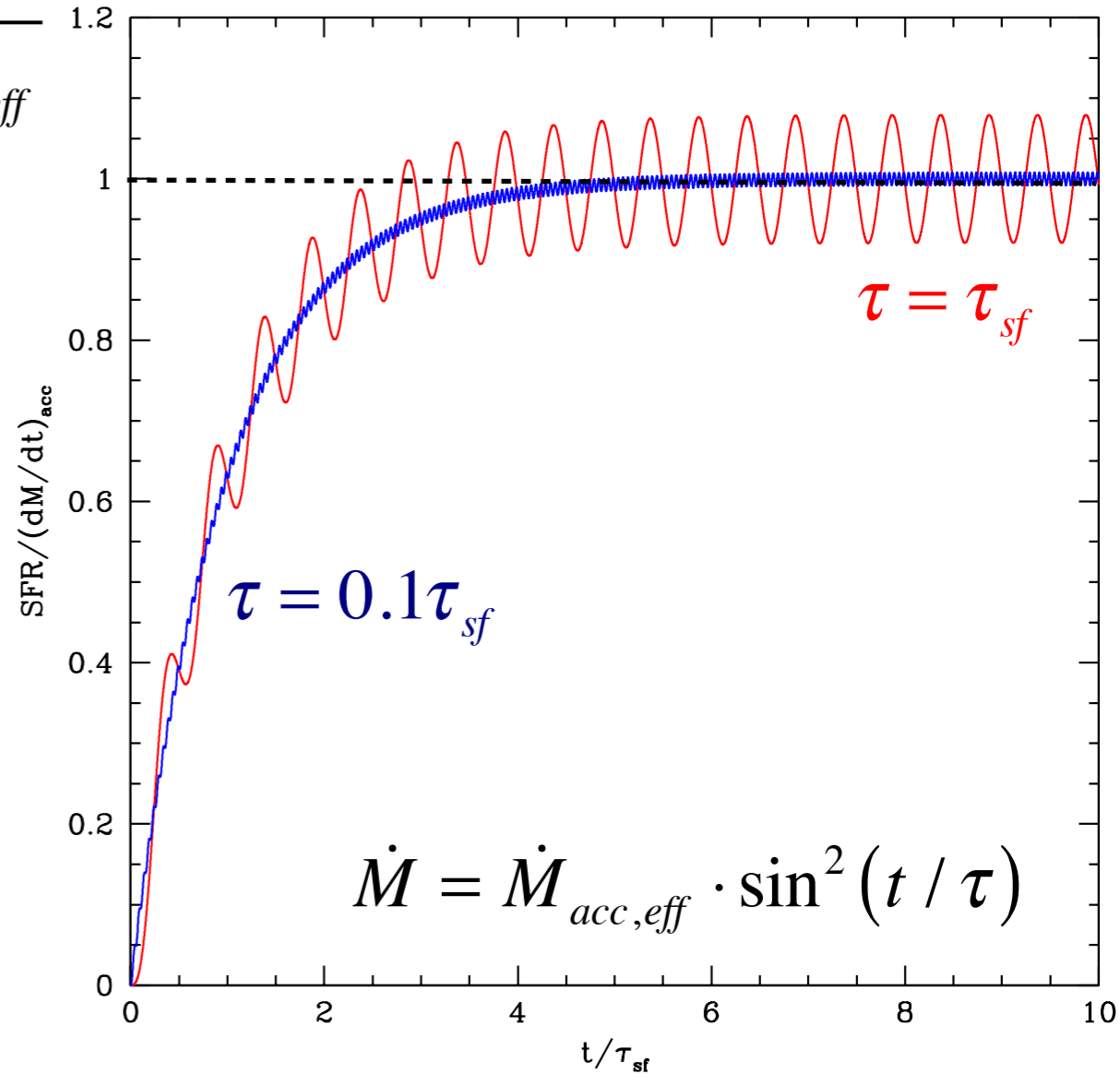
$$\frac{dM_g}{dt} = \left( \frac{dM_g}{dt} \right)_{acc} - \frac{M_g}{\tau_{sf}} (1 - R + \alpha_{wind})$$

$$\dot{M}_{acc,eff}$$

$$SFR = \frac{M_g}{\tau_{sf}} = \frac{1}{1 - R + \alpha_{wind}} \left( \frac{dM_g}{dt} \right)_{acc} \left( 1 - \exp\left( -\frac{t}{\tau_{sf}} \right) \right)$$

$$SFR = \dot{M}_{acc,eff}$$

$$\frac{SFR}{\dot{M}_{acc,eff}}$$



$$t / \tau_{sf}$$

- $\tau_{sf}$  does not determine SFR

# What determines the gas mass?

$$SFR = \dot{M}_{acc,eff}$$

- $\tau_{sf}$  determines the gas mass

$$M_g = \dot{M}_{acc,eff} \cdot \tau_{sf}$$

Why is  $\tau_{sf} \approx 10^9$  yrs?

# *Numerical simulations of the molecular web*

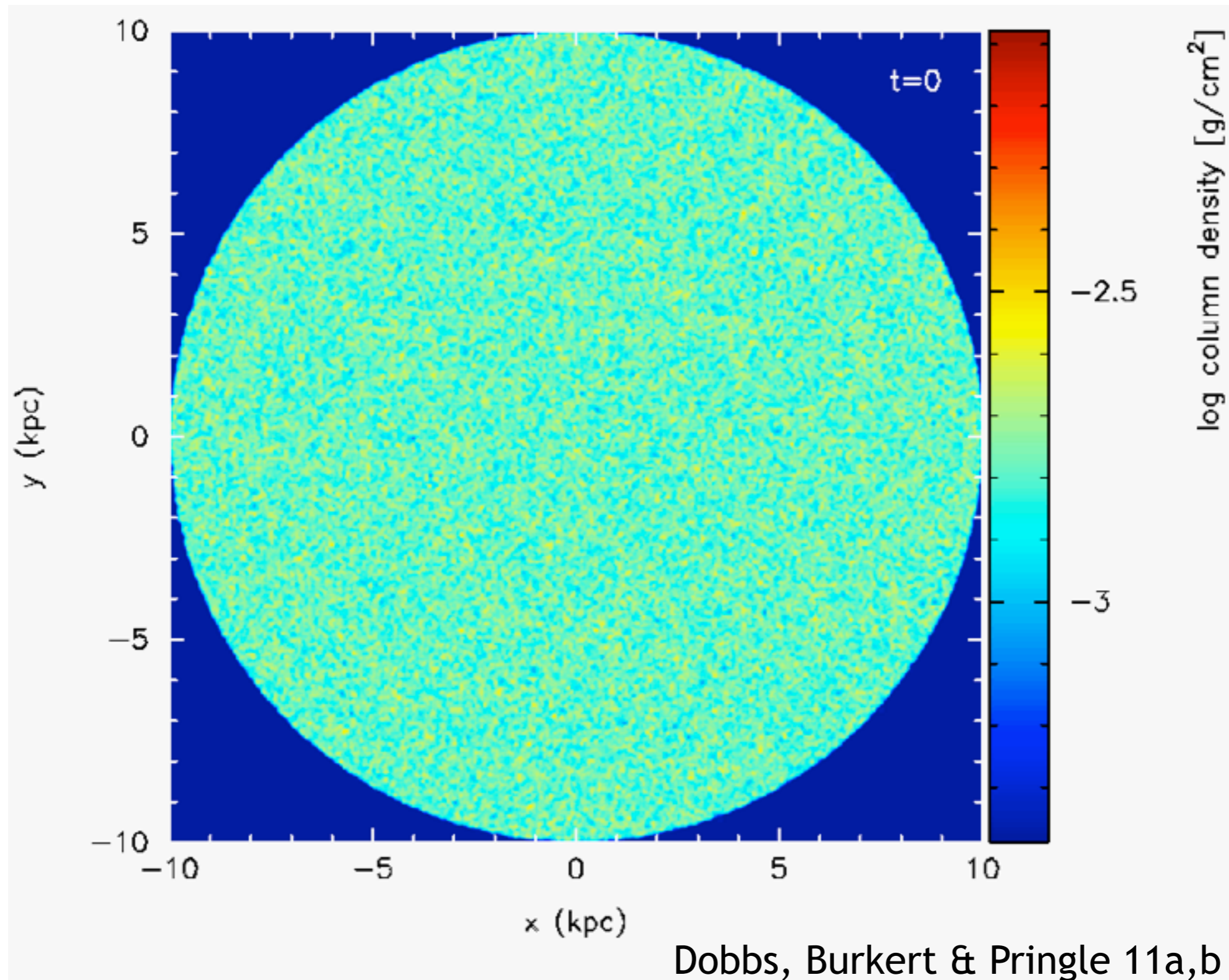
( Dobbs, Burkert & Pringle 11a,b)

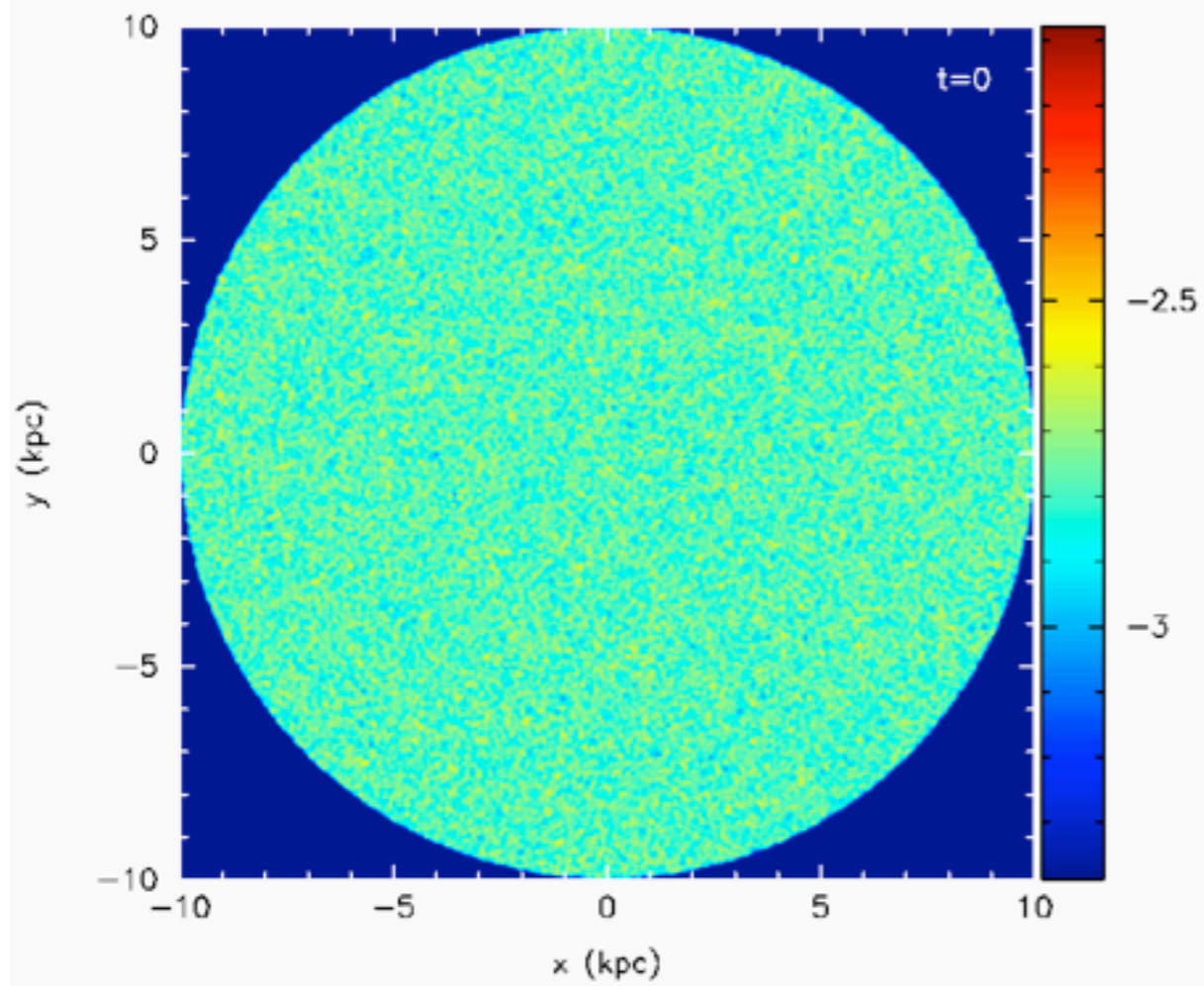
- 3d SPH simulations (Bate et al. 95)
- Fixed galactic gravitational potential (stellar disk + halo)
- Self-gravity of the gas component included
- Calculations with and without an additional 4 armed spiral potential
- Heating (supernovae + FUV background)
- Cooling: radiative + gas-grain energy transfer + recombination on grains
- Stars form when a local molecular region collapses and its density exceeds  $n_{crit} = 250 cm^{-3}$
- A fraction  $\epsilon$  of the gas is assumed to turn into stars that heat the environment with an energy (winds and SN) of

$$E_{SN} = \epsilon \frac{M_{dense}}{160 M_{\odot}} \cdot 10^{51} \text{ ergs}$$

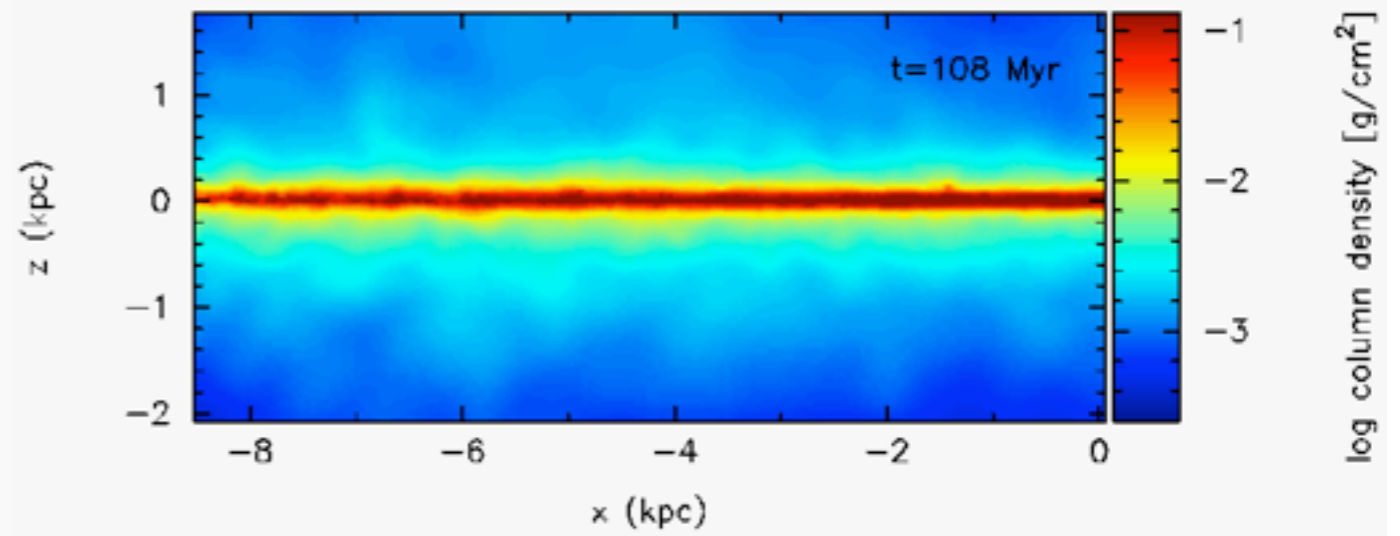
$$\epsilon \approx 2 - 5\%$$

# Calculation with 5 % efficiency

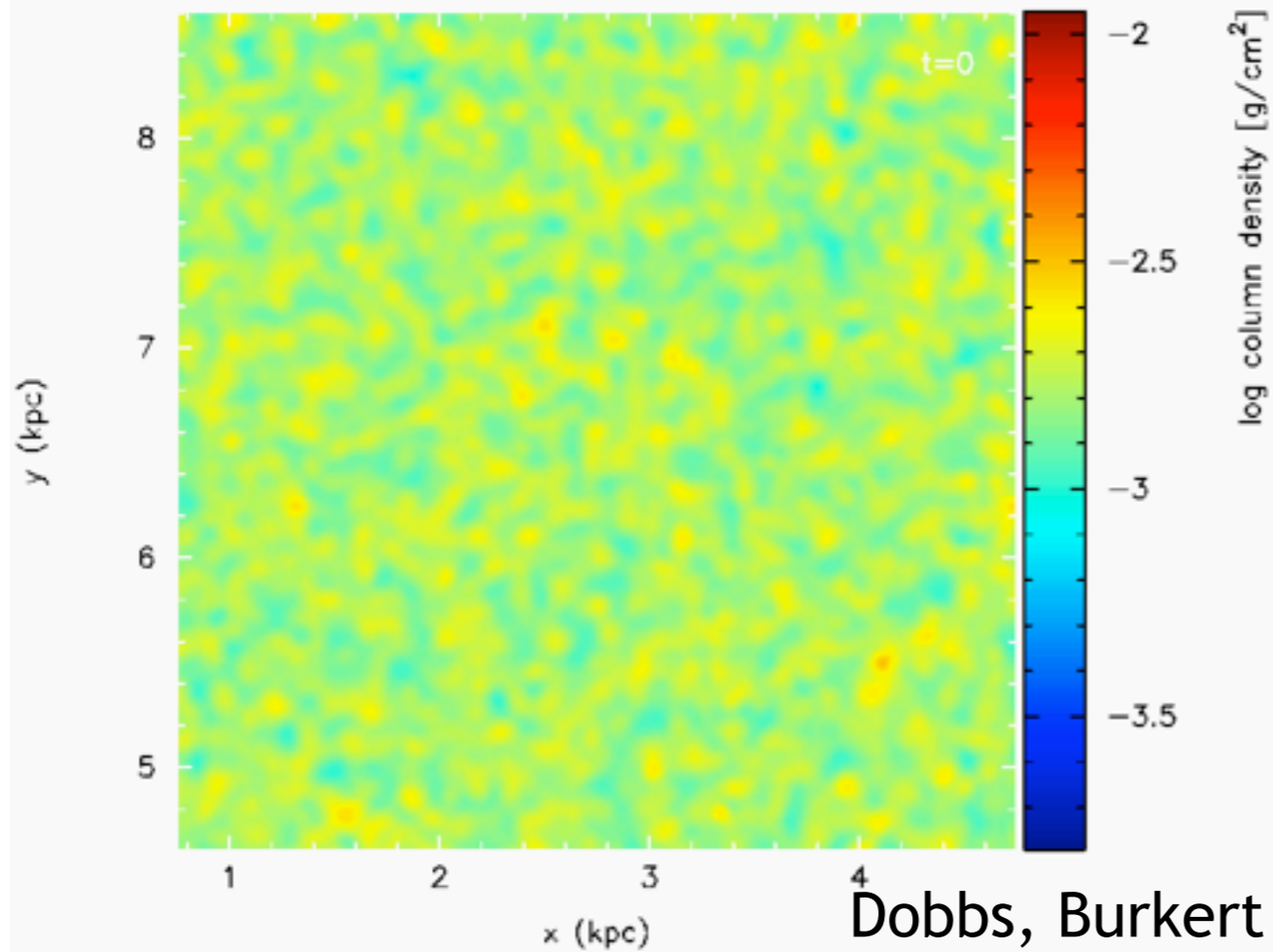




*Feedback puffs up disk*



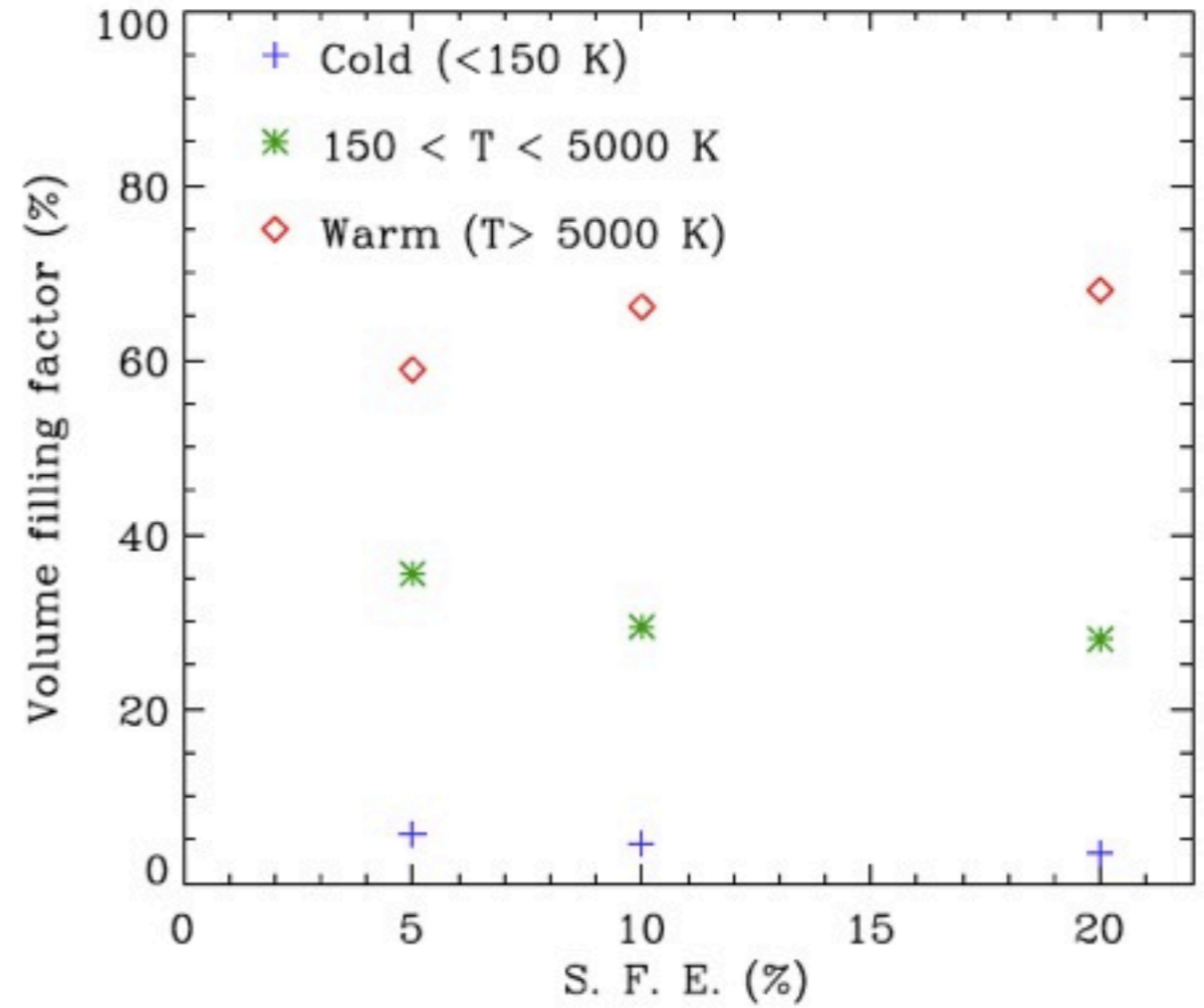
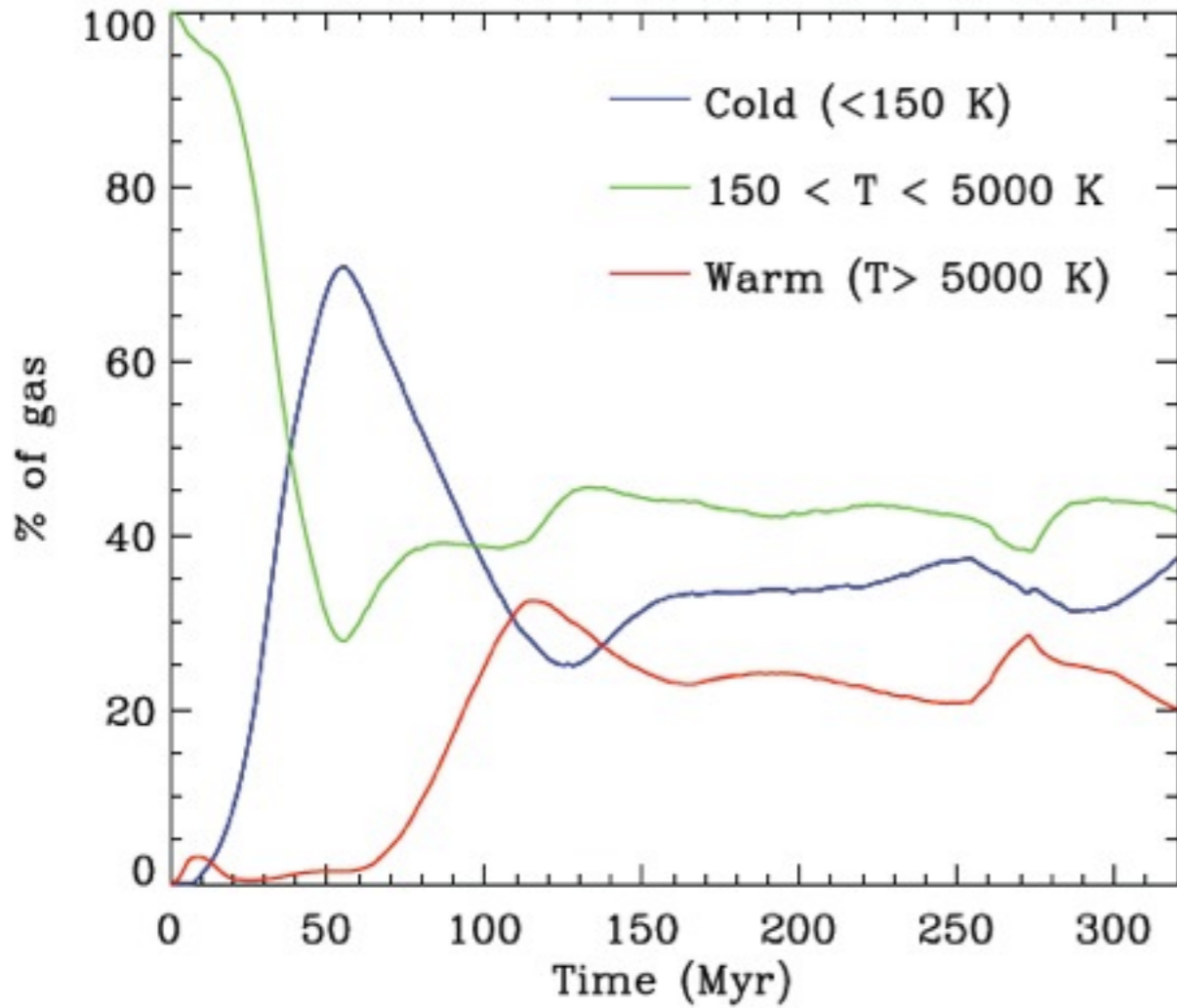
*Filamentary interarm features (spurs)*



Dobbs, Burkert & Pringle 11a,b

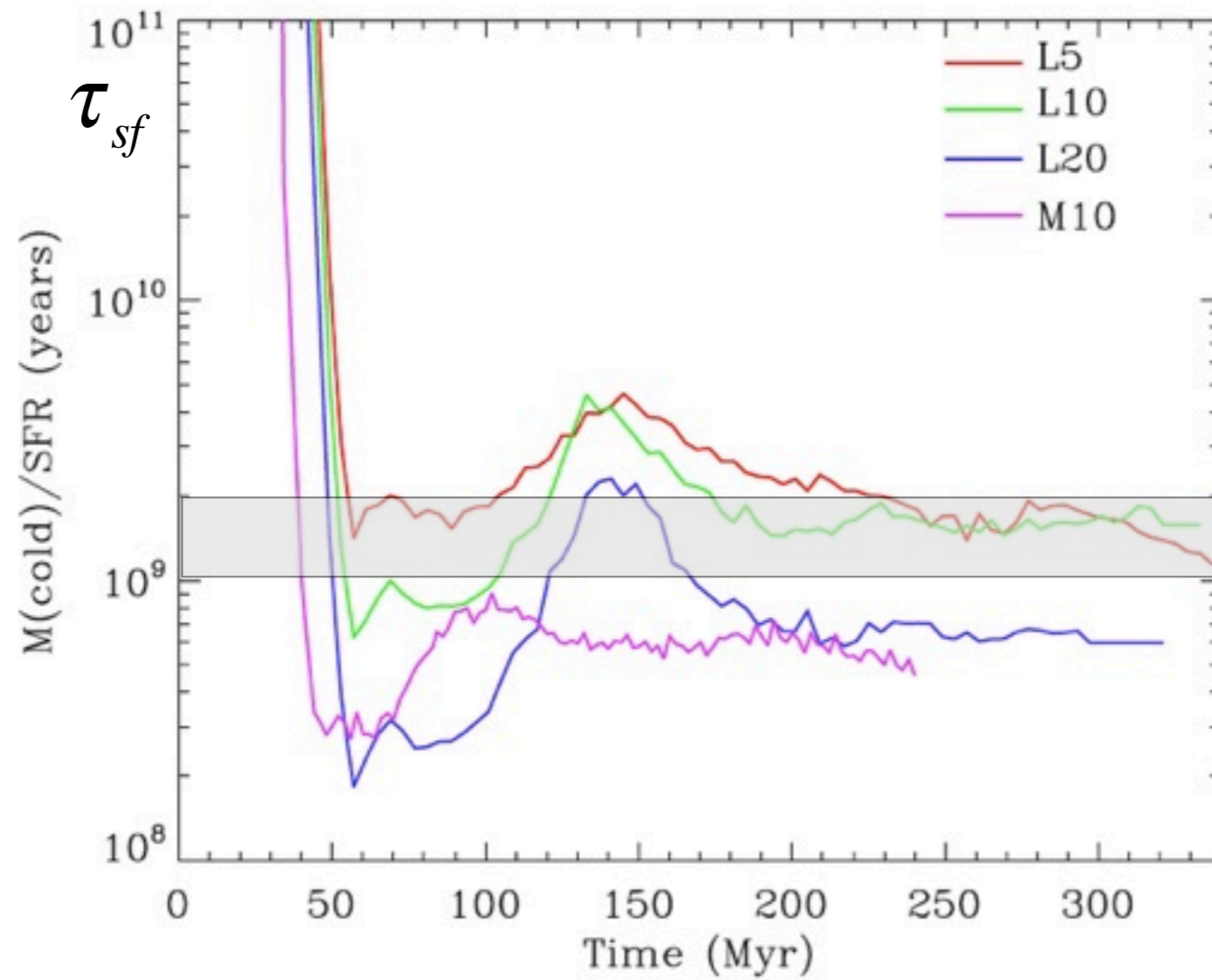


# Gas mass fraction and volume filling factor: 5% efficiency



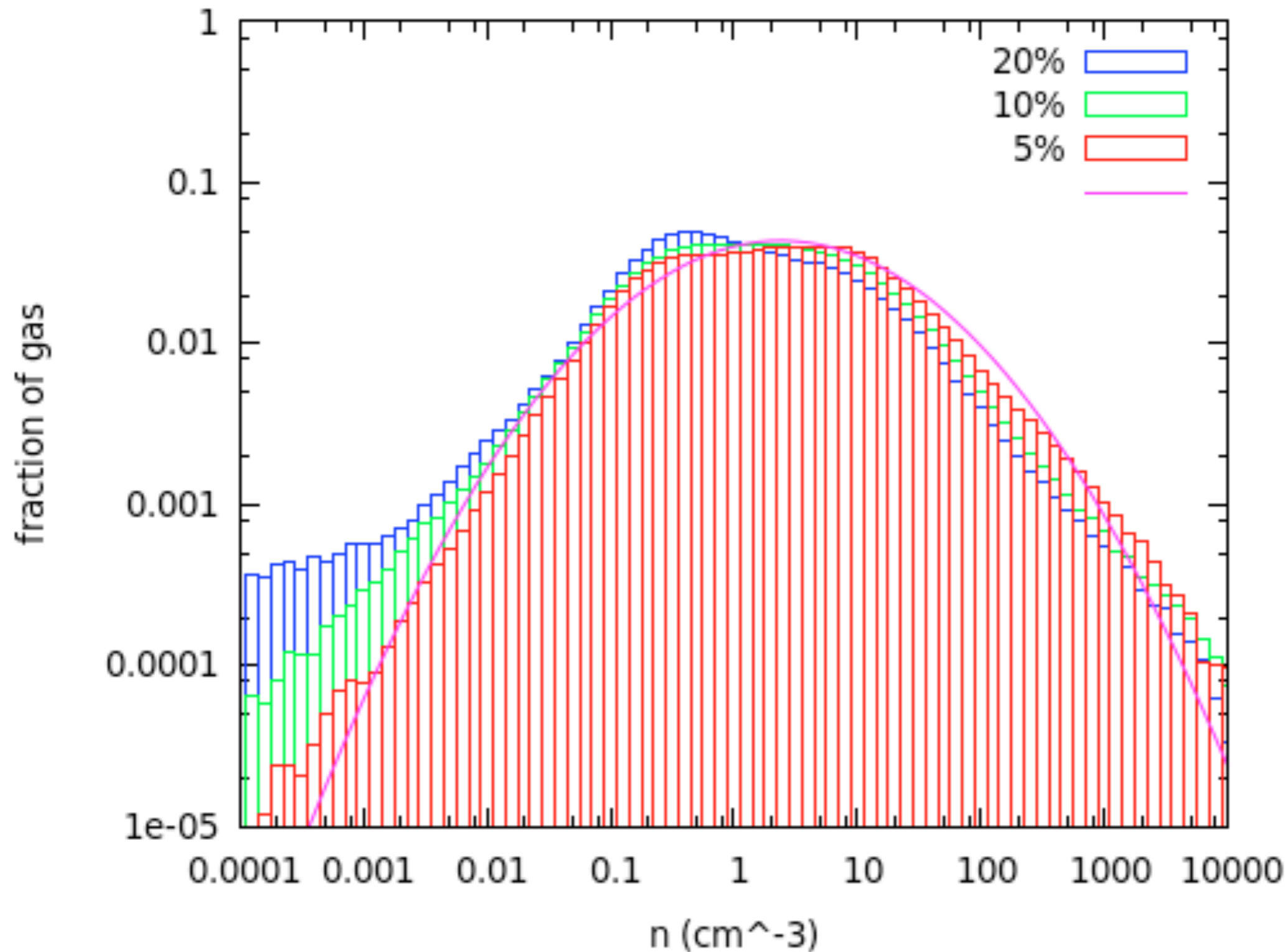
# Star formation timescale

$$SFR = \frac{M_{H_2}}{\tau_{sf}} \text{ with } \tau_{sf} \approx 1 - 2 \cdot 10^9 \text{ yrs}$$

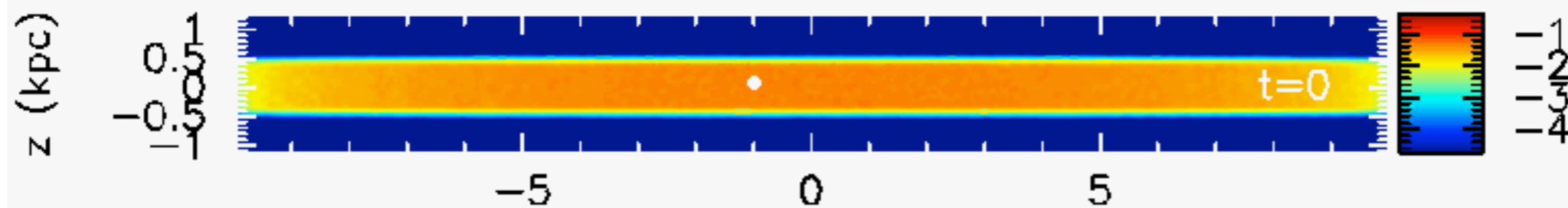
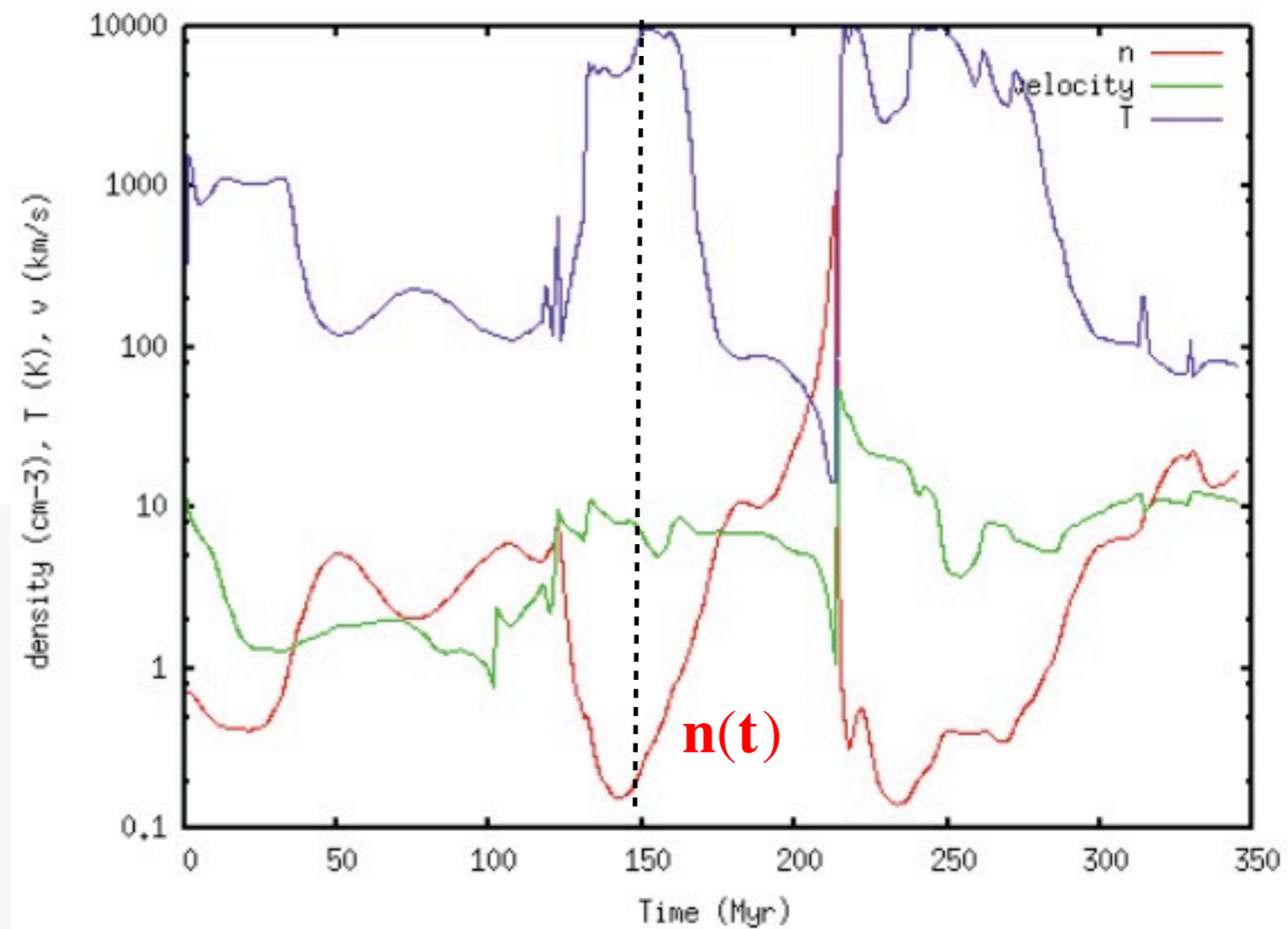
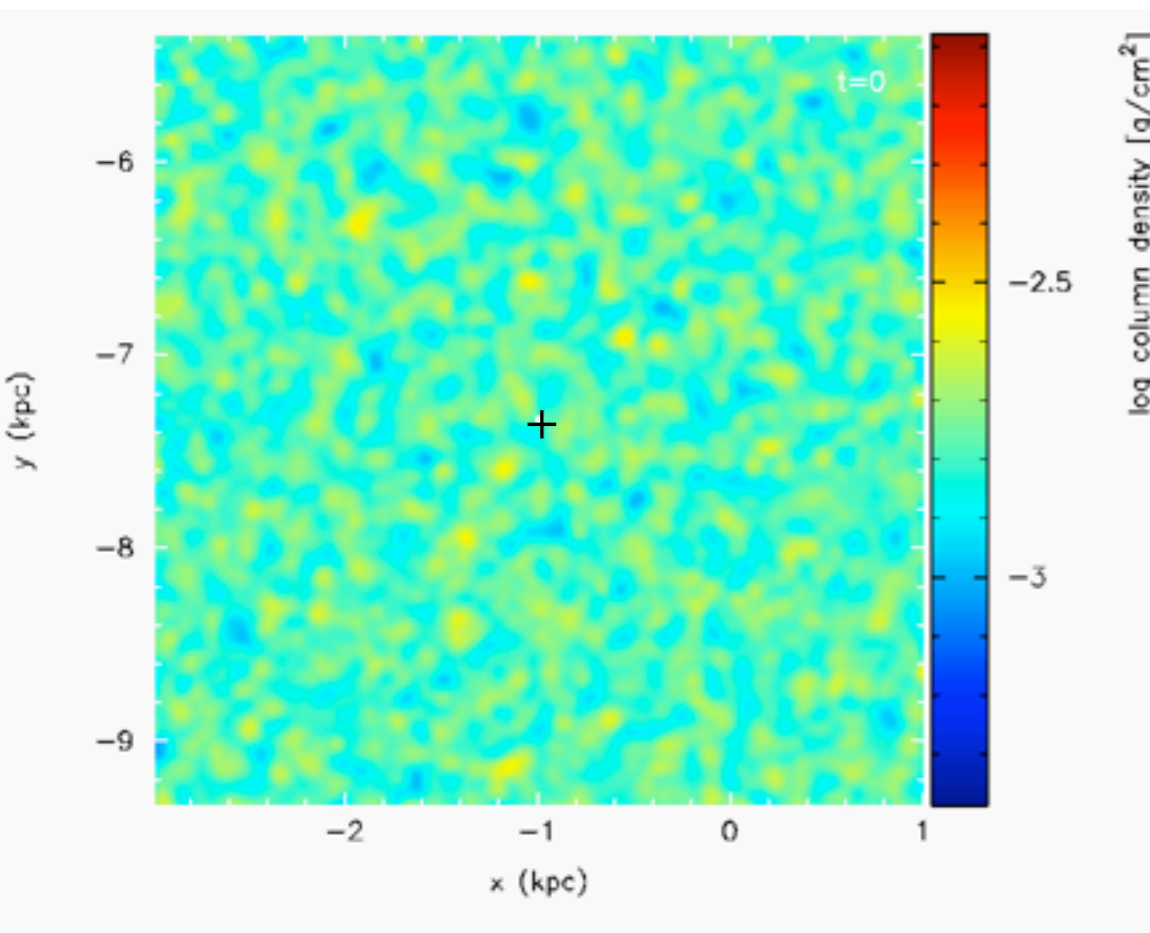


# Density Probability Distribution Function

(Elmegreen 02; Krumholz & McKee 05; Wada & Norman 07)



# Gravitational instabilities and star formation timescale



## Growth rate of gravitational instabilities:

$$\tau_{\text{Toomre}} = \frac{\sigma}{\pi G \Sigma} = \kappa^{-1} = \left(\sqrt{2\Omega}\right)^{-1} \rightarrow \tau_{\text{Toomre}} = 0.1 \cdot \tau_{\text{orb}} \approx 2 \cdot 10^7 \text{ yrs}$$

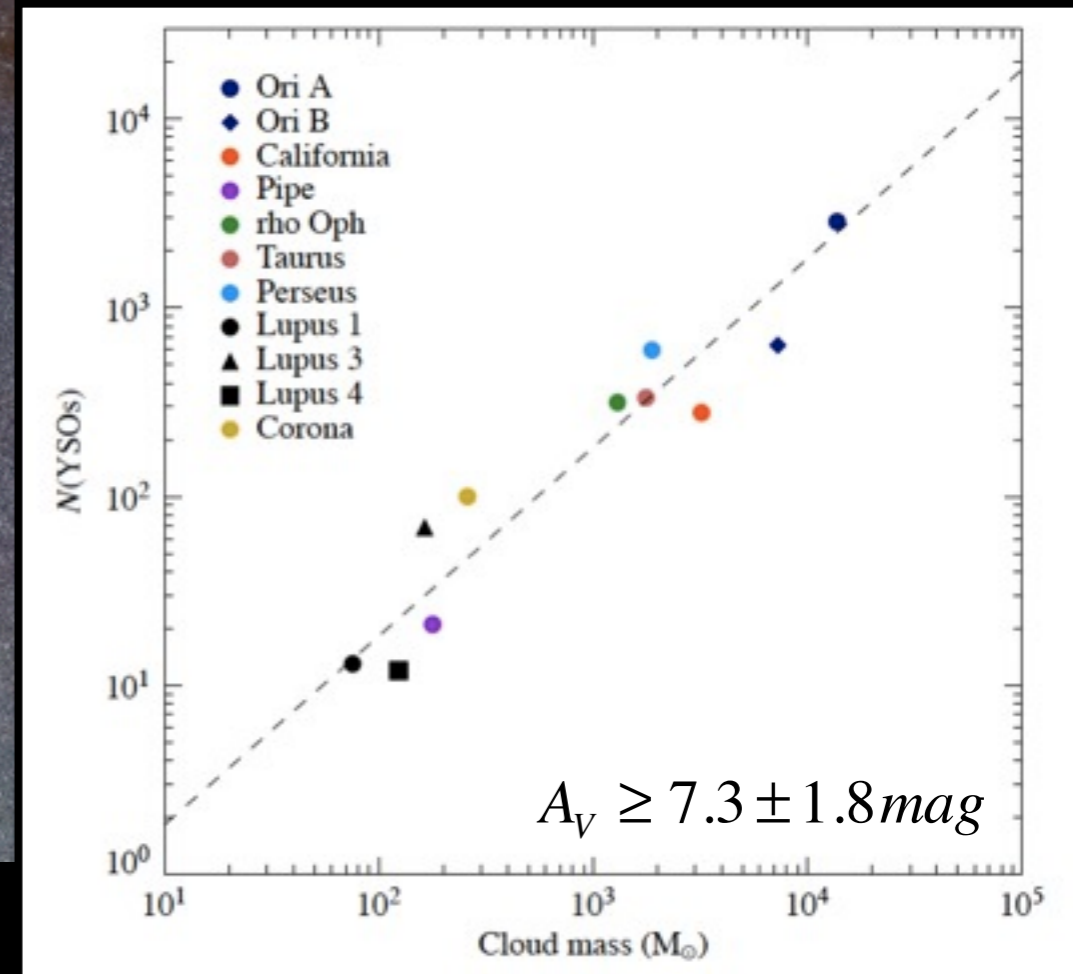
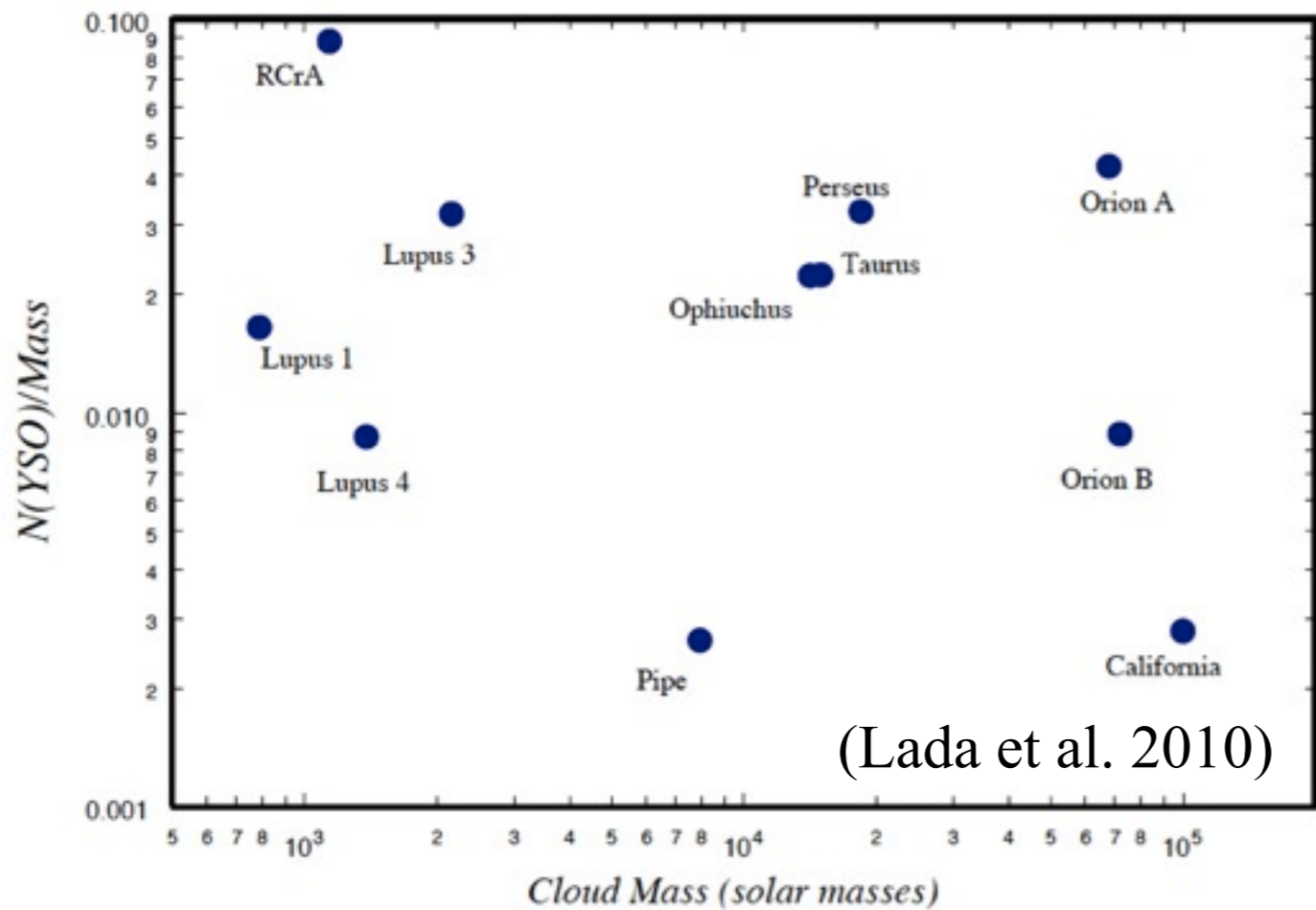
$$Q = 1$$

$$\tau_{\text{orb}} \sim \frac{R_{\text{vir}}}{V_{\text{vir}}} \sim H^{-1}$$

$$\tau_{\text{SF}} \approx 10^9 \text{ yrs} \approx 50 \cdot \tau_{\text{Toomre}} \approx \tau_{\text{Toomre}} / \epsilon$$

$$SFR = 0.02 \frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}}$$

$$N(\text{YSOs})_{\text{Oph}} = 15 \times N(\text{YSOs})_{\text{Pipe}}$$



$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{ff} \approx 3.5 \cdot 10^5 \text{ yrs}$$

$$SFR \approx 0.02 \frac{M_{dense}}{\tau_{ff}}$$

Large scales:

$$SFR \approx \frac{M_{H_2}}{10^9 \text{ yrs}} \approx 0.02 \frac{M_{H_2}}{\tau_{dyn}}$$

$$N(\text{YSOs})_{\text{Oph}} = 15 \times N(\text{YSOs})_{\text{Pipe}}$$

$$\frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}} = \frac{M_{\text{dense}}}{\tau_{\text{ff}}}$$

664  $M_{\odot}$   
14165  $M_{\odot}$

316 YSOs

Lada & Alves

$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 3.5 \cdot 10^5 \text{ yrs} \rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}}}{\tau_{\text{ff}}}$$

Large scales:

$$\text{SFR} \approx \frac{M_{\text{H}_2}}{10^9 \text{ yrs}} \approx 0.02 \frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}}$$

$$N(\text{YSOs})_{\text{Oph}} = 15 \times N(\text{YSOs})_{\text{Pipe}}$$

$$\frac{M_{H_2}}{\tau_{\text{Toomre}}} = \frac{M_{\text{dense}}}{\tau_{\text{ff}}}$$

$$\frac{dM_{\text{dense}}}{dt} = \frac{M_{\text{diff}}}{\tau_{\text{Toomre}}} - \frac{M_{\text{dense}}}{\tau_{\text{ff}}} \approx \frac{M_{H_2}}{\tau_{\text{Toomre}}} - \frac{M_{\text{dense}}}{\tau_{\text{Toomre}}} - \frac{M_{\text{dense}}}{\tau_{\text{ff}}} = 0$$

$M_{\text{dense}} \ll M_{H_2}$

664  $M_{\odot}$   
14165  $M_{\odot}$

316 YSOs

Lada & Alves

Large scales:

$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 3.5 \cdot 10^5 \text{ yrs} \rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}}}{\tau_{\text{ff}}}$$

$$\text{SFR} \approx \frac{M_{H_2}}{10^9 \text{ yrs}} \approx 0.02 \frac{M_{H_2}}{\tau_{\text{Toomre}}}$$



$$N(\text{YSOs})_{\text{Oph}} = 15 \times N(\text{YSOs})_{\text{Pipe}}$$

$$\frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}} = \frac{M_{\text{dense}}}{\tau_{\text{ff}}} \rightarrow M_{\text{dense}} = \frac{\tau_{\text{ff}}}{\tau_{\text{Toomre}}} M_{\text{H}_2} \approx 0.02 M_{\text{H}_2}$$

$$M_{\text{diff}} \approx 0.98 M_{\text{H}_2}$$

664  $M_{\odot}$   
14165  $M_{\odot}$

316 YSOs

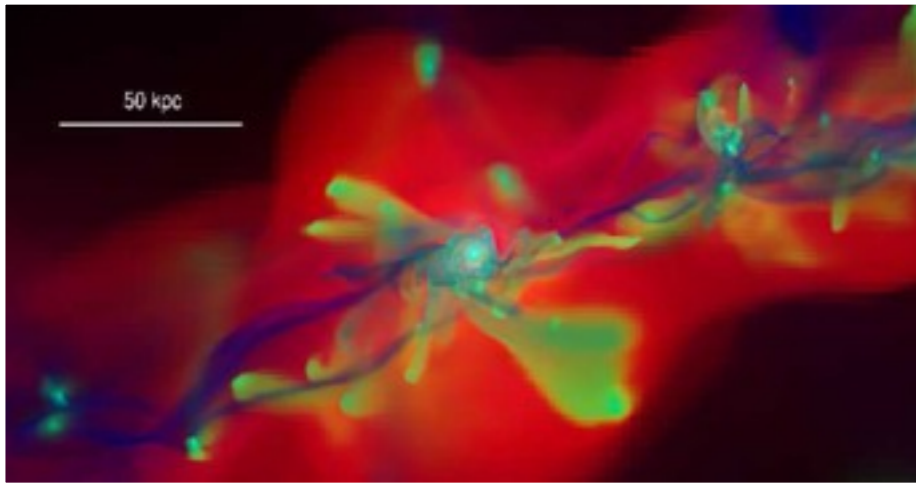
Lada & Alves

$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 3.5 \cdot 10^5 \text{ yrs} \rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}}}{\tau_{\text{ff}}}$$

Large scales:

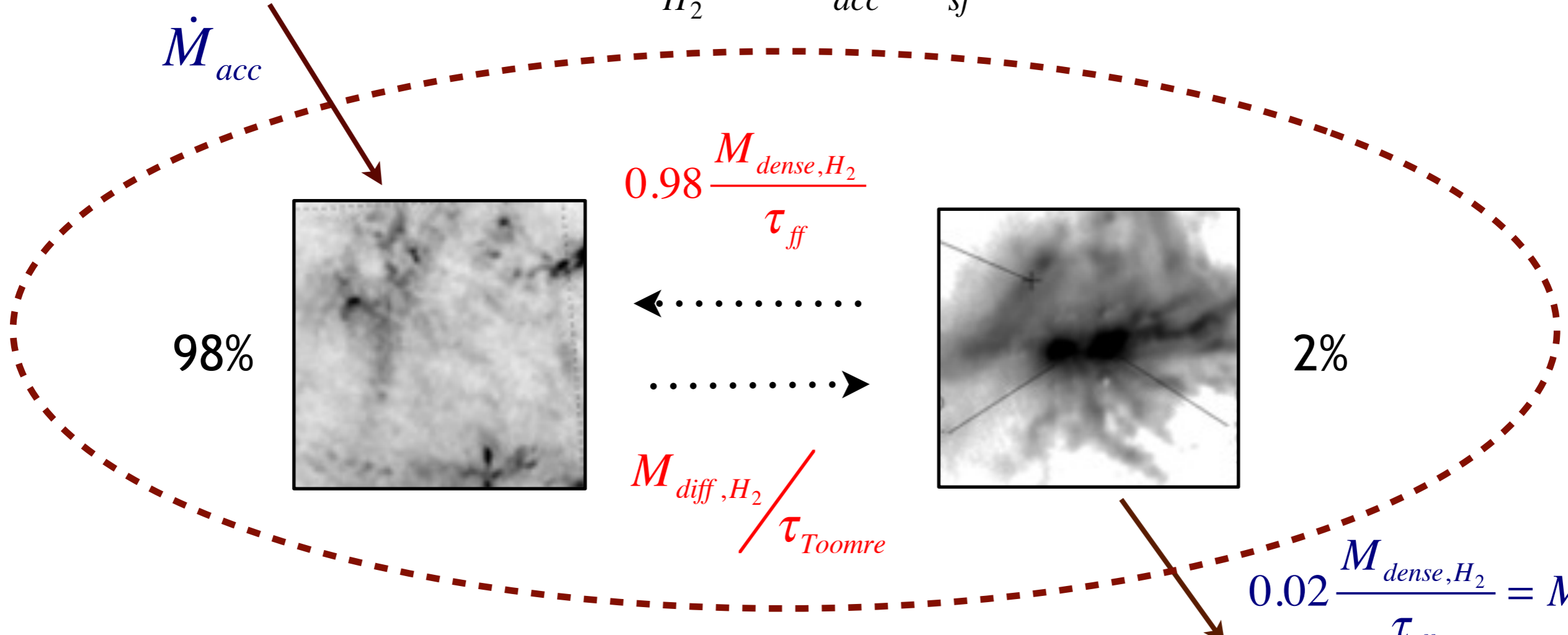
$$\text{SFR} \approx \frac{M_{\text{H}_2}}{10^9 \text{ yrs}} \approx 0.02 \frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}}$$

# Self-regulated star formation



$$M_{H_2} = \dot{M}_{acc} \cdot \tau_{sf}$$

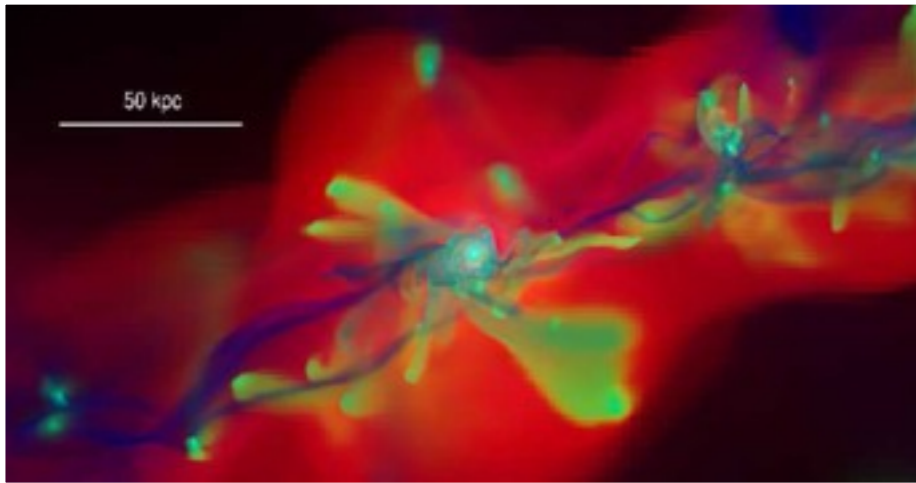
$\dot{M}_{acc}$



$$\tau_{sf} = 50 \cdot \tau_{Toomre} = \frac{1}{\epsilon \cdot \kappa}$$

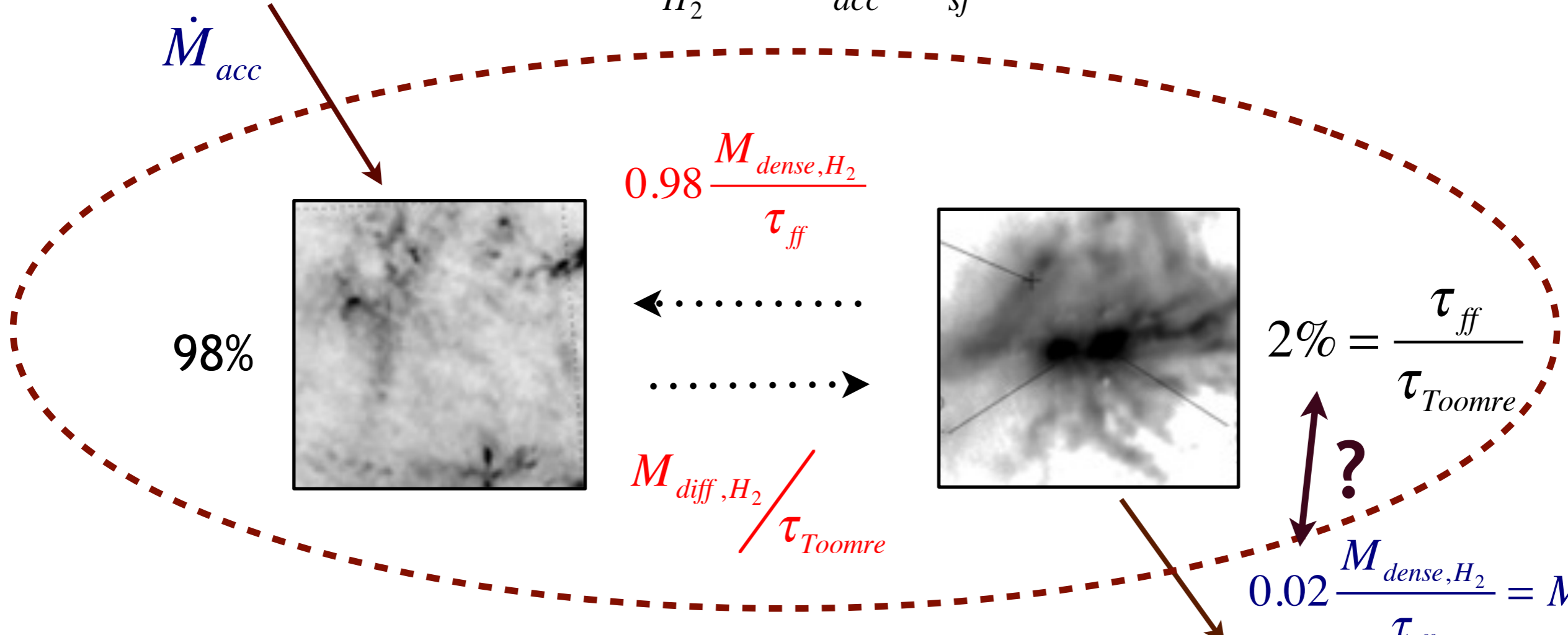


# Self-regulated star formation



$$M_{H_2} = \dot{M}_{acc} \cdot \tau_{sf}$$

$\dot{M}_{acc}$



$$0.98 \frac{M_{dense, H_2}}{\tau_{ff}}$$

98%

$$2\% = \frac{\tau_{ff}}{\tau_{Toomre}}$$

$$\frac{M_{diff, H_2}}{\tau_{Toomre}}$$

$$0.02 \frac{M_{dense, H_2}}{\tau_{ff}} = \dot{M}_{acc}$$

$$\tau_{sf} = 50 \cdot \tau_{Toomre} = \frac{1}{\epsilon \cdot \kappa}$$

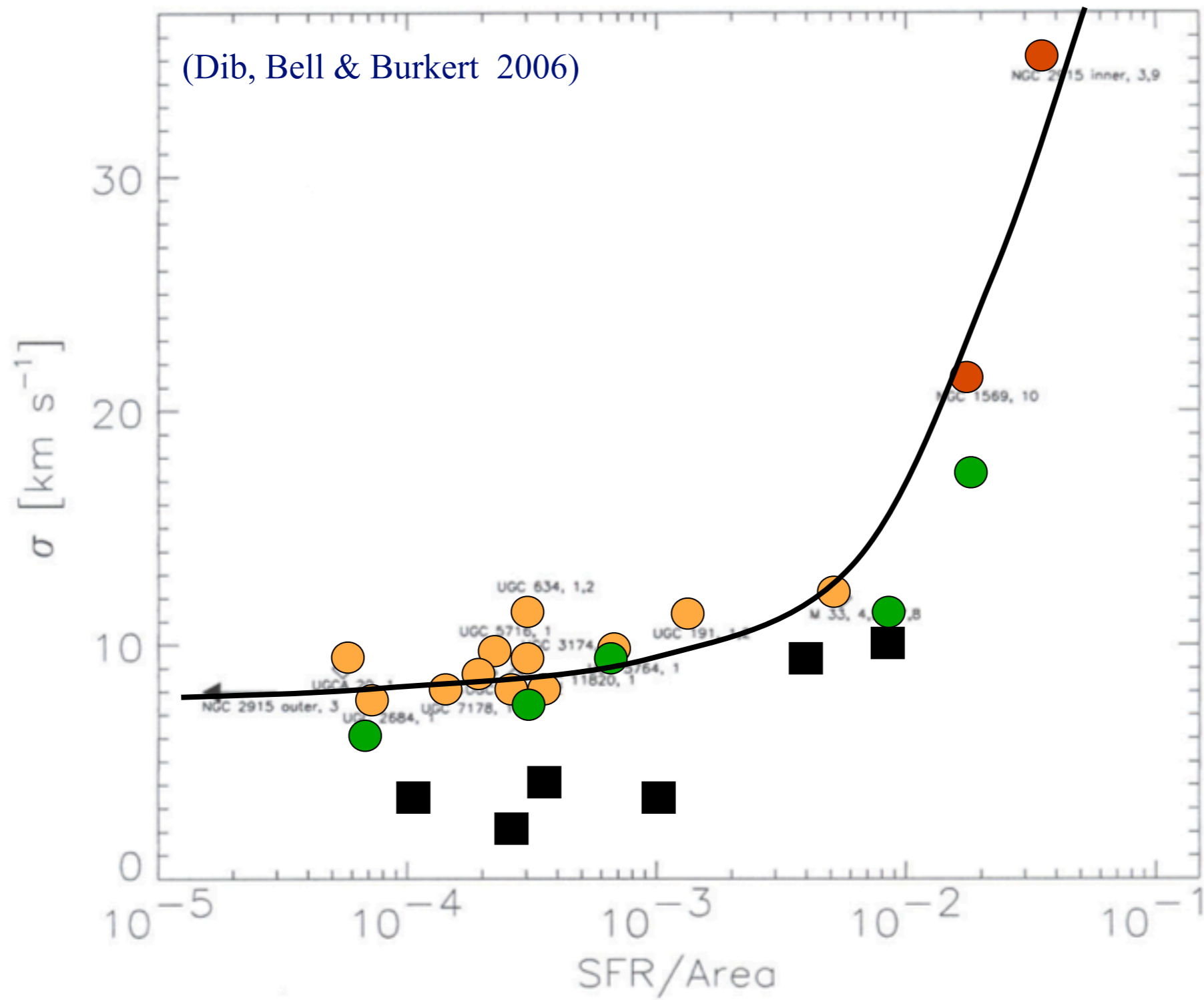


$$SFR \approx \varepsilon \frac{M_{H_2}}{\tau_{Toomre}} \approx \frac{M_{H_2}}{10^9 \text{ yrs}}$$

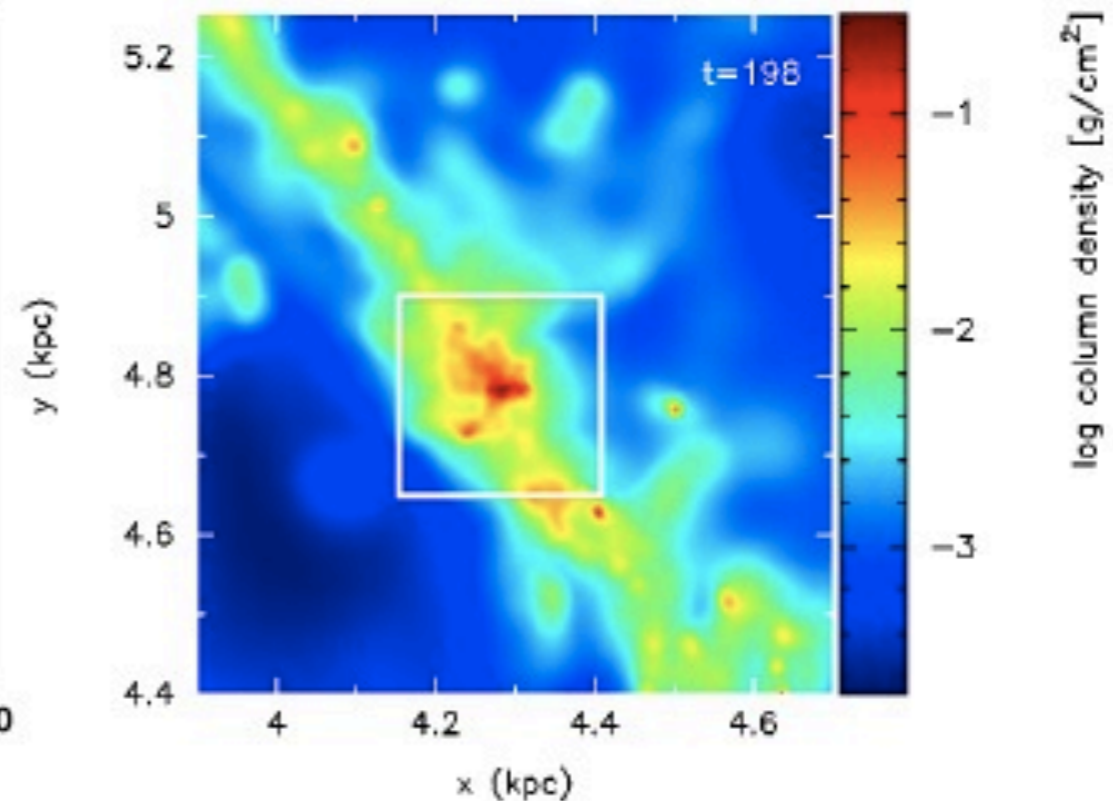
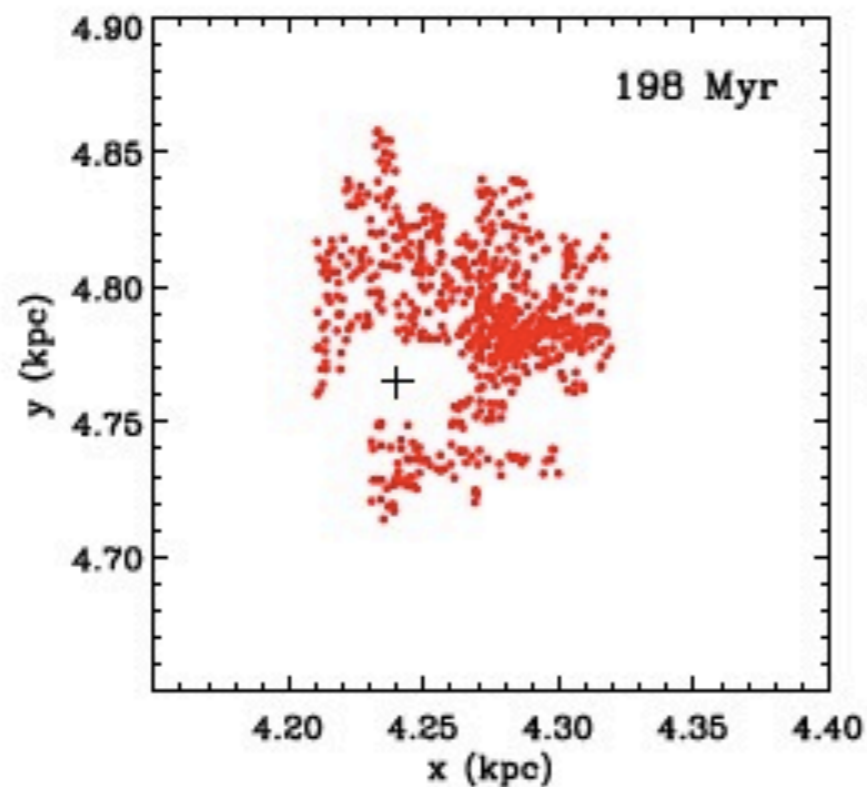
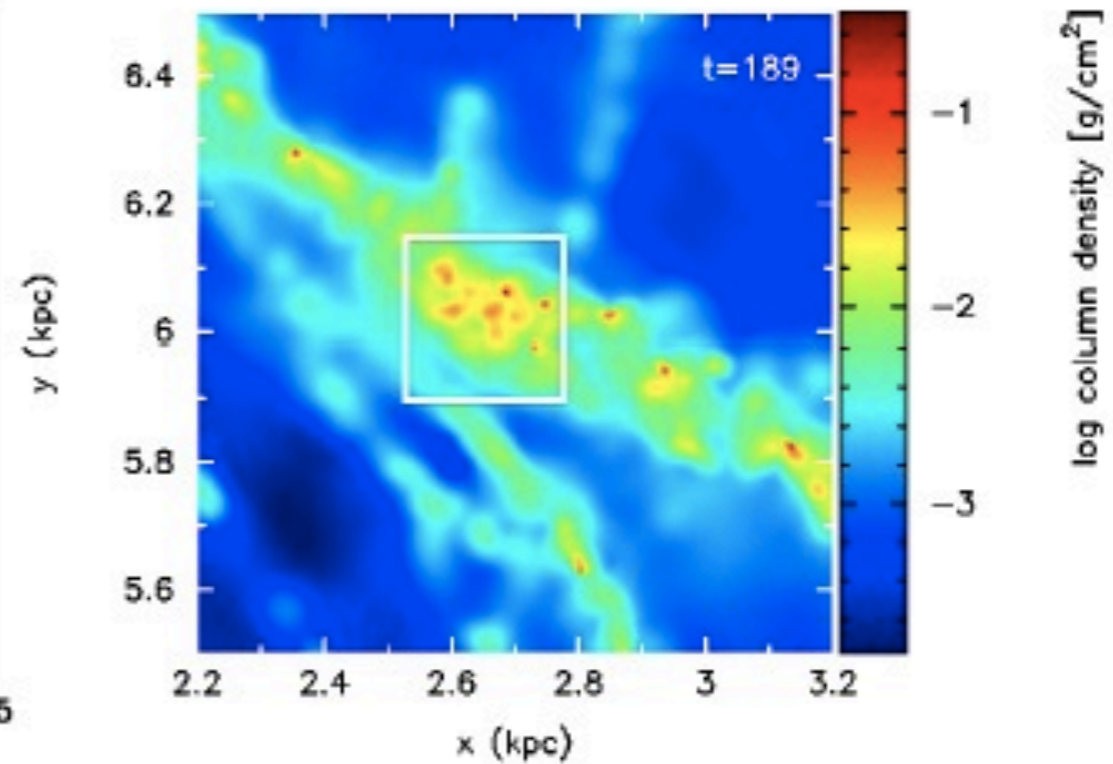
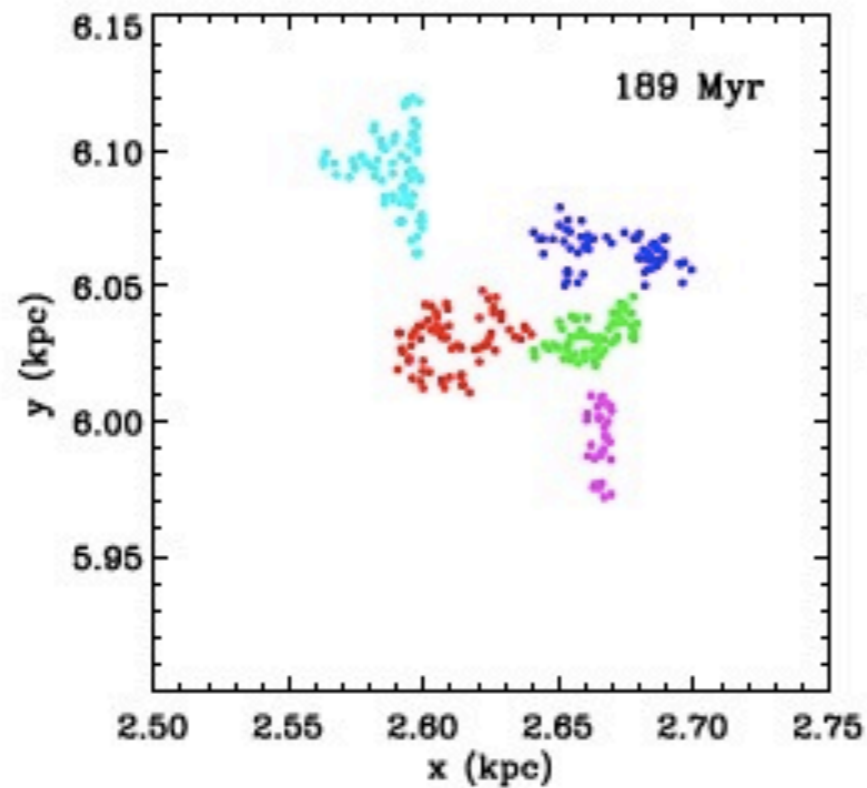
$$\varepsilon \approx 0.02$$

***What determines the star formation efficiency?***

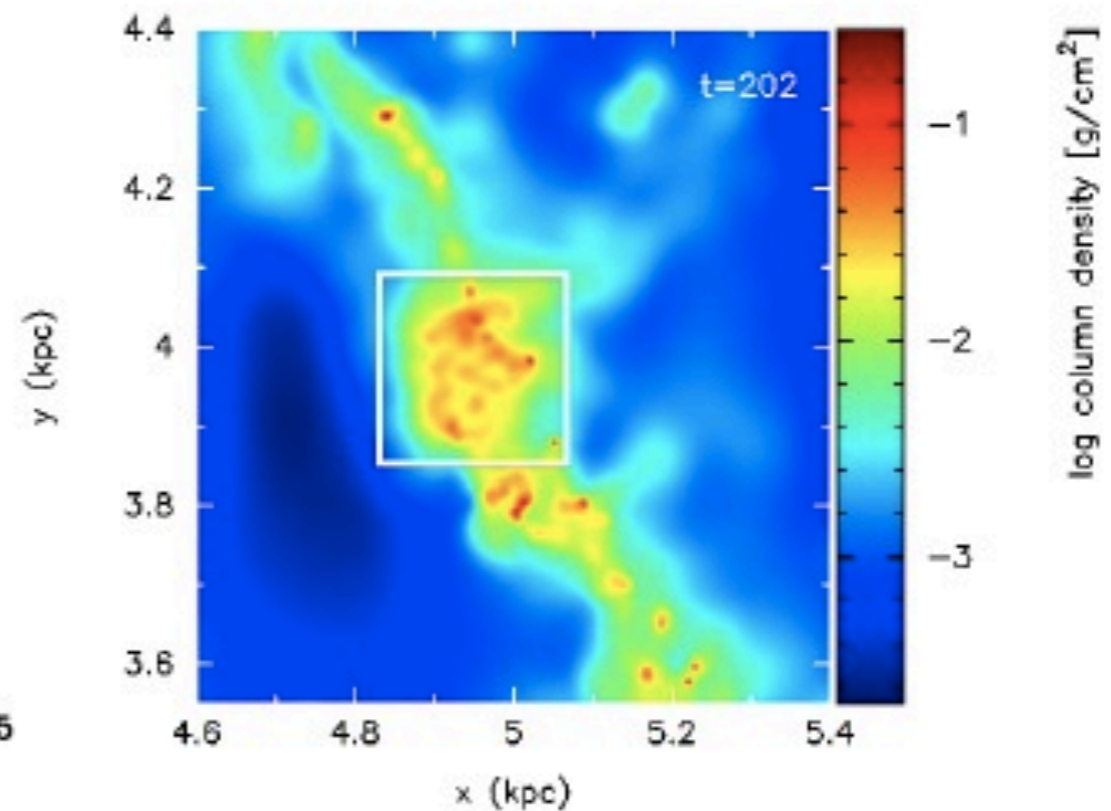
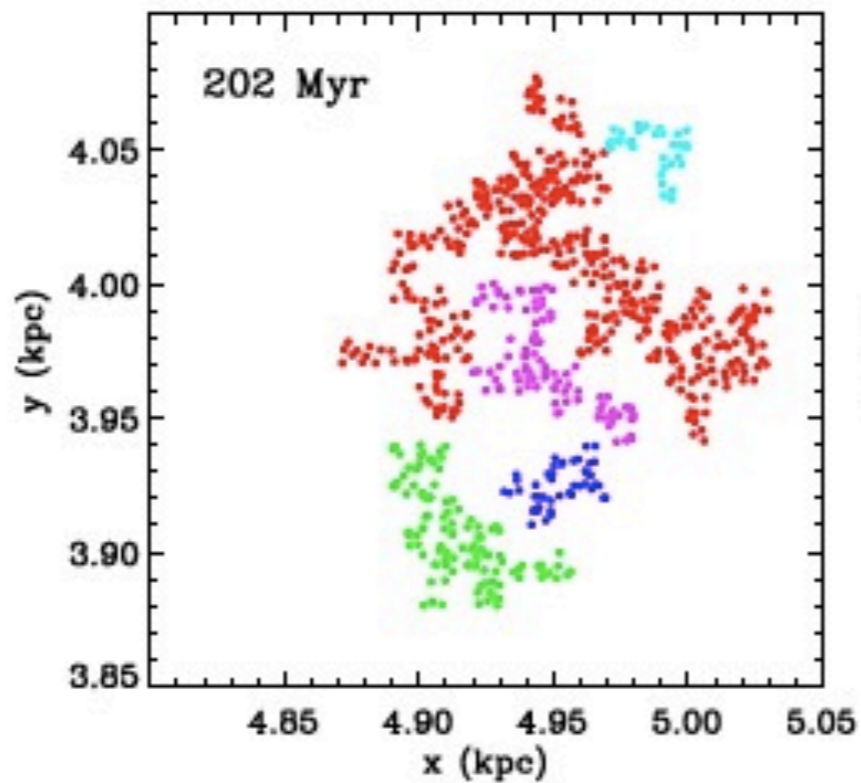
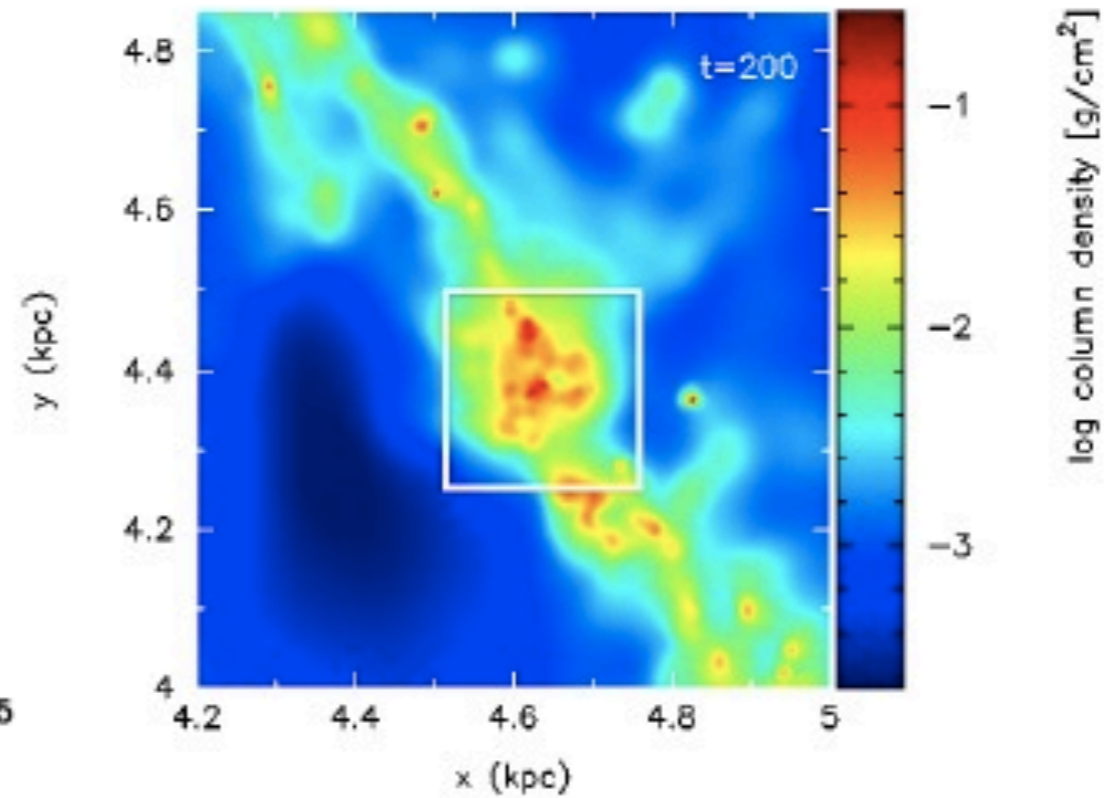
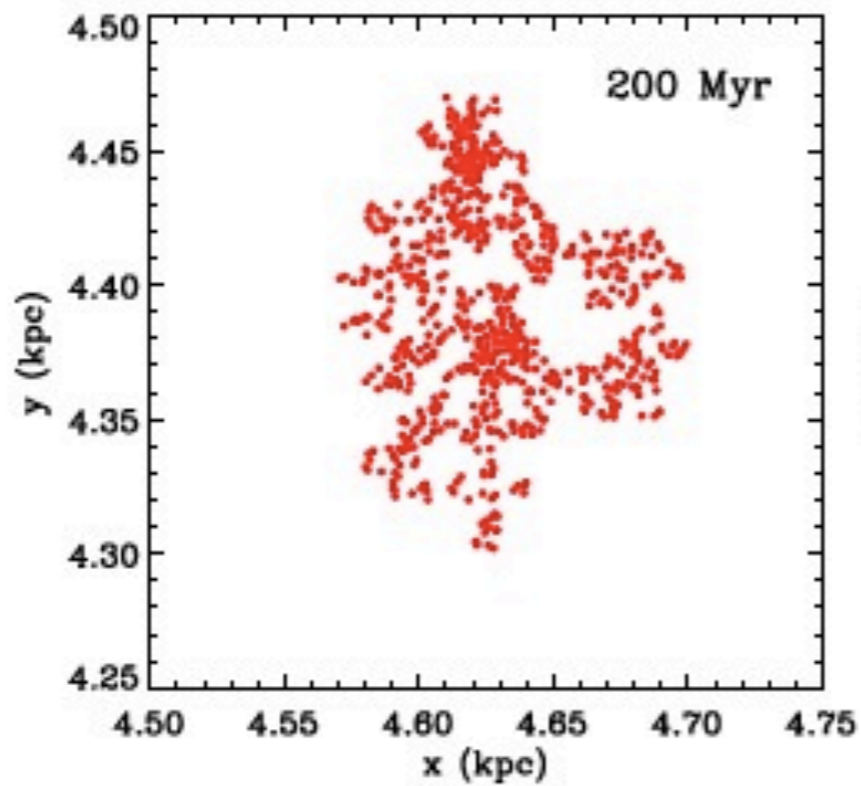
# Turbulence in the ISM



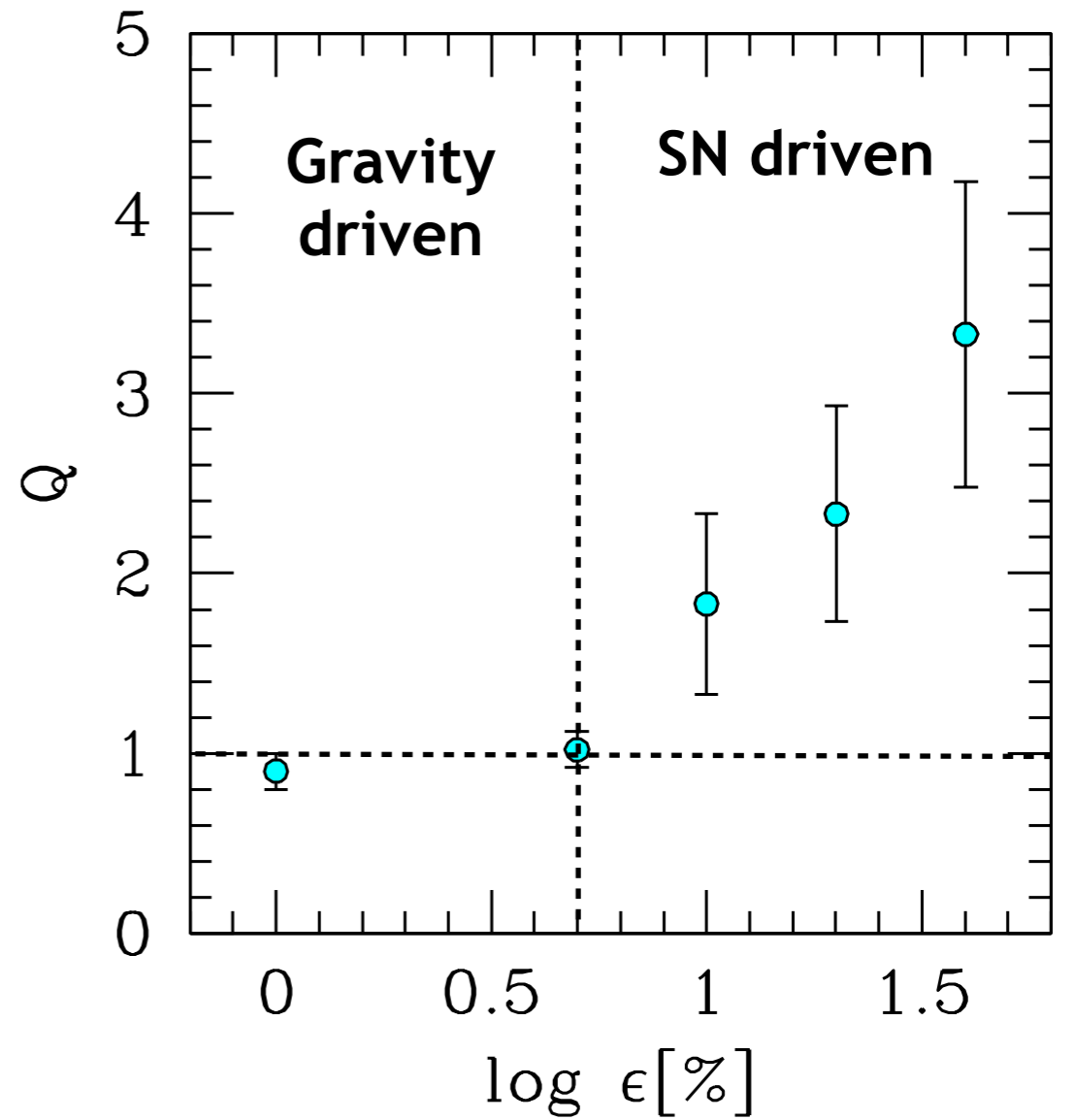
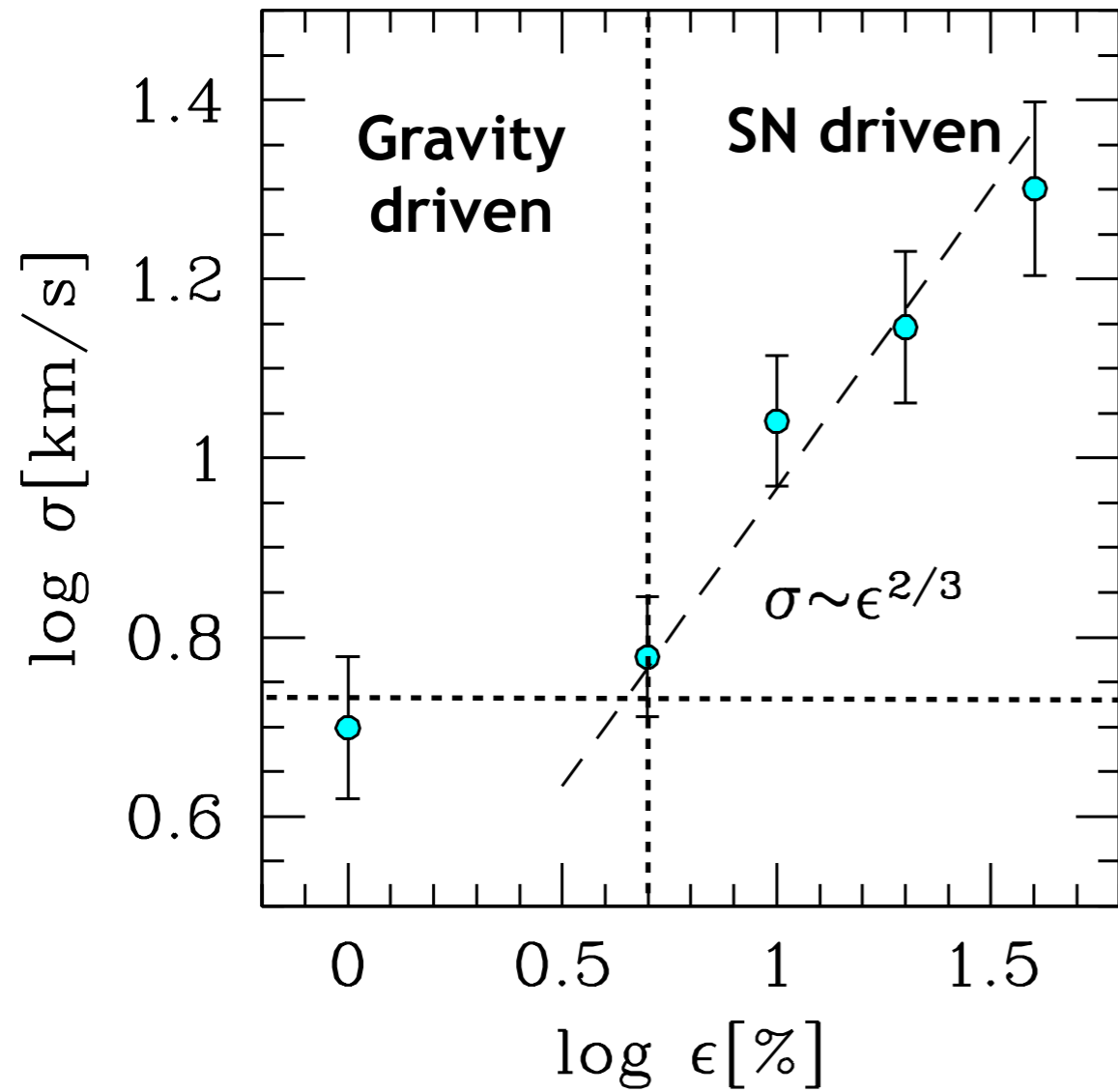
# 1. Collisions by local gravitational instability and irregular gas motions generate massive clouds and drive internal turbulence



## 2. *Stellar feedback disperses clouds and drives irregular gas motions in the molecular web.*



# Gas velocity dispersion



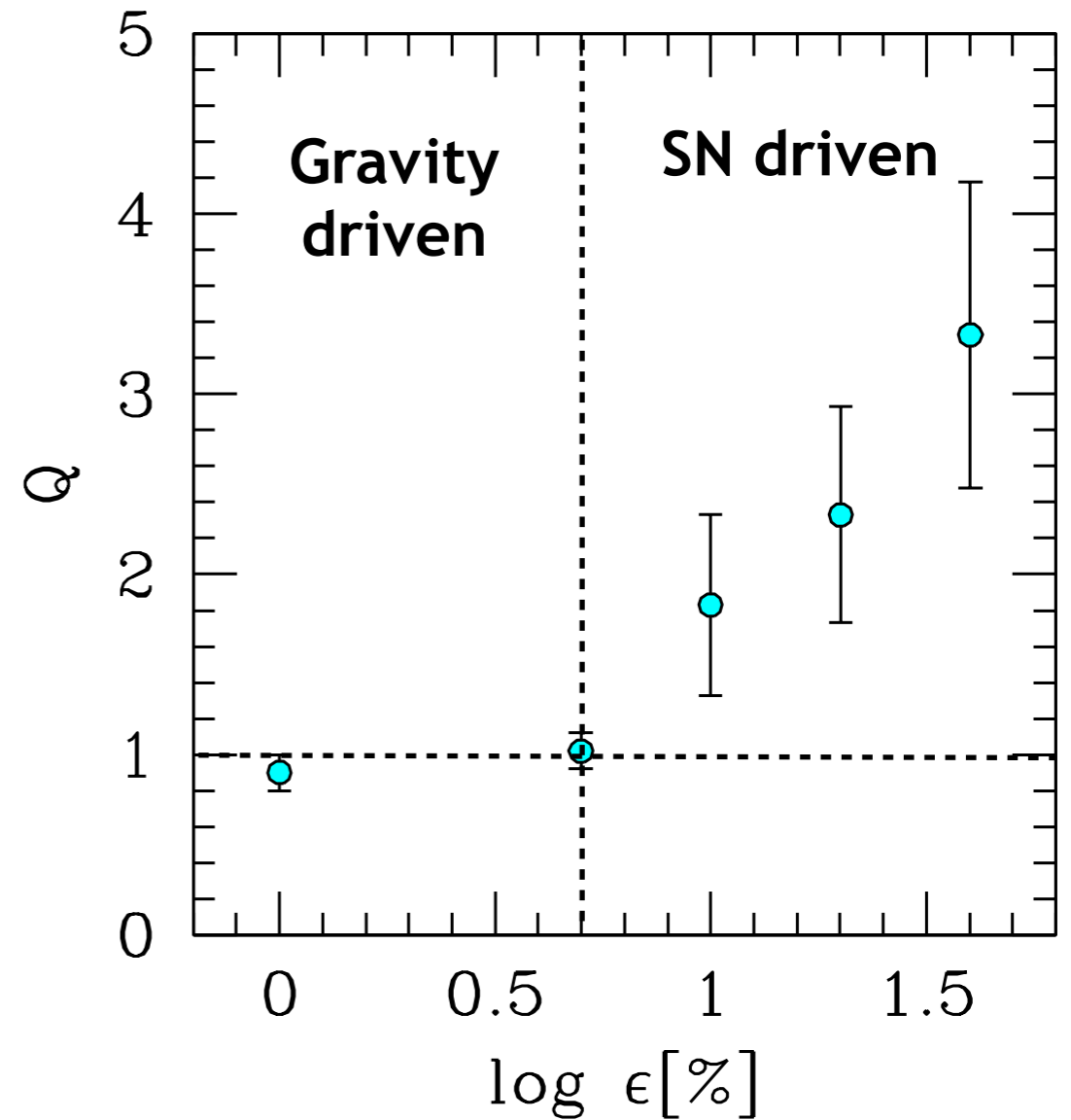
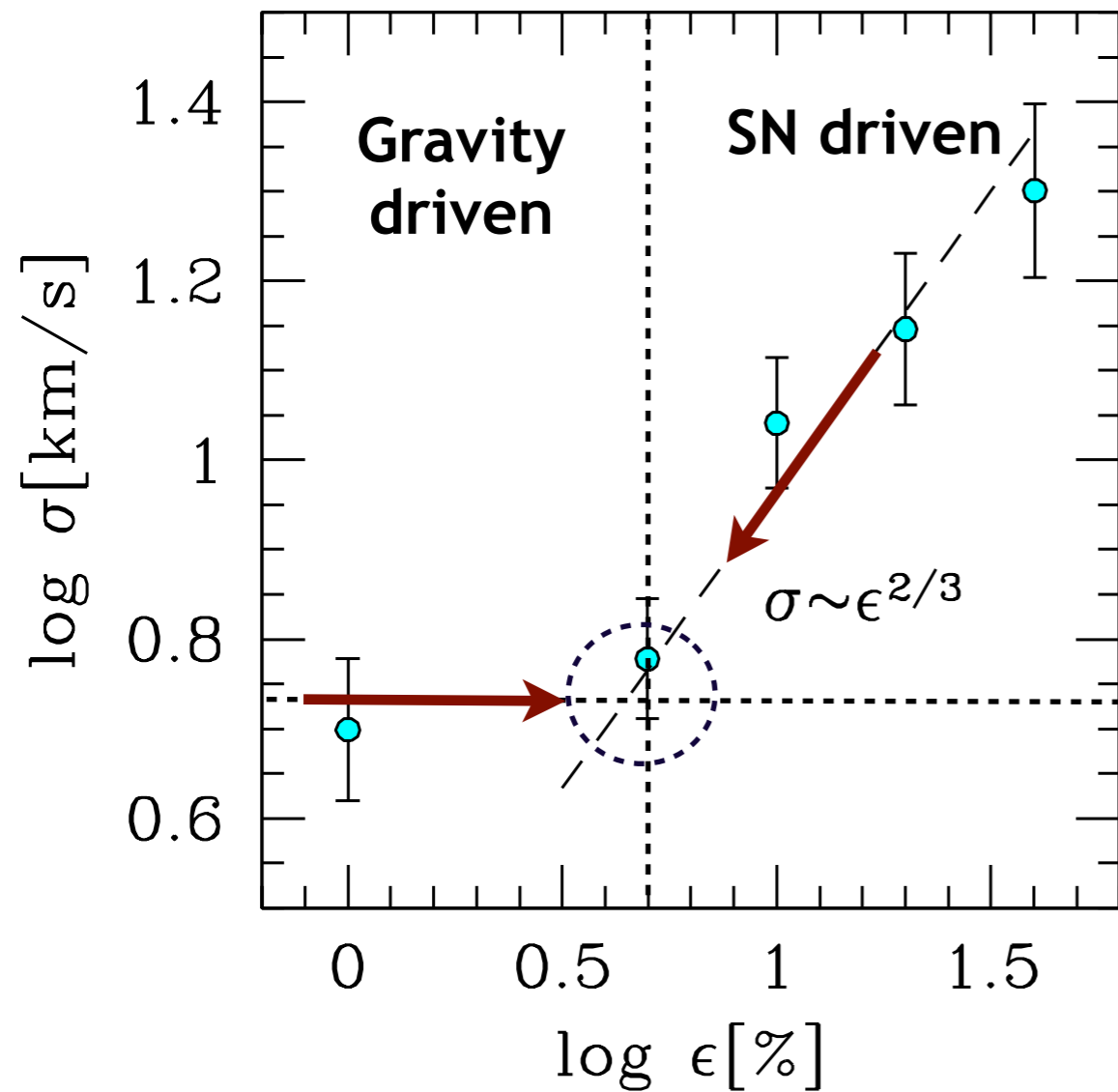


## *Summary*

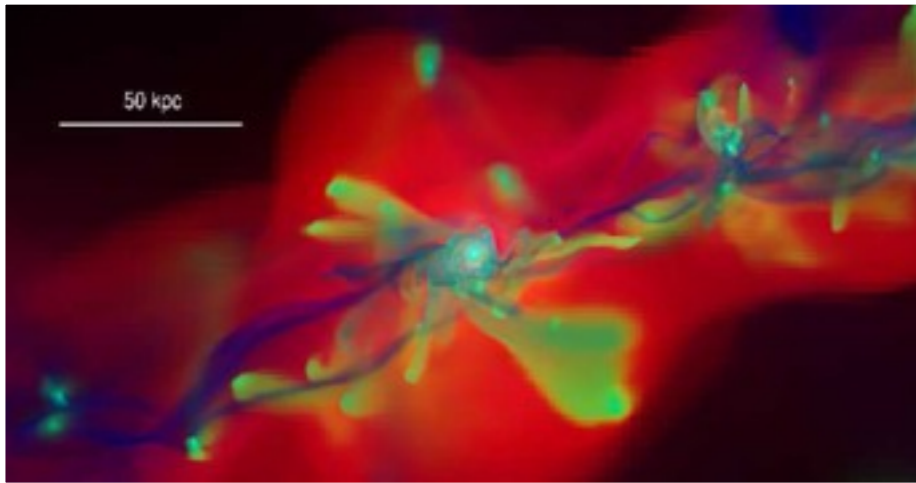
- The **molecular web** is regulated by **gravitational instabilities** and **stellar feedback**.
- The **star formation timescale** is set by the timescale of **global disk instabilities** and the **efficiency** of star formation.
- Molecular clouds are **transient structures** in the **molecular web**
  - **Stellar feedback** disrupts bound cloud regions and drives **turbulence** in the molecular web
  - **Cloud-cloud collisions** drive internal cloud turbulence, **stabilising** clouds against gravitational collapse
- In the **gravity-driven mode** turbulence is regulated by  $Q \approx 1$  leading to **massive, rotating cloud complexes** and **massive star clusters**
- In the **feedback-driven mode** turbulence is regulated by **stellar feedback** leading to  $Q > 1$  and a **power-spectrum** of cloud masses, with highly turbulent clouds and **negligible rotation**.

**Galaxies might prefer to live in the transition region from gravity-driven to stellar feedback driven turbulence  star formation efficiency**

# What determines the star formation efficiency?

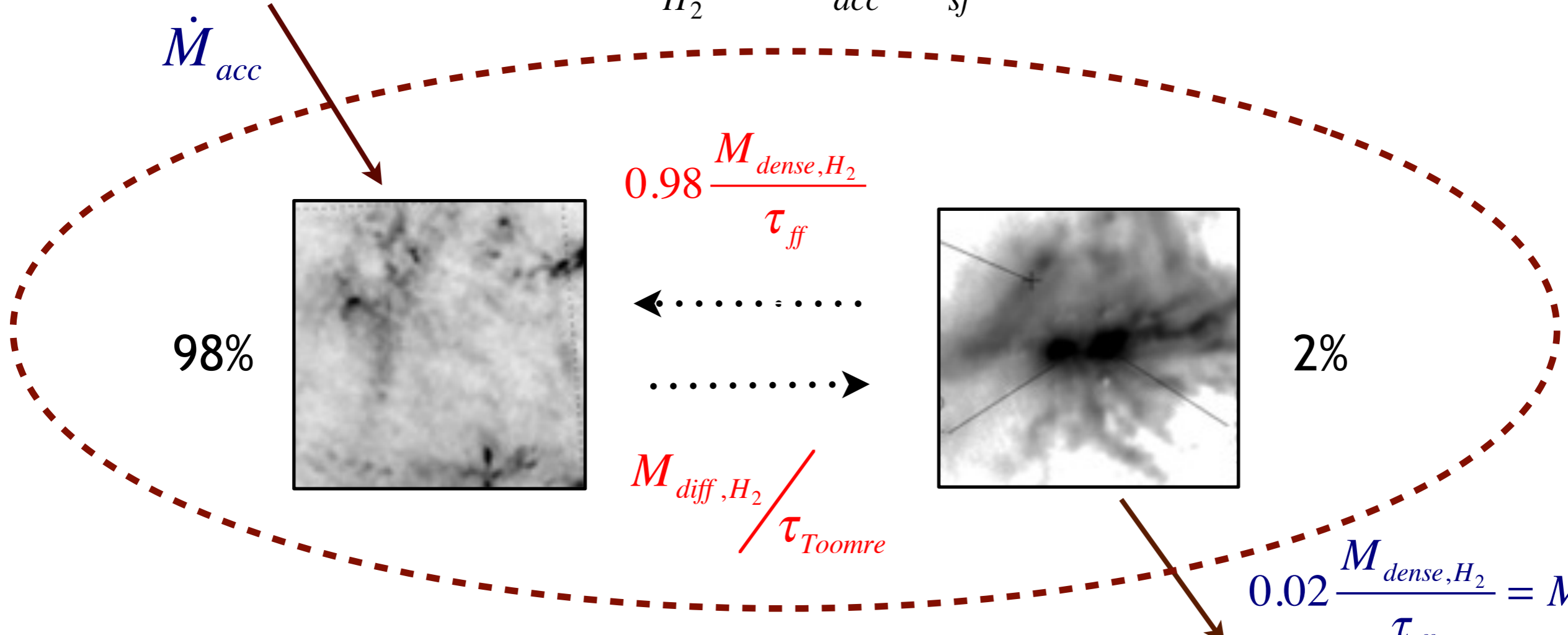


# Self-regulated star formation



$$M_{H_2} = \dot{M}_{acc} \cdot \tau_{sf}$$

$\dot{M}_{acc}$

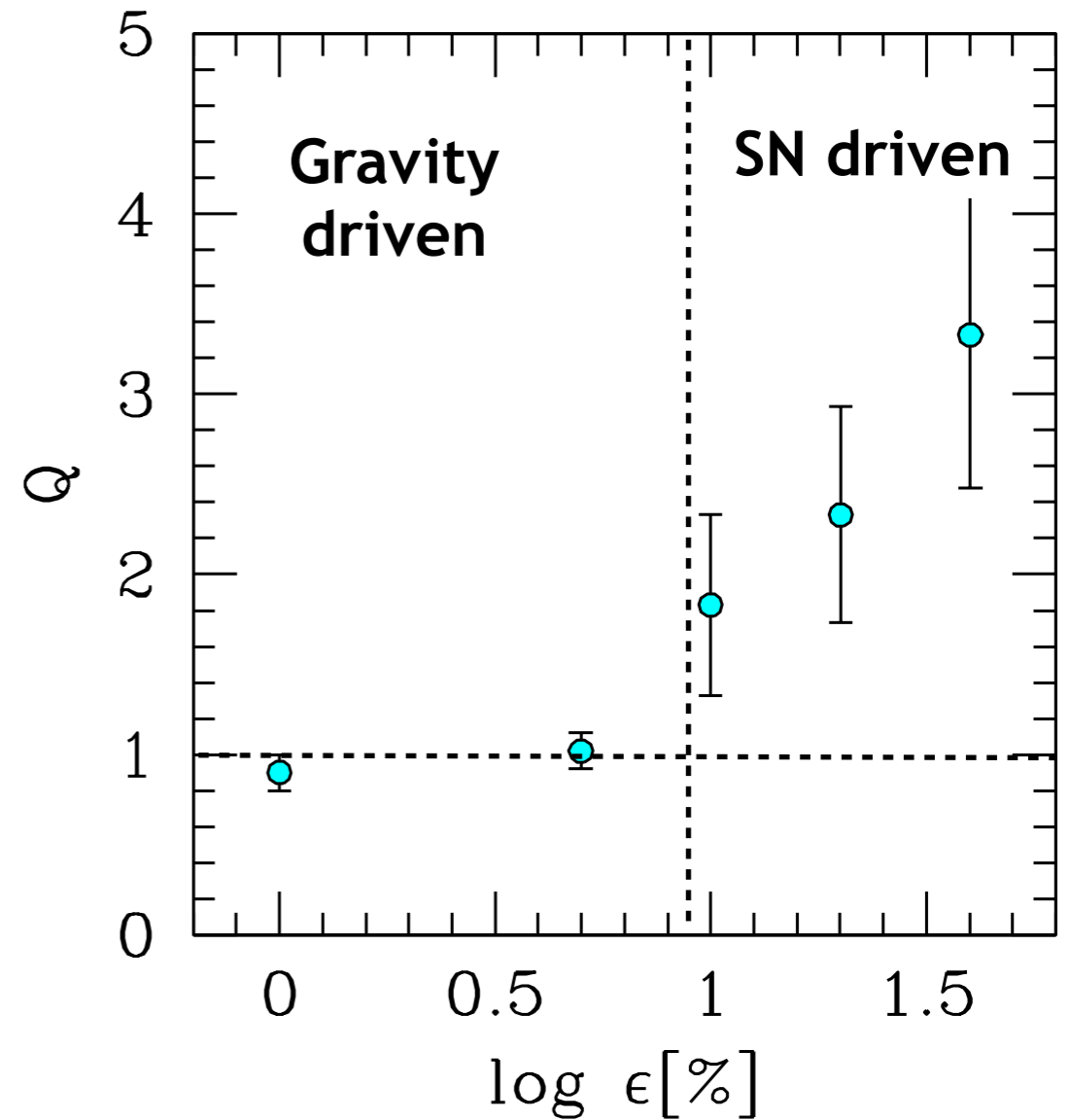
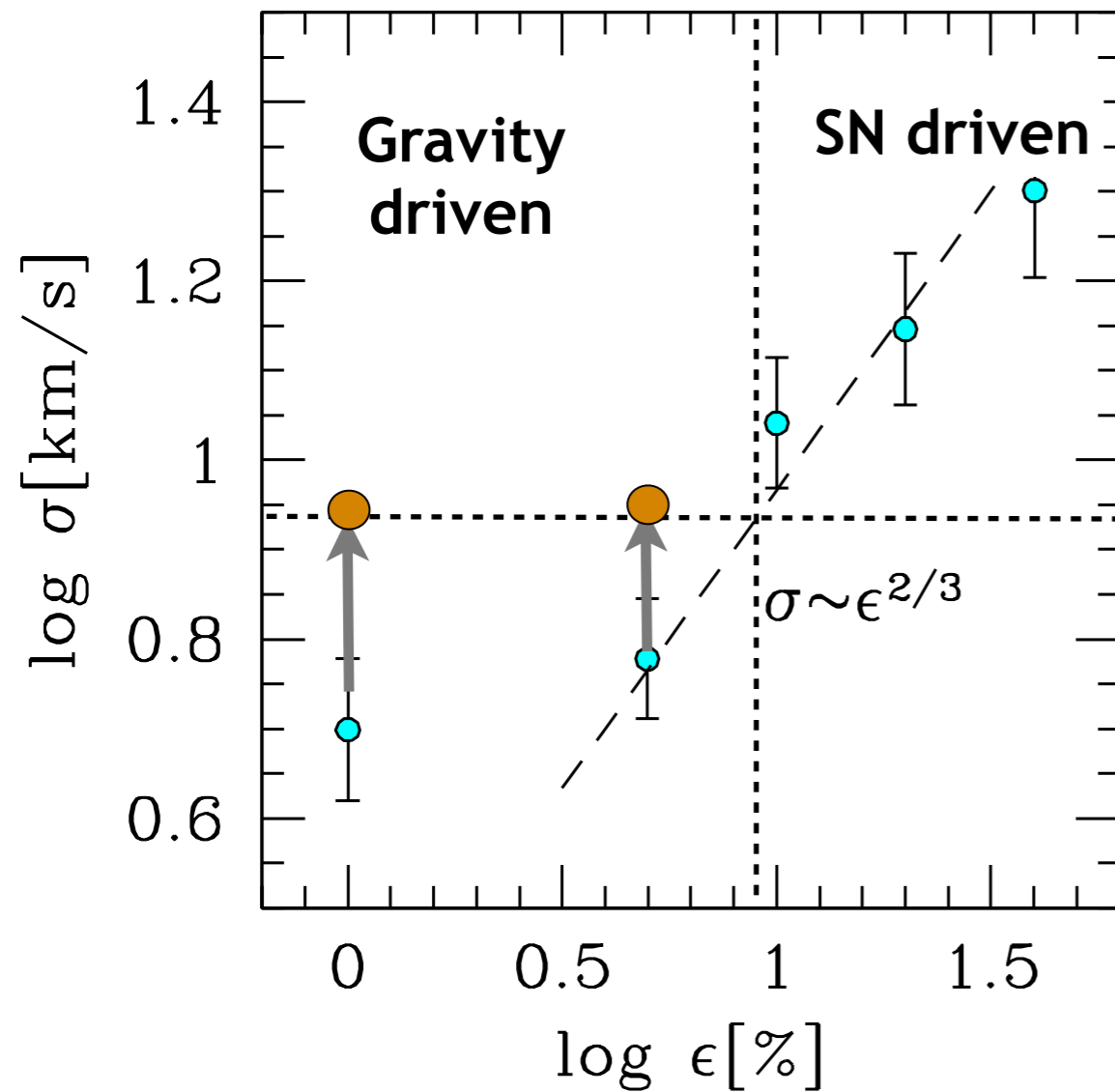


$$0.02 \frac{M_{dense,H_2}}{\tau_{ff}} = \dot{M}_{acc}$$

$$\tau_{sf} = 50 \cdot \tau_{Toomre} = \frac{1}{\epsilon \cdot \kappa}$$



## Higher gas surface densities



The **gravity driven mode** becomes more dominant for higher gas surface densities.

## Growth rate of gravitational instabilities:

$$\tau_{\text{Toomre}} = \frac{\sigma}{\pi G \Sigma} = \kappa^{-1} = \left(\sqrt{2} \Omega\right)^{-1} \rightarrow \tau_{\text{Toomre}} = 0.1 \cdot \tau_{\text{orb}} \approx 2 \cdot 10^7 \text{ yrs}$$

$$Q = 1$$

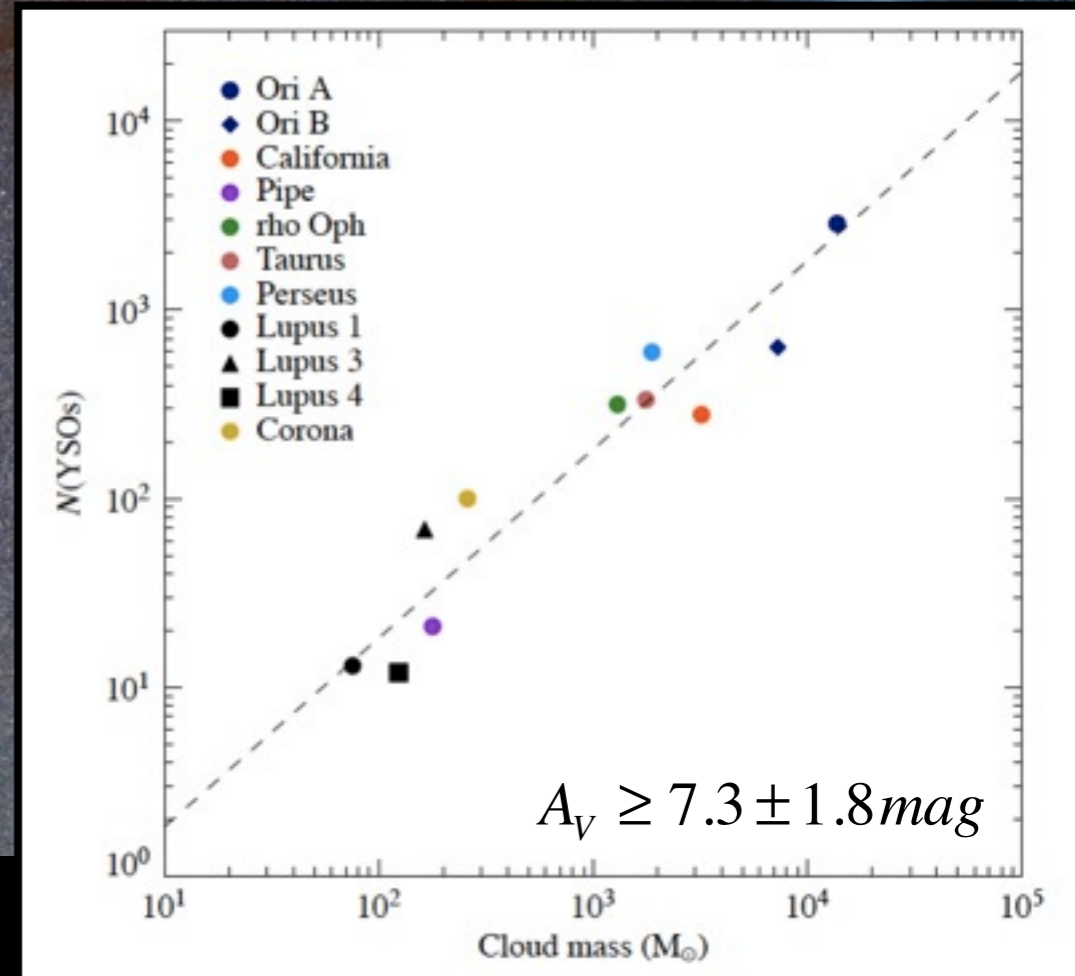
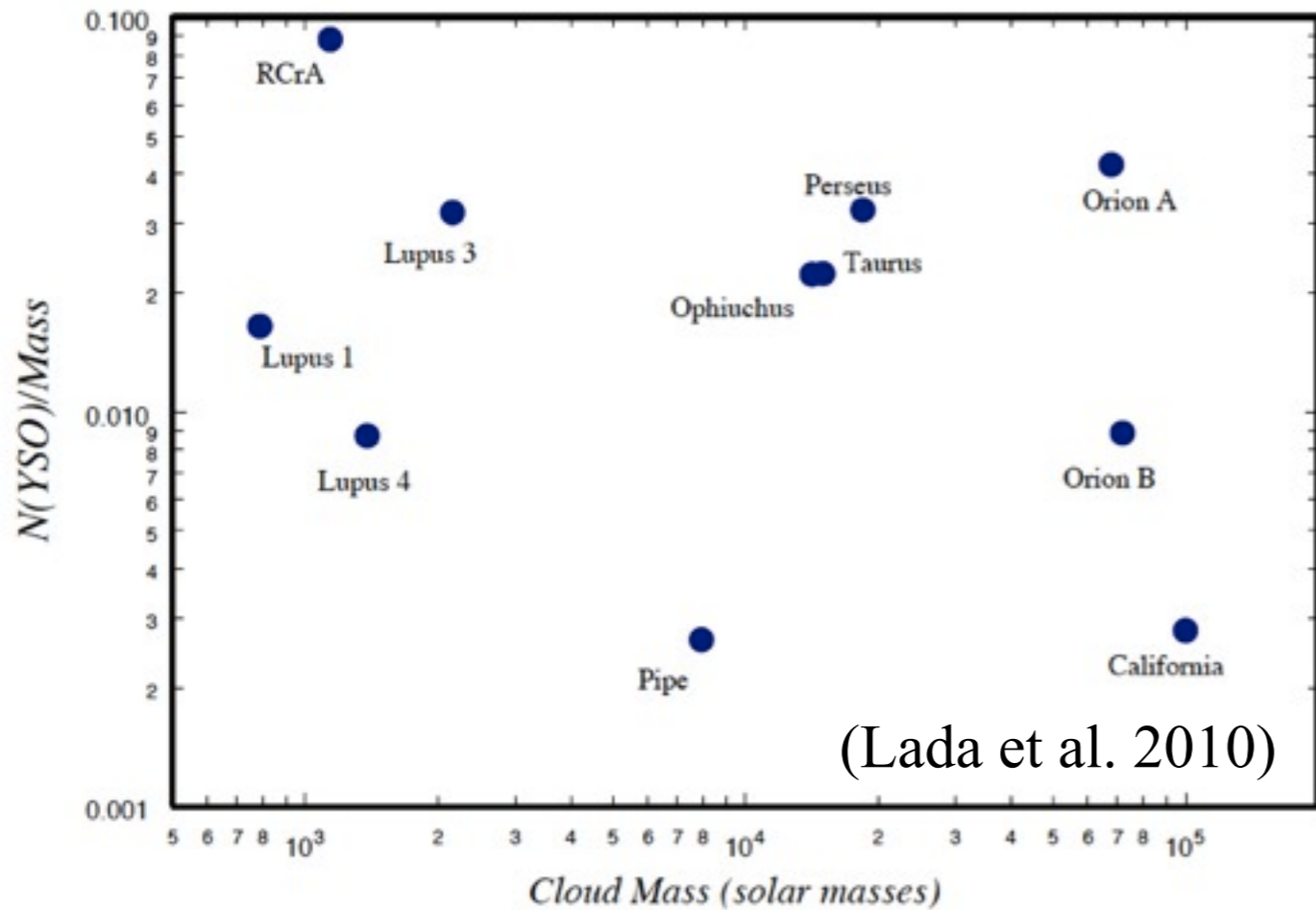
$$\tau_{\text{orb}} \sim \frac{R_{\text{vir}}}{V_{\text{vir}}} \sim H^{-1}$$

$$SFR \approx \varepsilon \frac{M_{\text{H}_2}}{\tau_{\text{Toomre}}} \approx \frac{M_{\text{H}_2}}{10^9 \text{ yrs}}$$

$$\varepsilon \approx 0.02$$

**What determines the star formation efficiency?**

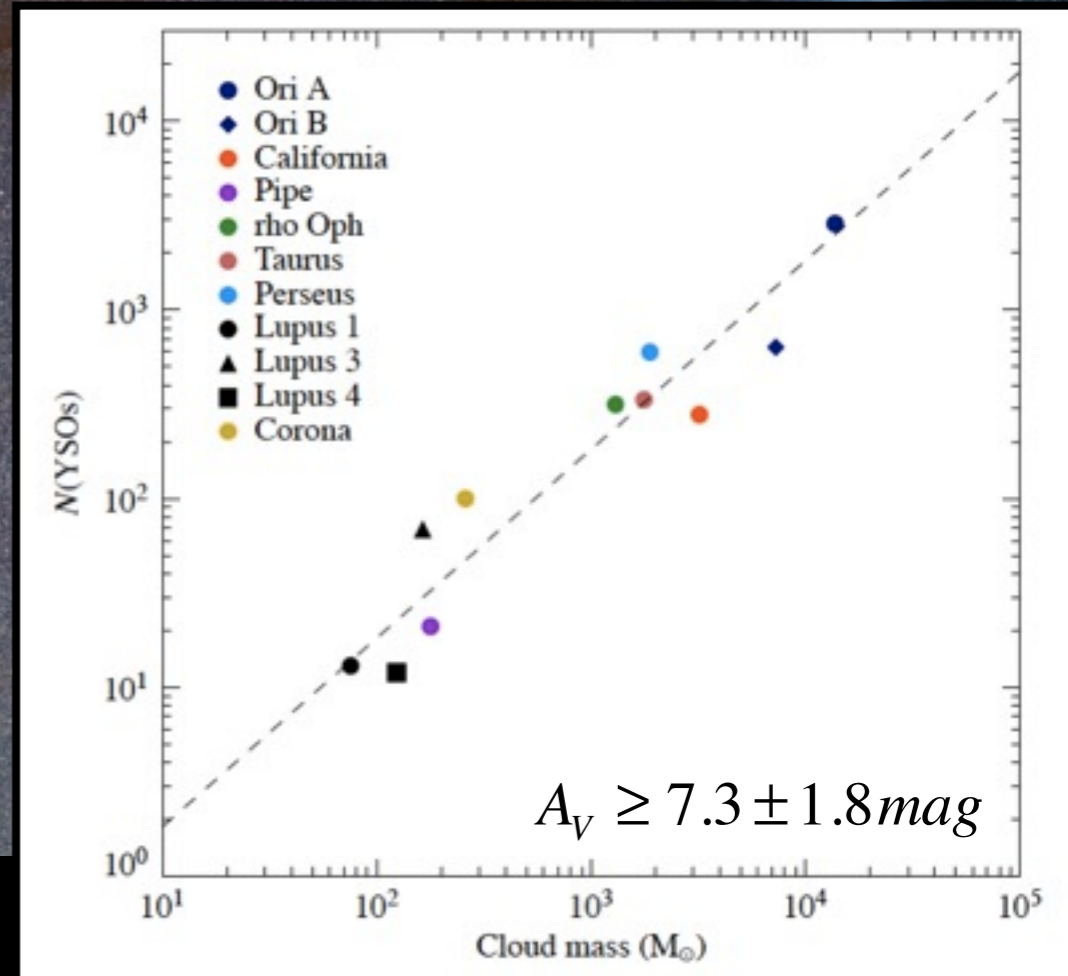
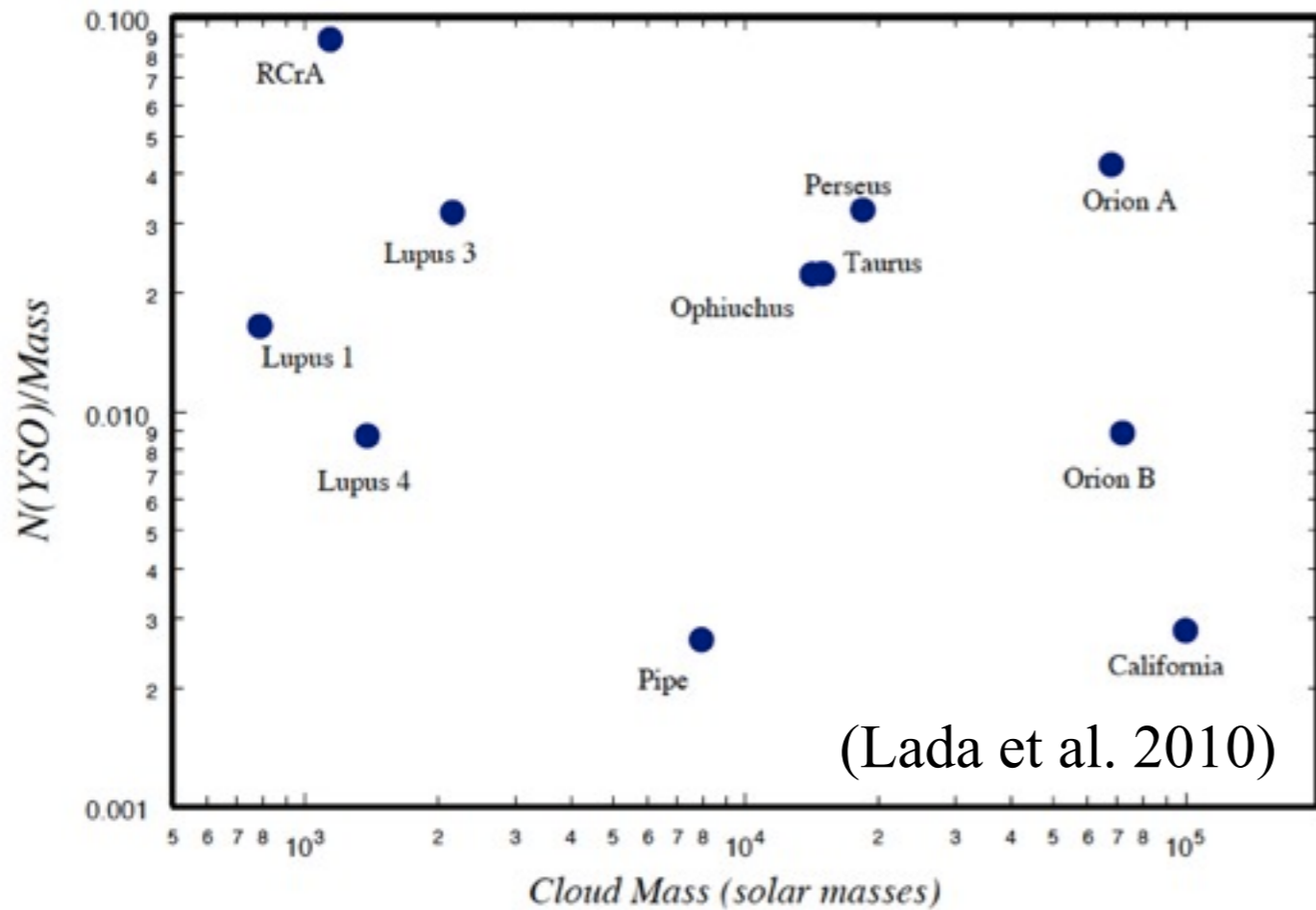
$$N(\text{YSOs})_{\text{Oph}} = 10 \times N(\text{YSOs})_{\text{Pipe}}$$



$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 4 \cdot 10^5 \text{ yrs}$$

$$\rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}, \text{H}_2}}{\tau_{\text{ff}}}$$

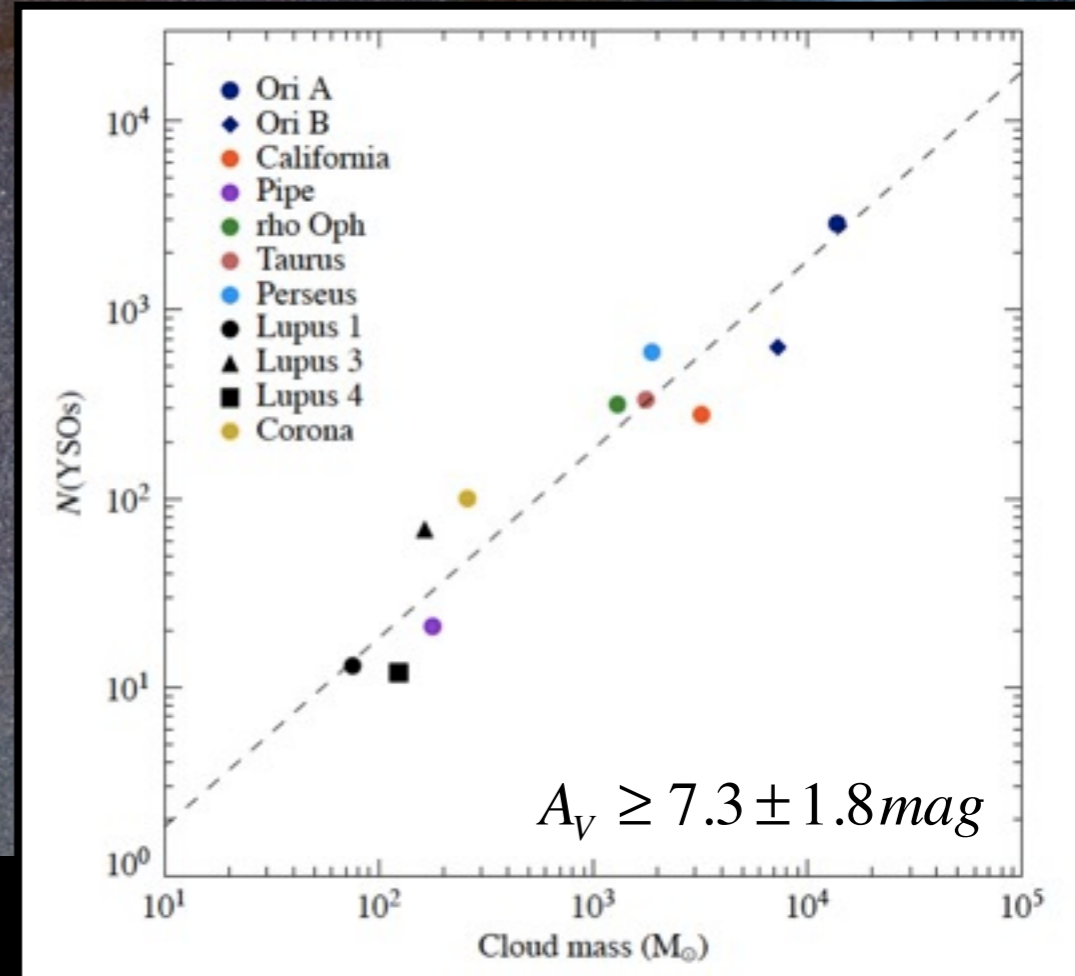
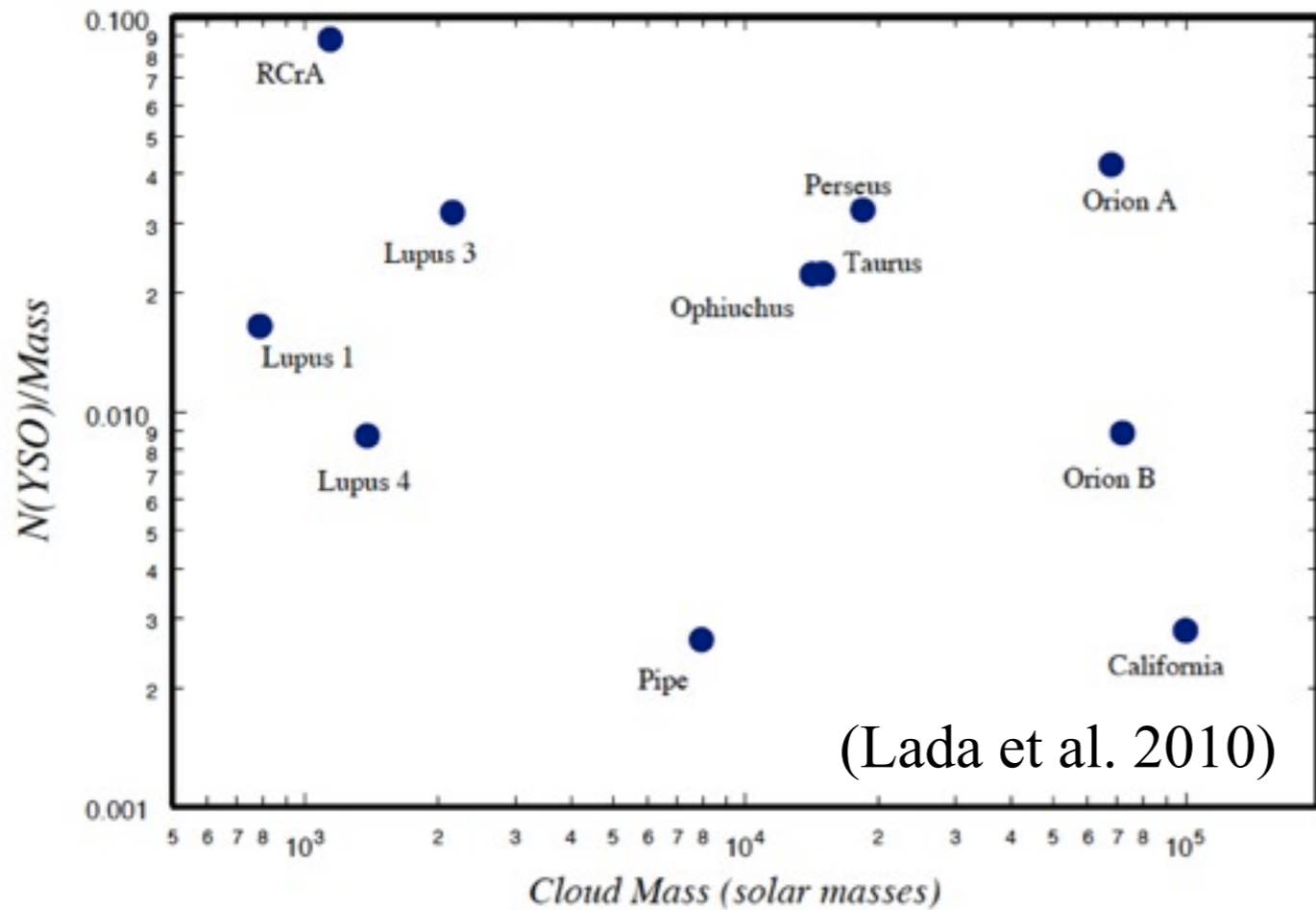
$$N(\text{YSOs})_{\text{Oph}} = 10 \times N(\text{YSOs})_{\text{Pipe}}$$



$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 4 \cdot 10^5 \text{ yrs}$$

$$\rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}, \text{H}_2}}{\tau_{\text{ff}}} \equiv \frac{M_{\text{dense}, \text{H}_2}}{\tau_{\text{sf}, \text{dense}}}$$

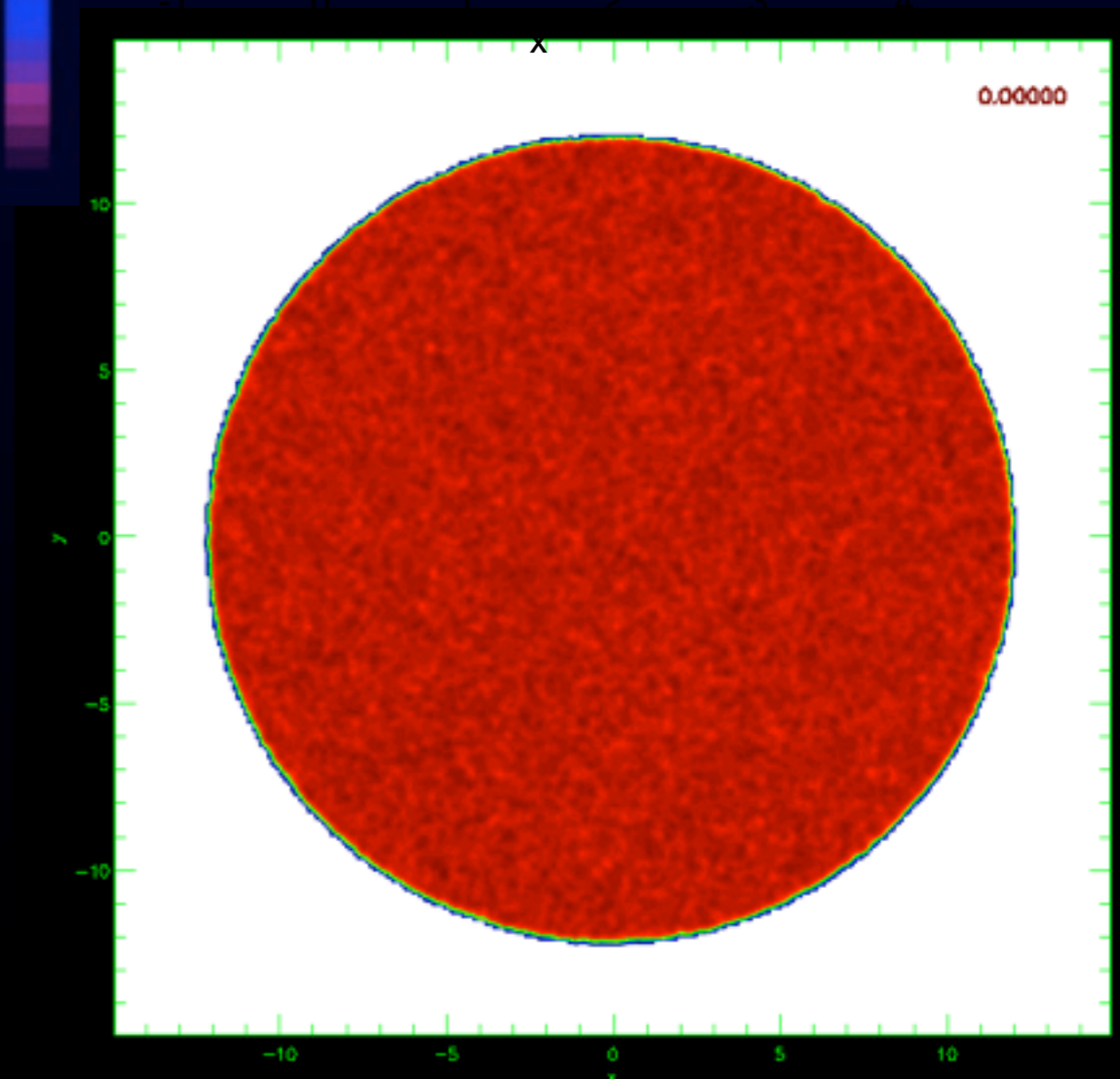
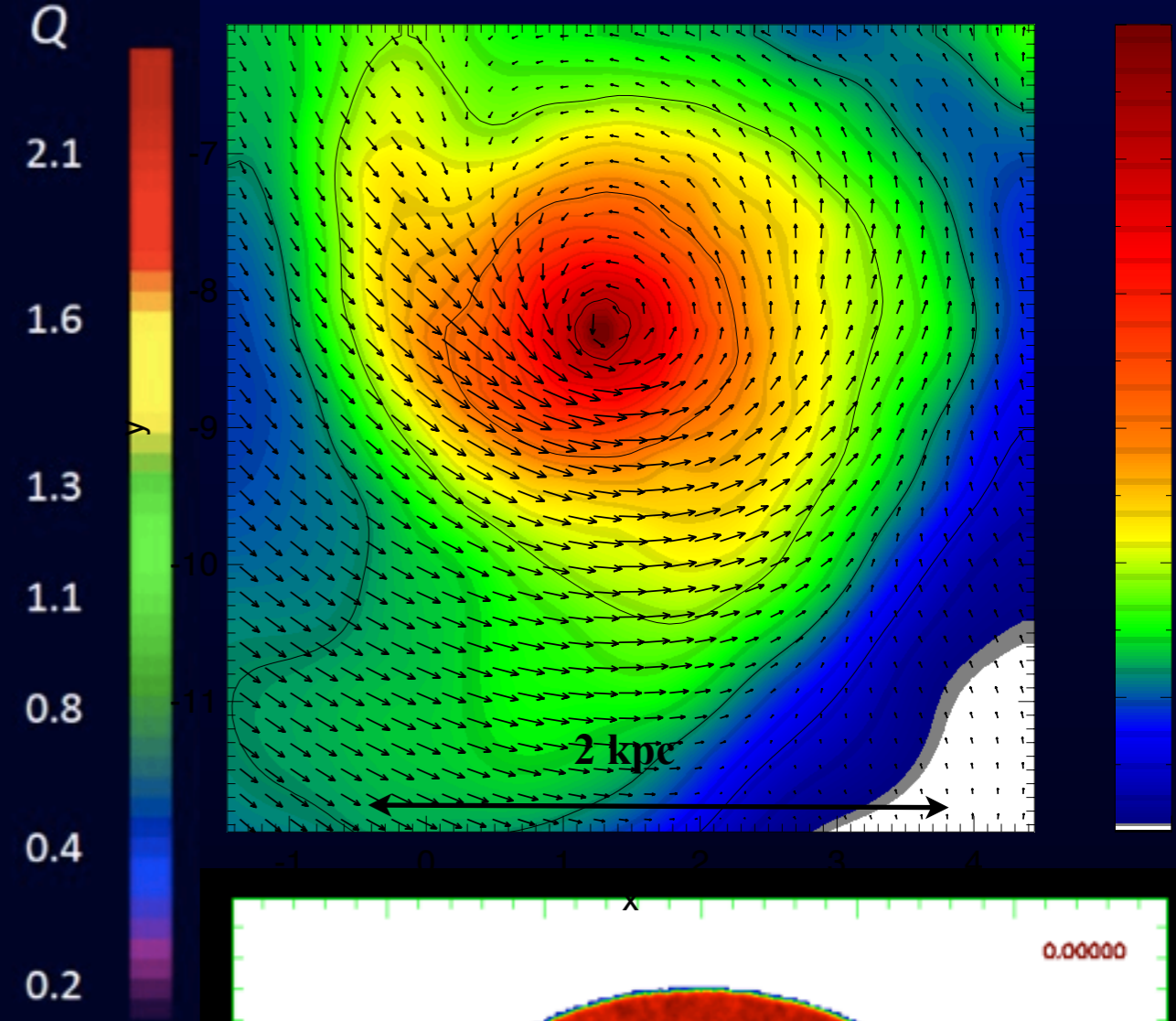
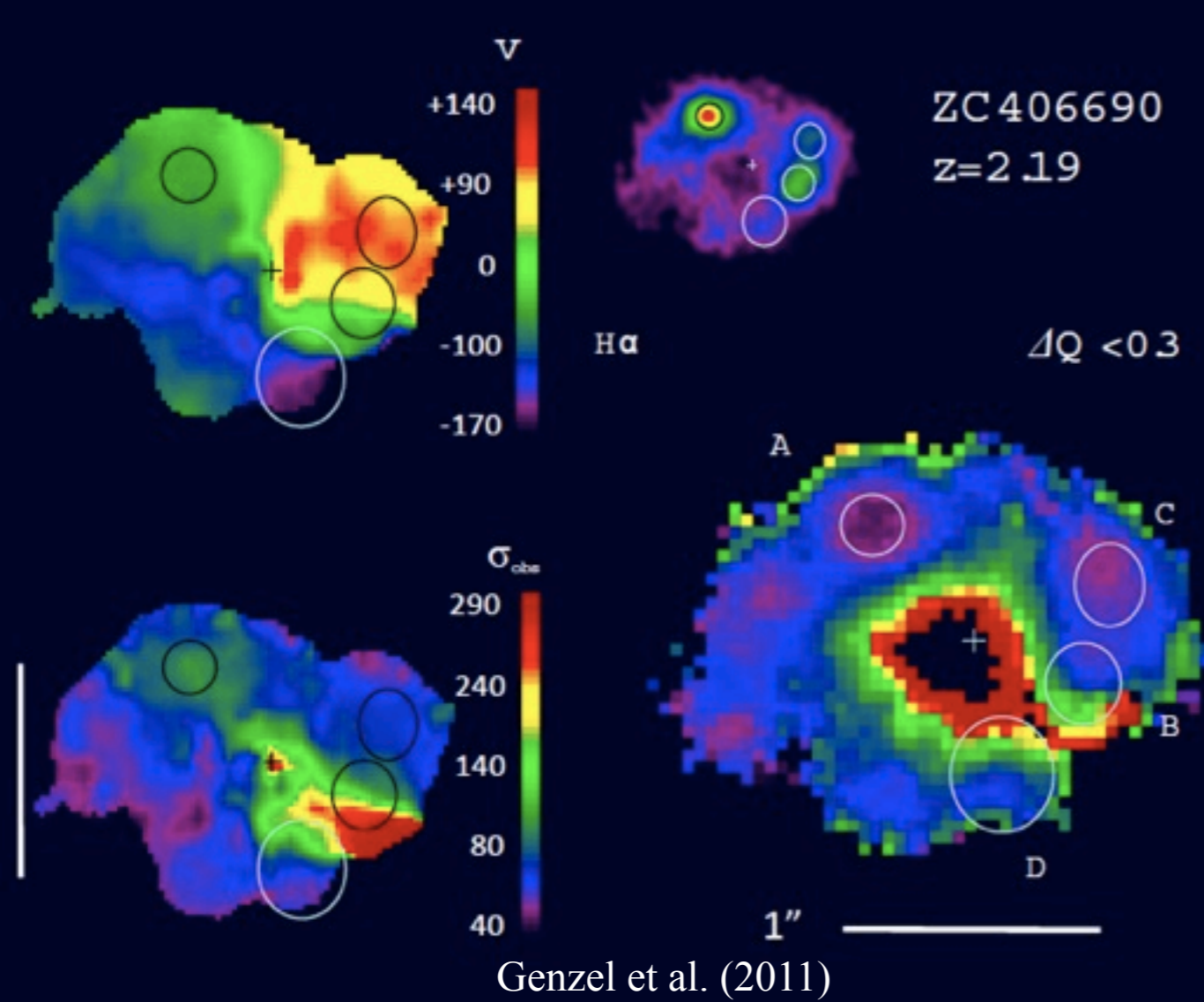
$$N(\text{YSOs})_{\text{Oph}} = 10 \times N(\text{YSOs})_{\text{Pipe}}$$

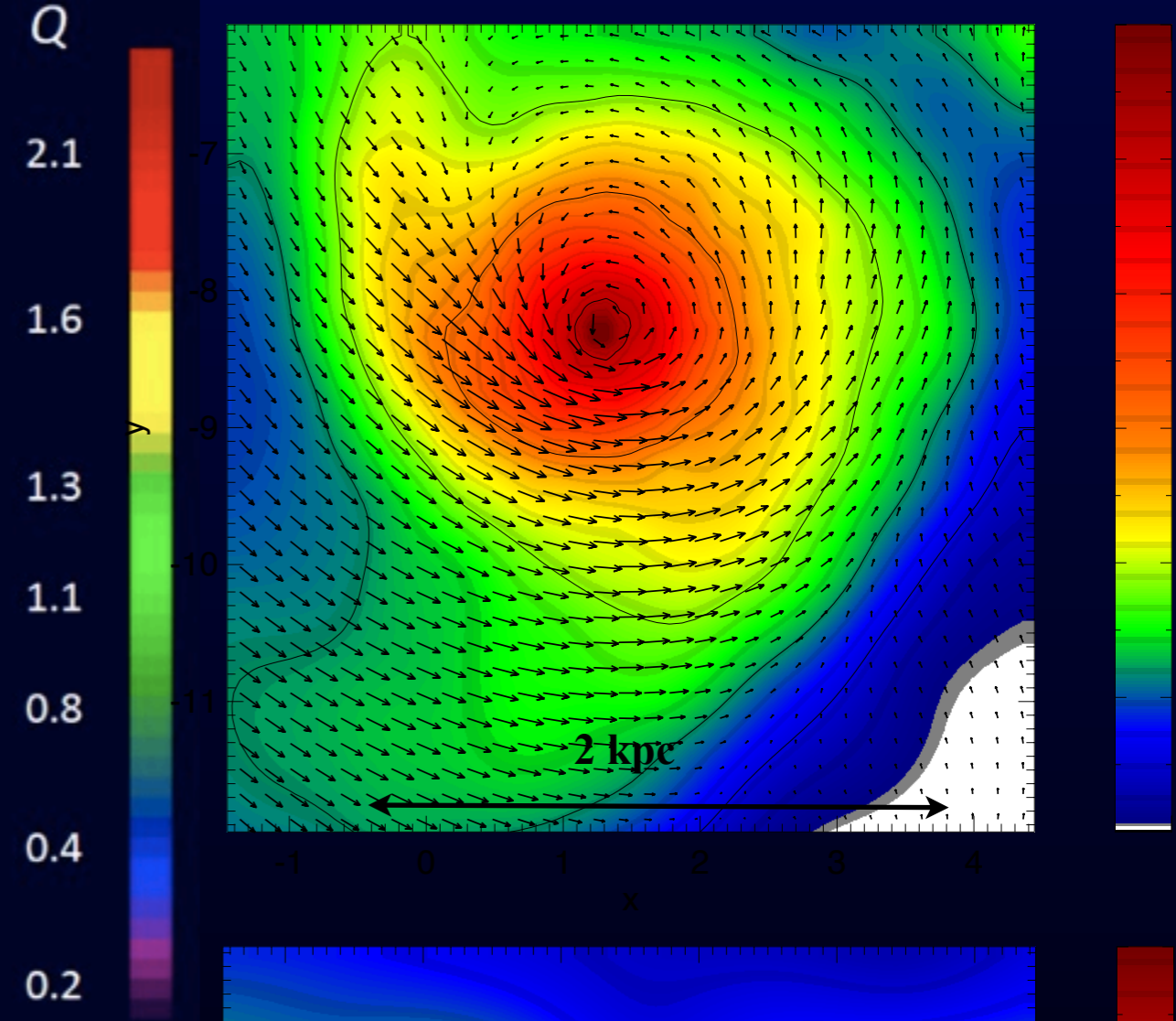
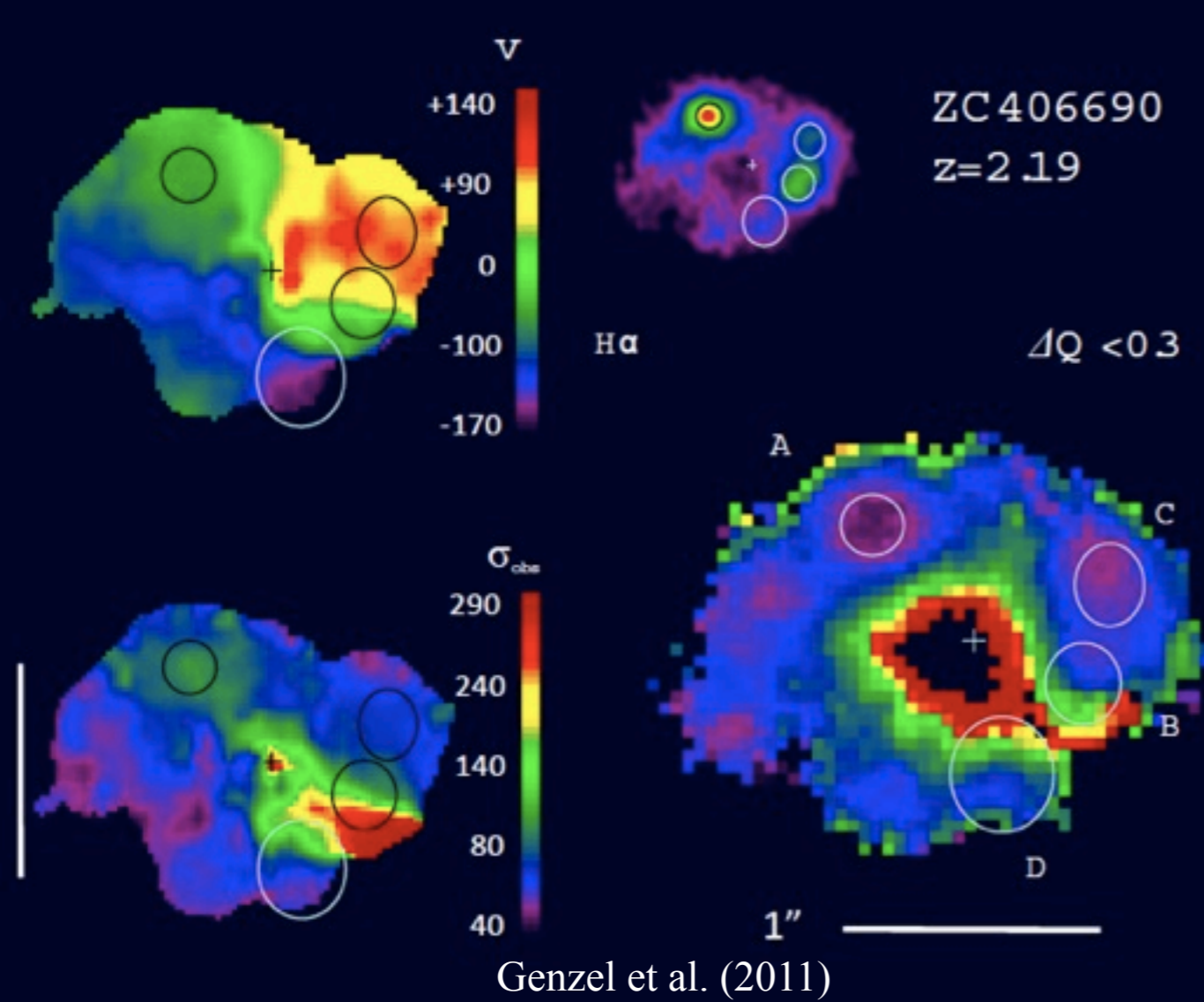


$$n_{A_V=7.3} \approx 10^4 \text{ cm}^{-3} \rightarrow \tau_{\text{ff}} \approx 4 \cdot 10^5 \text{ yrs}$$

$$\rightarrow \text{SFR} \approx 0.02 \frac{M_{\text{dense}, \text{H}_2}}{\tau_{\text{ff}}} \rightarrow \tau_{\text{sf}, \text{dense}} \approx 2 \cdot 10^7 \text{ yrs}$$



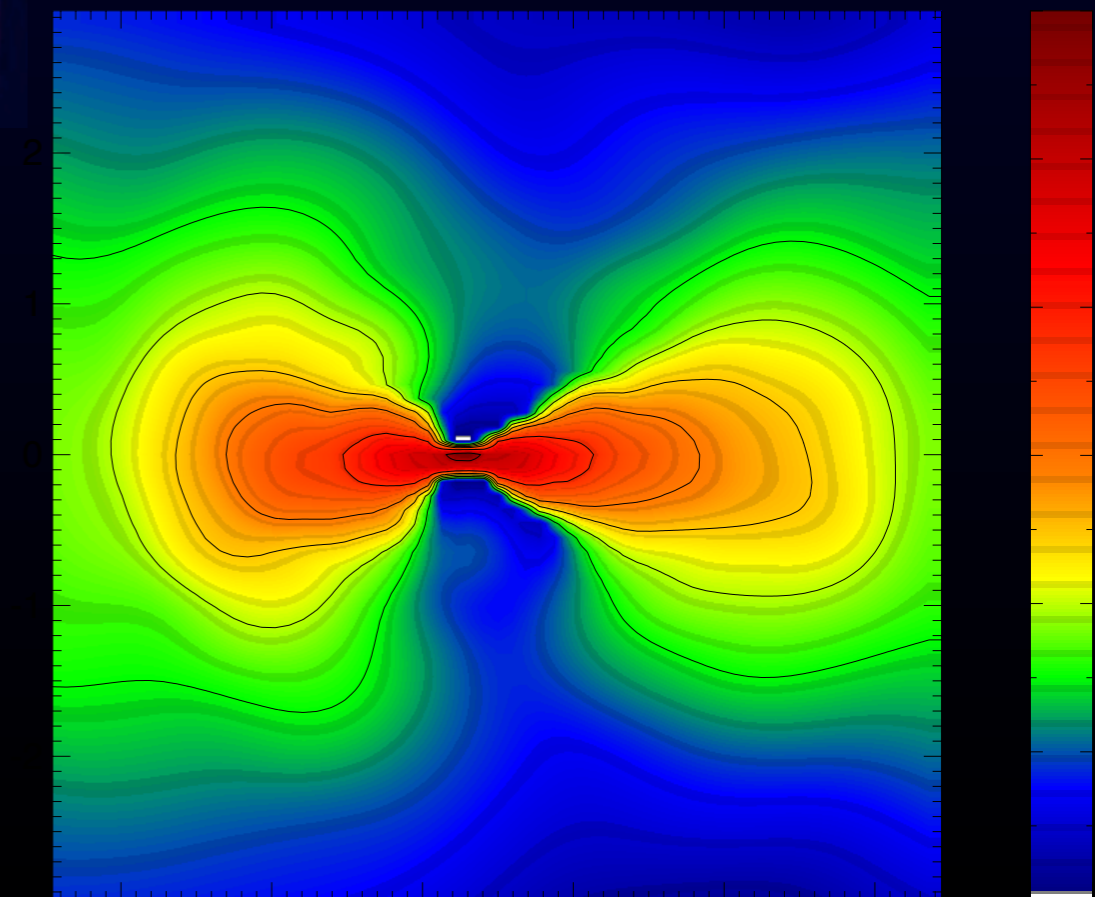




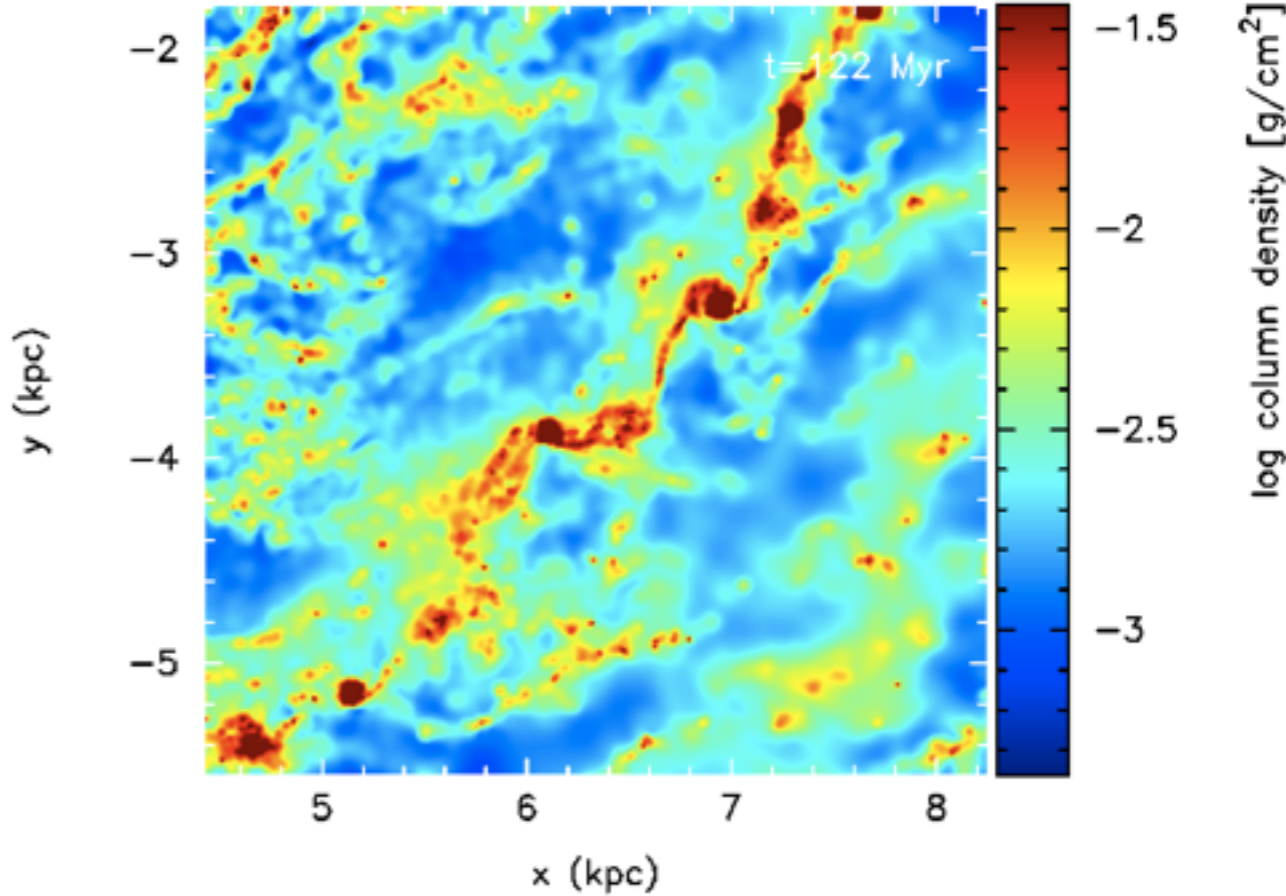
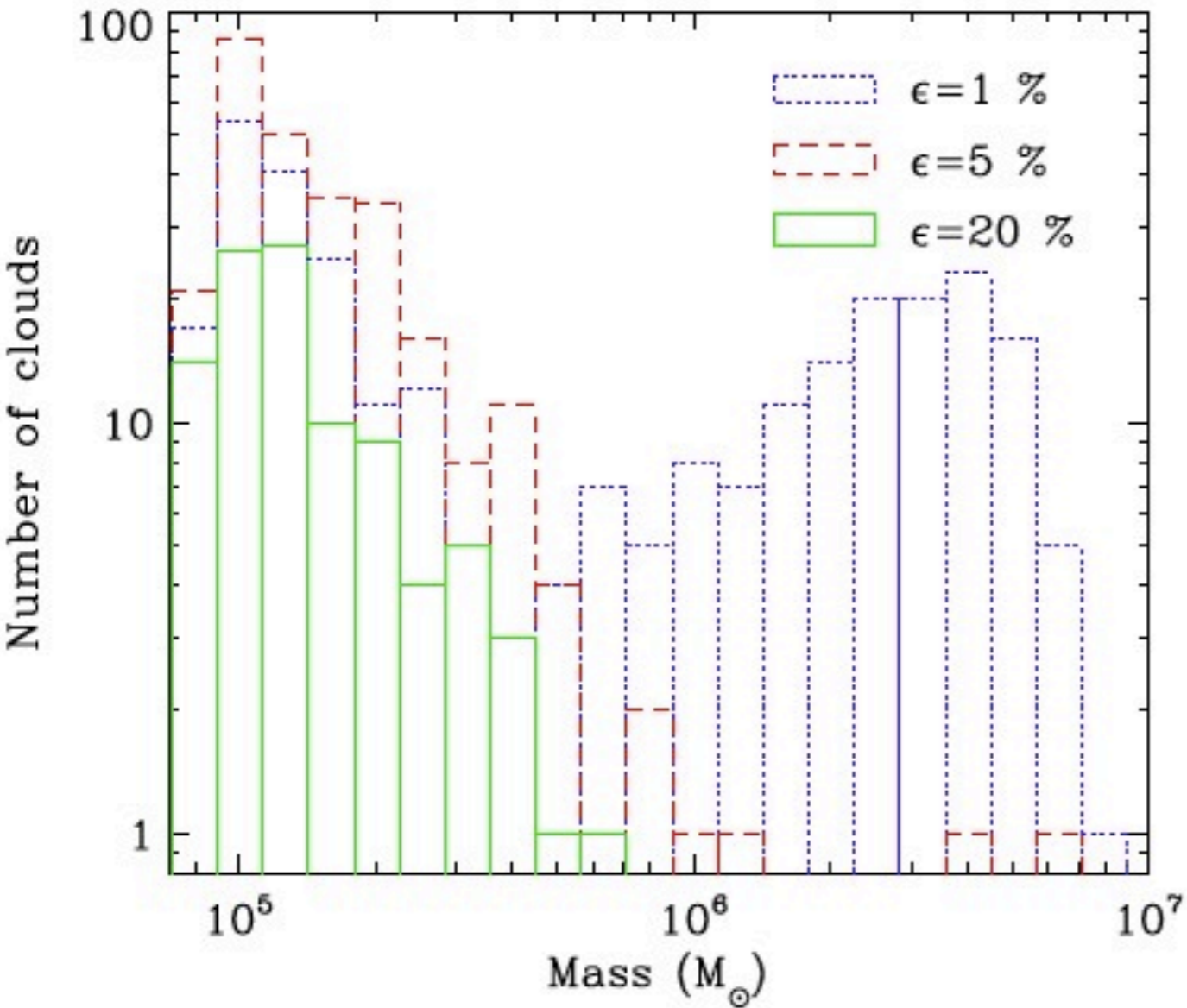
## Rotationally supported minidisks

Expected:  $v_{\text{rot}} \approx 200 \text{ km / s}$

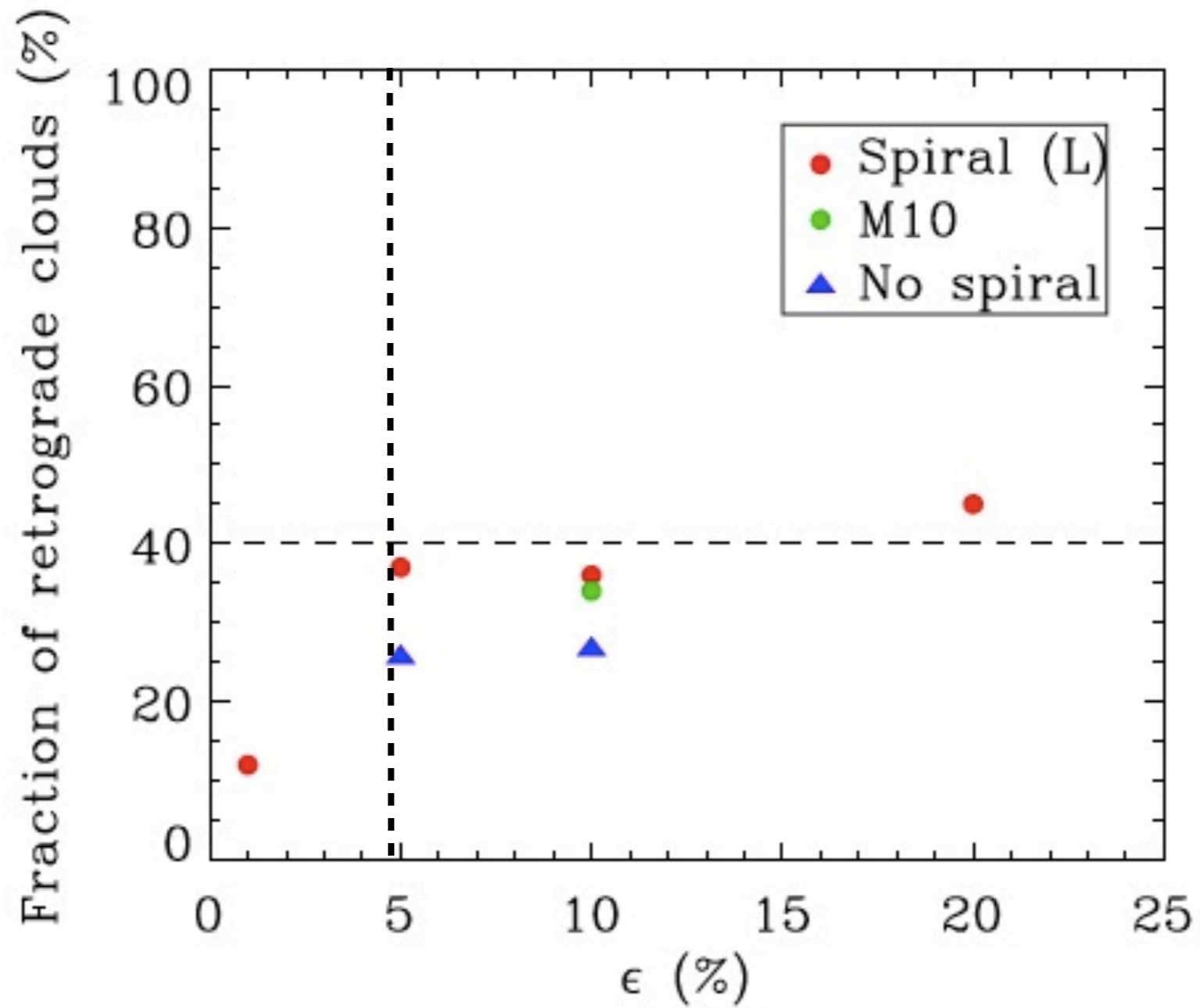
Observed:  $v_{\text{rot}} \approx 10 - 40 \text{ km / s}$



# Gravity driven mode: formation of giant clumps



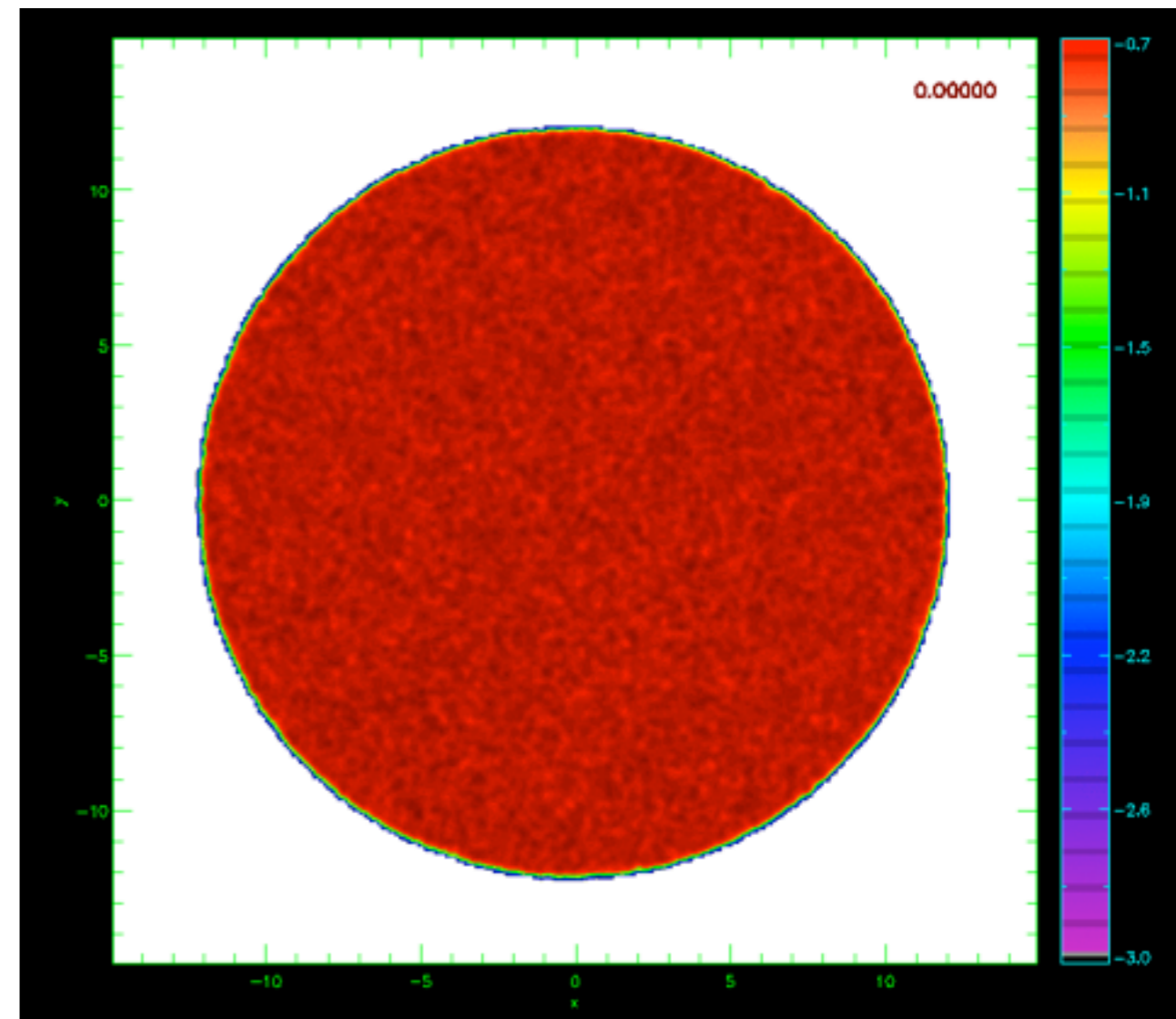
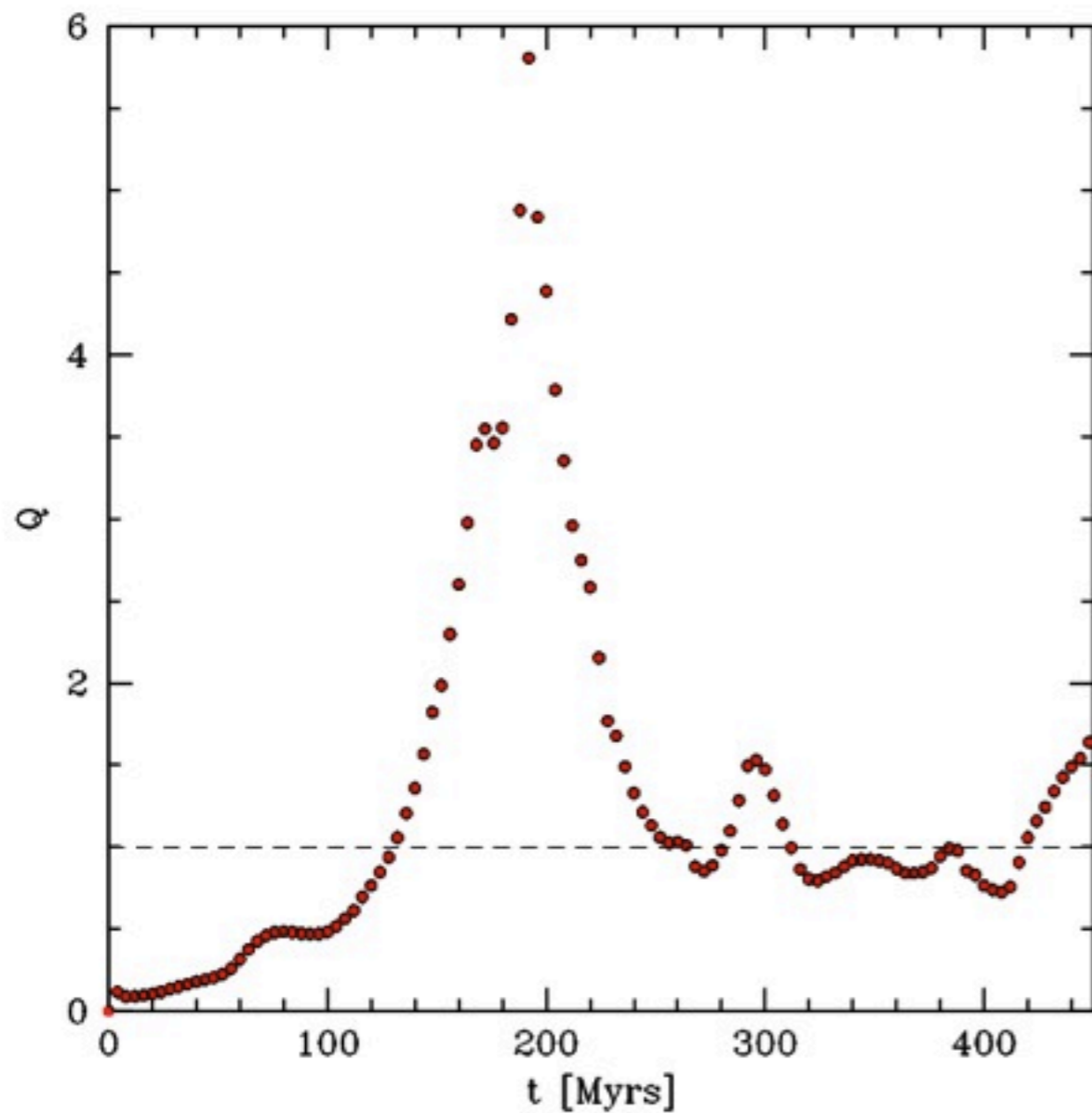
## *Fraction of retrograde clouds*



# Gravitational disk instabilities

(Toomre 1964; Goldreich & Lynden-Bell 65; Elmegreen 94; Kim & Ostriker 01, 06)

Gaseous disks will **self-regulate** themselves into a state of **marginal stability** (Dekel et al. 09; Bournaud et al. 09; Krumholz & Burkert 10; Elmegreen & Burkert 10; Genzel et al. 10, Burkert et al. 11; Dobbs et al. 11a,b)



## Growth rate of gravitational instabilities:

$$\tau_{\text{Toomre}} = \frac{\sigma}{\pi G \Sigma} = \kappa^{-1} = \left( \sqrt{2} \Omega \right)^{-1} \rightarrow \tau_{\text{Toomre}} = 0.1 \cdot \tau_{\text{orb}} \approx 2 \cdot 10^7 \text{ yrs}$$

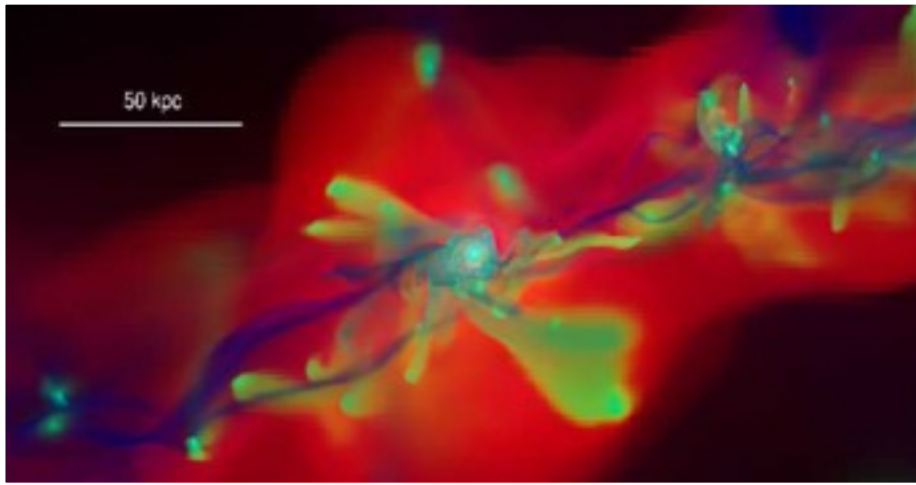
$$Q = 1$$

$$\tau_{\text{orb}} \sim \frac{R_{\text{vir}}}{V_{\text{vir}}} \sim H^{-1}$$

$$\tau_{\text{Toomre}} \approx \tau_{\text{sf, dense}}$$

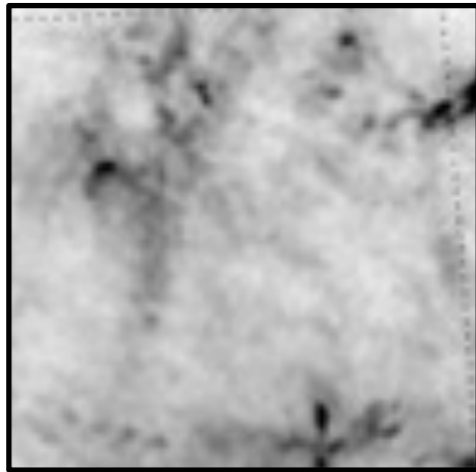
$$\frac{M_{\text{diff, H}_2}}{\tau_{\text{Toomre}}} \approx \frac{M_{\text{dense, H}_2}}{\tau_{\text{ff}}} \rightarrow \frac{M_{\text{dense, H}_2}}{M_{\text{diff, H}_2}} = 0.02$$

# Self-regulated star formation



$$M_{H_2} = \dot{M}_{acc} \cdot \tau_{sf}$$

$\dot{M}_{acc}$

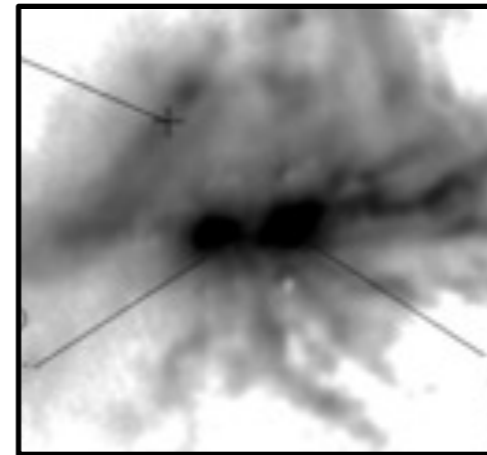


98%

$$0.98 \frac{M_{dense,H_2}}{\tau_{ff}}$$



$$\frac{M_{diff,H_2}}{\tau_{Toomre}}$$



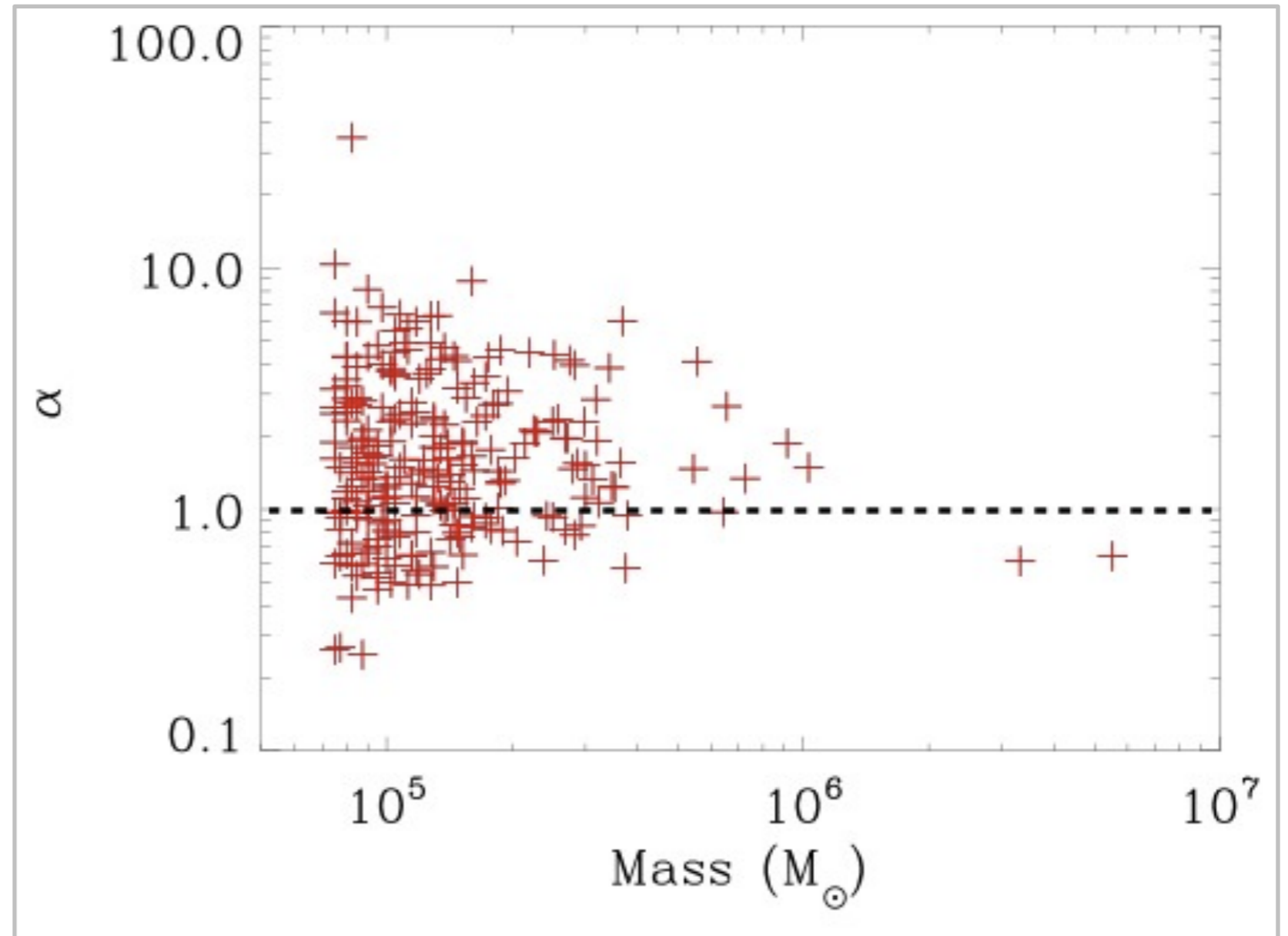
2%

$$0.02 \frac{M_{dense,H_2}}{\tau_{ff}} = \dot{M}_{acc}$$



**Most clouds are gravitationally unbound!**

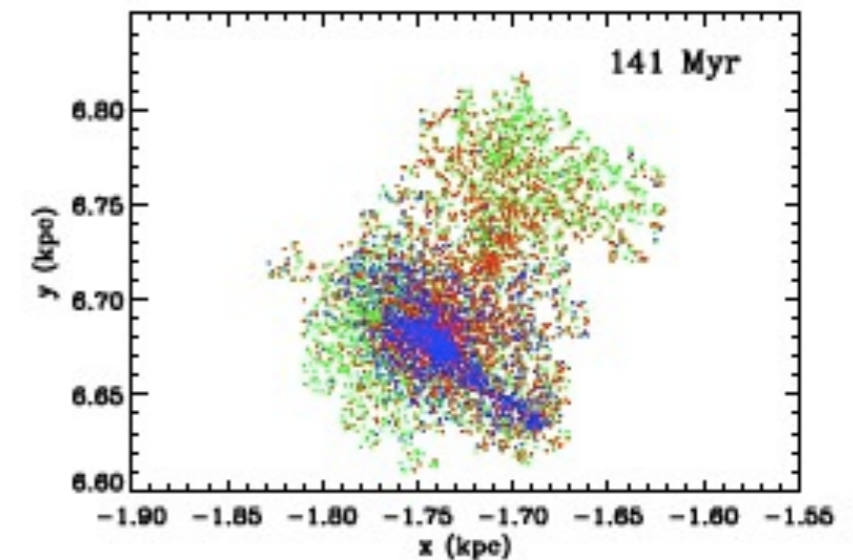
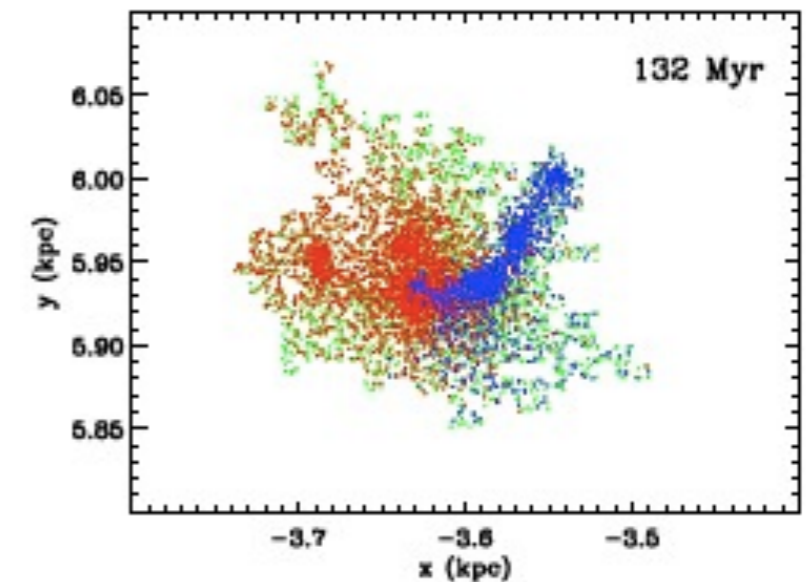
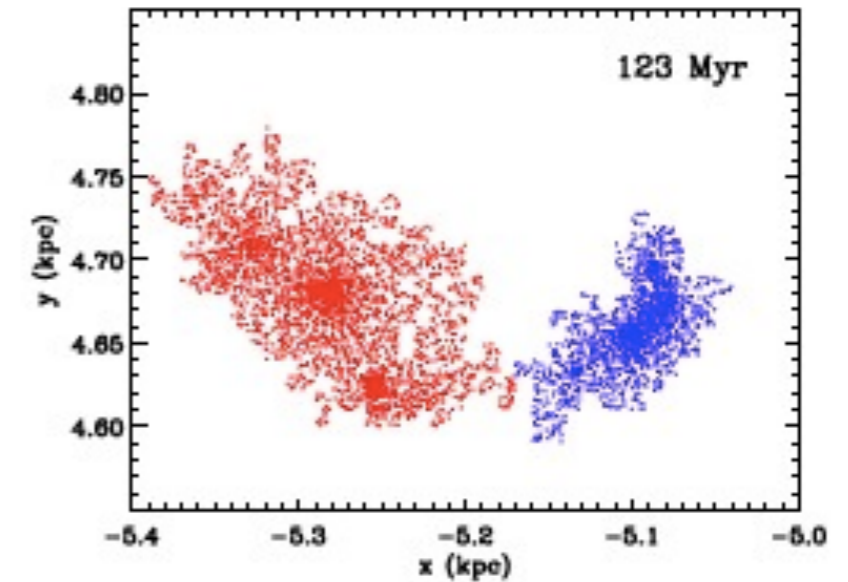
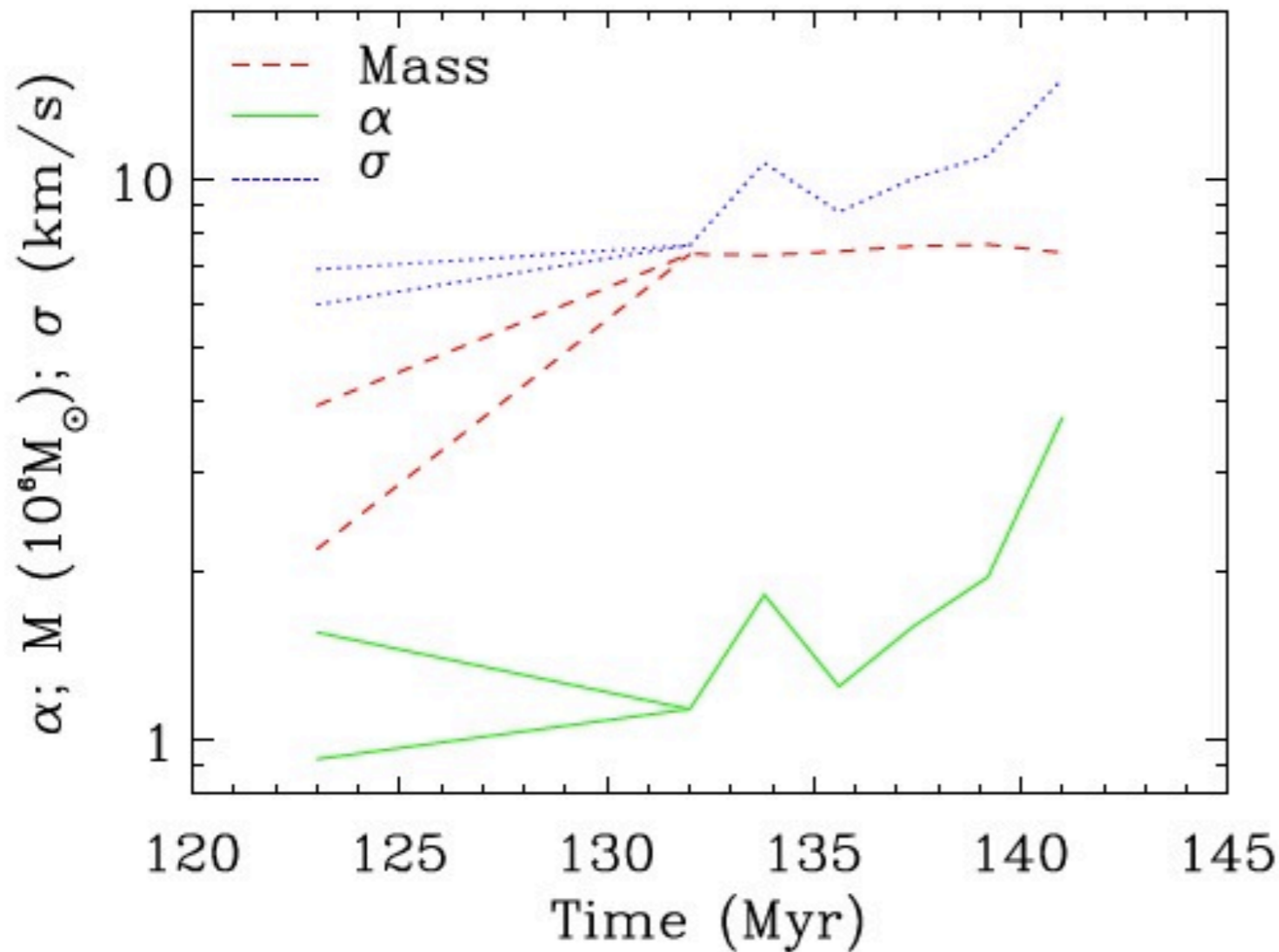
$$\alpha = \frac{5\sigma_v^2 R}{GM}$$





# Why are most clouds gravitationally unbound?

## 1. Collisions drive internal turbulence

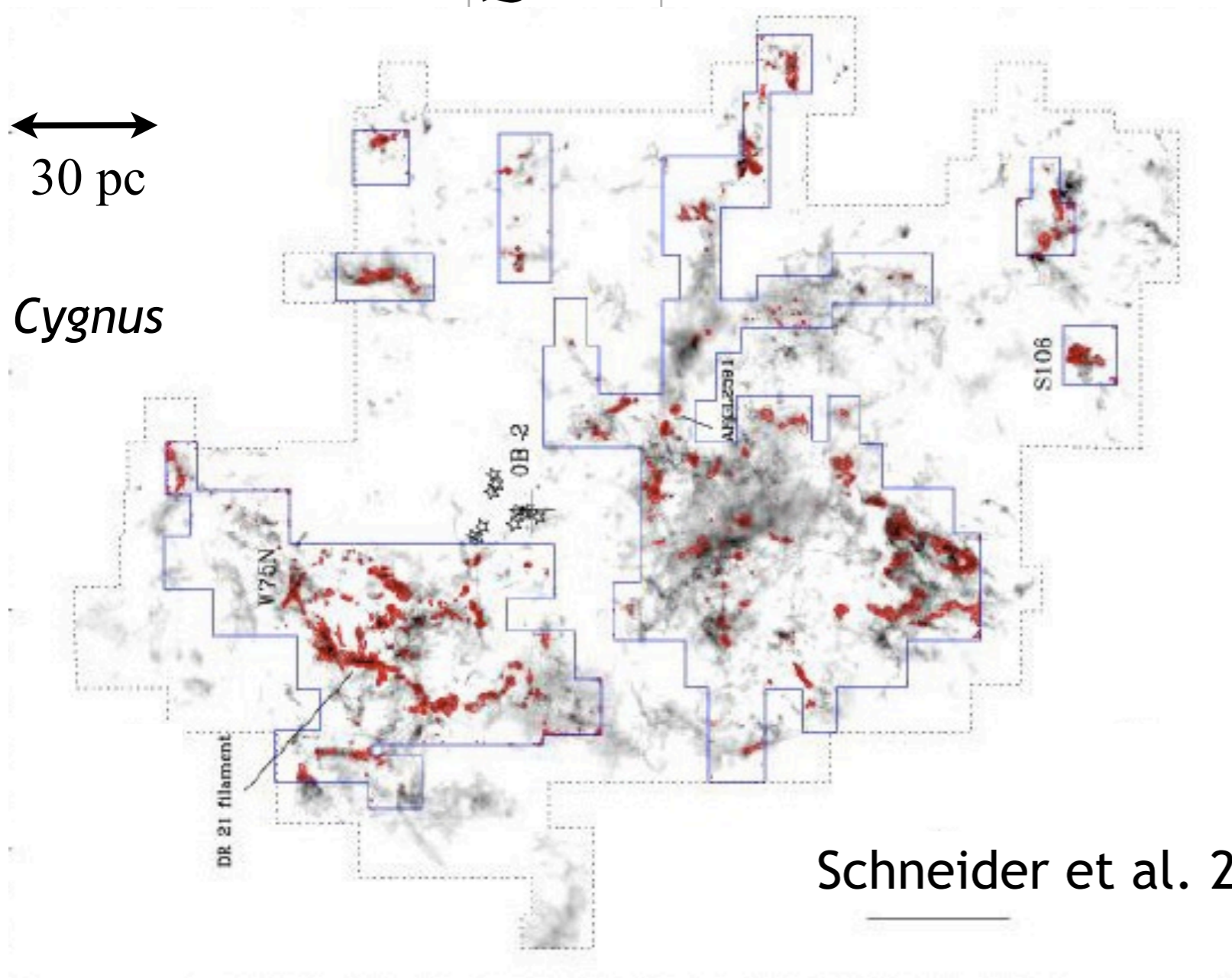


**Growth rate of gravitational instabilities:** (Silk 01; Silk&Norman 09)

$$\tau_{\text{Toomre}} = \frac{\sigma}{\pi G \Sigma} = \kappa^{-1} = \left(\sqrt{2\Omega}\right)^{-1} \rightarrow \tau_{\text{Toomre}} = 0.1 \cdot \tau_{\text{orb}} \approx 2 - 5 \cdot 10^7 \text{ yrs}$$

$$Q = 1$$

$$\tau_{\text{orb}} \sim \frac{R_{\text{vir}}}{V_{\text{vir}}} \sim H^{-1}$$



**The molecular web**

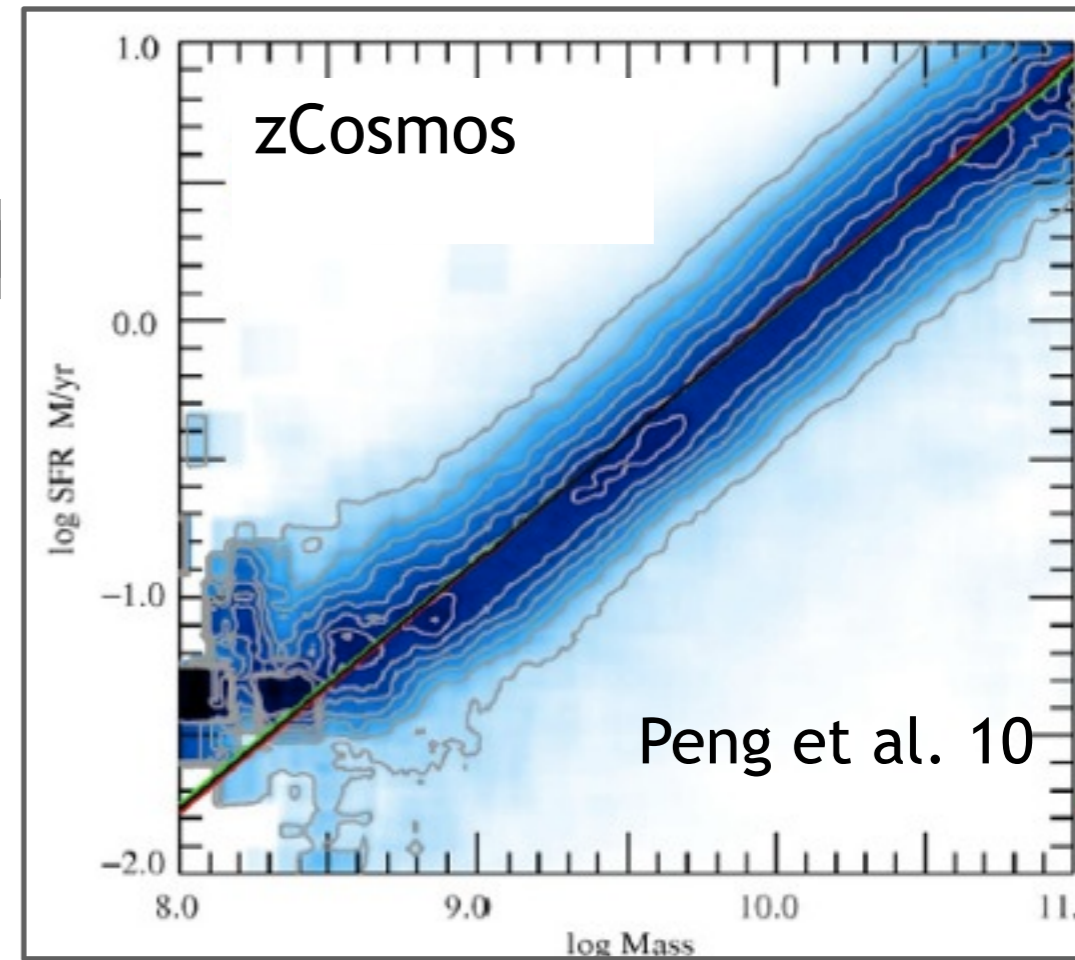
Schneider et al. 2010

# What determines the SFR?

$$\log \text{SFR} \left[ M_{\odot} / \text{yr} \right]$$

**Star formation main sequence** (Noeske et al. 07; Daddi et al. 07, Peng et al. 10, Bouche et al. 10):

$$\text{SFR} \approx 6 \left( \frac{M_*}{10^{11} M_{\odot}} \right)^{0.8..1} (1+z)^{2.5} \frac{M_{\odot}}{\text{yr}}$$



**Cosmic baryonic accretion rate** (Neistein & Dekel 08):  $\log M_* [M_{\odot}]$

$$\left( \frac{dM_g}{dt} \right)_{acc} \approx 7 \cdot \epsilon_g \left( \frac{M_{DM}}{10^{12} M_{\odot}} \right)^{1.1} (1+z)^{2.2} \frac{M_{\odot}}{\text{yr}}$$